# Self-similar solutions of the $p$-Laplace heat equation: the case $p > 2$.

Marie Françoise Bidaut-Véron*

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**Abstract**

We study the self-similar solutions of the equation

$$u_t - div(|\nabla u|^{p-2} \nabla u) = 0,$$

in $\mathbb{R}^N$, when $p > 2$. We make a complete study of the existence and possible uniqueness of solutions of the form

$$u(x,t) = (\pm t)^{-\alpha/\beta} w((\pm t)^{-1/\beta} |x|)$$

of any sign, regular or singular at $x = 0$. Among them we find solutions with an expanding compact support or a shrinking hole (for $t > 0$), or a spreading compact support or a focussing hole (for $t < 0$). When $t < 0$, we show the existence of positive solutions oscillating around the particular solution $U(x,t) = C_{N,p}(|x|^p / (-t))^{1/(p-2)}$.

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*Laboratoire de Mathématiques et Physique Théorique, CNRS UMR 6083, Faculté des Sciences, Parc Grandmont, 37200 Tours, France. e-mail:veronmf@univ-tours.fr

## 1 Introduction and main results

Here we consider the self-similar solutions of the degenerate heat equation involving the $p$-Laplace operator

$$u_t - div(|\nabla u|^{p-2} \nabla u) = 0. \qquad (\mathbf{E}_u)$$

in $\mathbb{R}^N$, with $p > 2$. This study is the continuation of the work started in [4], relative to the case $p < 2$. It can be read independently. We set

$$\gamma = \frac{p}{p-2}, \qquad \eta = \frac{N-p}{p-1}, \qquad (1.1)$$

thus $\gamma > 1$, $\eta < N$,

$$\frac{N+\gamma}{p-1} = \eta + \gamma = \frac{N-\eta}{p-2}. \qquad (1.2)$$

If $u$ is a solution, then for any $\alpha, \beta \in \mathbb{R}$, $u_\lambda(x,t) = \lambda^\alpha u(\lambda x, \lambda^\beta t)$ is a solution of $(\mathbf{E}_u)$ if and only if

$$\beta = \alpha(p-2) + p = (p-2)(\alpha + \gamma); \qquad (1.3)$$

notice that $\beta > 0 \Longleftrightarrow \alpha > -\gamma$. Given $\alpha \in \mathbb{R}$ such that $\alpha \neq -\gamma$, we search self-similar solutions, radially symmetric in $x$, of the form:

$$u = u(x,t) = (\varepsilon\beta t)^{-\alpha/\beta} w(r), \qquad r = (\varepsilon\beta t)^{-1/\beta} |x|, \qquad (1.4)$$

where $\varepsilon = \pm 1$. By translation, for any real $T$, we obtain solutions defined for any $t > T$ when $\varepsilon\beta > 0$, or $t < T$ when $\varepsilon\beta < 0$. We are lead to the equation

$$\left(|w'|^{p-2} w'\right)' + \frac{N-1}{r} |w'|^{p-2} w' + \varepsilon(rw' + \alpha w) = 0 \qquad \text{in } (0,\infty). \qquad (\mathbf{E}_w)$$

Our purpose is to give a complete description of all the solutions, with constant or changing sign. Equation $(\mathbf{E}_w)$ is very interesting, because it is singular at any zero of $w'$, since $p > 2$, implying a nonuniqueness phenomena.

For example, concerning the constant sign solutions near the origin, it can happen that

$$\lim_{r\to 0} w = a \neq 0, \qquad \lim_{r\to 0} w' = 0,$$

we will say that $w$ is *regular*, or

$$\lim_{r\to 0} w = \lim_{r\to 0} w' = 0,$$

we say that $w$ is *flat*. Or different kinds of singularities may occur, either at the level of $w$ :

$$\lim_{r\to 0} w = \infty,$$

or at the level of the gradient:

$$\begin{array}{lll} \lim_{r\to 0} w = a \in \mathbb{R}, & \lim_{r\to 0} w' = \pm\infty, & \text{when } p > N > 1, \\ \lim_{r\to 0} w = a \in \mathbb{R}, & \lim_{r\to 0} w' = b \neq 0 & \text{when } p > N = 1. \end{array}$$

We first show that any local solution $w$ of $(\mathbf{E}_w)$ can be defined on $(0, \infty)$, thus any solution $u$ of equation $(\mathbf{E}_u)$ associated to $w$ by (1.4) is defined on $\mathbb{R}^N \backslash \{0\} \times (0, \pm\infty)$. Then we prove the existence of regular solutions, flat ones, and of all singular solutions mentioned above.

Moreover, for $\varepsilon = 1$, there exist solutions $w$ with a compact support $(0, \bar{r})$; then $u \equiv 0$ on the set

$$D = \left\{(x,t) : x \in \mathbb{R}^N, \quad \beta t > 0, \quad |x| > (\beta t)^{1/\beta} \bar{r}\right\}.$$

For $\varepsilon = -1$, there exist solutions with a hole: $w(r) = 0 \Longleftrightarrow r \in (0, \bar{r})$. Then $u \equiv 0$ on the set

$$H = \left\{(x,t) : x \in \mathbb{R}^N, \quad \beta t < 0, \quad |x| < (-\beta t)^{1/\beta} \bar{r}\right\}.$$

The free boundary is of parabolic type for $\beta > 0$, of hyperbolic type for $\beta < 0$. This leads to four types of solutions, and we prove their existence:

• If $t > 0$, with $\varepsilon = 1, \beta > 0$, we say that $u$ has an *expanding support*; the support increases from $\{0\}$ as $t$ increases from 0.

• If $t > 0$, with $\varepsilon = -1, \beta < 0$, we say that $u$ has a *shrinking hole*: the hole decreases from infinity as $t$ increases from 0;

• If $t < 0$, with $\varepsilon = 1, \beta < 0$, we say that $u$ has a *spreading support*: the support increases to be infinite as $t$ increases to 0.

• If $t < 0$, with $\varepsilon = -1, \beta > 0$, we say that $u$ has a *focussing hole*: the hole disappears as $t$ increases to 0.

Up to our knowledge, some of them seem completely new, as for example the solutions with a shrinking hole or a spreading support. In particular we find again and improve some results of [8] concerning the existence of focussing type solutions.

Finally for $t < 0$ we also show the existence of positive solutions turning around the fundamental solution $U$ given at (1.8) with a kind of periodicity, and also the existence of changing sign solutions doubly oscillating in $|x|$ near 0 and infinity.

As in [4] we reduce the problem to dynamical systems.

When $\varepsilon = -1$, a critical negative value of $\alpha$ is involved:

$$\alpha^* = -\gamma + \frac{\gamma(N+\gamma)}{(p-1)(N+2\gamma)}. \tag{1.5}$$

## 1.1 Explicit solutions

Obviously if $w$ is a solution of $(\mathbf{E}_w)$, $-w$ is also a solution. Some particular solutions are well-known.

**The solution** $U$. For any $\alpha$ such that $\varepsilon(\alpha+\gamma) < 0$, that means $\varepsilon\beta < 0$, there exist flat solutions of $(\mathbf{E}_w)$, given by

$$w(r) = \pm \ell r^{\gamma}, \tag{1.6}$$

where

$$\ell = \left(\frac{|\alpha+\gamma|}{\gamma^{p-1}(\gamma+N)}\right)^{1/(p-2)} > 0. \tag{1.7}$$

They correspond to a unique solution of $(\mathbf{E}_u)$ called $U$, defined for $t < 0$, such that $U(0,t) = 0$, flat, blowing up at $t = 0$ for fixed $x \neq 0$ :

$$U(x,t) = C\left(\frac{|x|^p}{-t}\right)^{1/(p-2)}, \qquad C = ((p-2)\gamma^{p-1}(\gamma+N))^{1/(2-p)}. \tag{1.8}$$

**The case** $\alpha = N$. Then $\beta = \beta_N = N(p-2)+p > 0$, and the equation has a first integral

$$w + \varepsilon r^{-1}\left|w'\right|^{p-2} w' = Cr^{-N}. \tag{1.9}$$

All the solutions corresponding to $C = 0$ are given by

$$\begin{aligned} w &= w_{K,\varepsilon}(r) = \pm\left(K - \varepsilon\gamma^{-1} r^{p'}\right)_+^{(p-1)/(p-2)}, \qquad K \in \mathbb{R}, \\ u &= \pm u_{K,\varepsilon}(x,t) = \pm(\varepsilon\beta_N t)^{-N/\beta_N}\left(K - \varepsilon\gamma^{-1}(\varepsilon\beta_N t)^{-p'/\beta_N}\,|x|^{p'}\right)_+^{(p-1)/(p-2)}. \end{aligned} \tag{1.10}$$

For $\varepsilon = 1$, $K > 0$, they are defined for $t > 0$, called *Barenblatt solutions,* regular with a compact support. Given $c > 0$, the function $u_{K,1}$, defined on $\mathbb{R}^N \times (0,\infty)$, is the unique solution of equation $(\mathbf{E}_u)$ with initial data $u(0) = c\delta_0$, where $\delta_0$ is the Dirac mass at 0, and $K$ being linked by $\int_{\mathbb{R}^N} u_K(x,t)dt = c$. The $u_{K,1}$ are the only nonnegative solutions defined on $\mathbb{R}^N \times (0,\infty)$, such that $u(x,0) = 0$ for any $x \neq 0$. For $\varepsilon = -1$, the $u_{K,-1}$ are defined for $t < 0$; for $K > 0$, $w$ does not vanish on $(0,\infty)$ ; for $K<0$, $w$ is flat with a hole near 0. For $K=0$, we find again the function $w$ given at (1.6).

**The case** $\alpha=\eta\neq0$. We exhibit a family of solutions of $(\mathbf{E}_w)$ :

$$w(r)=Cr^{-\eta},\qquad u(t,x)=C\left|x\right|^{-\eta},\qquad C\neq0. \tag{1.11}$$

The solutions $u$, independent of $t$, are $p$-harmonic in $\mathbb{R}^N$; they are fundamental solutions when $p<N$. When $p>N$, $w$ satisfies $\lim_{r\to0}w=0$, and $\lim_{r\to0}w'=\infty$ for $N>1$, $\lim_{r\to0}w'=b$ for $N=1$.

**The case** $\alpha=-p'$. Equation $(\mathbf{E}_w)$ admits regular solutions of the form

$$w(r)=\pm K\left(N(Kp')^{p-2}+\varepsilon r^{p'}\right),\qquad u(x,t)=\pm K\left(N(Kp')^{p-2}t+|x|^{p'}\right),\qquad K>0. \tag{1.12}$$

Here $\beta>0$; in the two cases $\varepsilon=1,t>0$ and $\varepsilon=-1,t<0$, $u$ is defined for any $t\in\mathbb{R}$ and of the form $\psi(t)+\Phi(|x|)$ with $\Phi$ nonconstant, and $u(.,t)$ has a constant sign for $t>0$ and changing sign for $t<0$.

**The case** $\alpha=0$. Equation $(\mathbf{E}_w)$ can be explicitely solved: either $w'\equiv0$, thus $w\equiv a\in\mathbb{R}$, $u$ is a constant solution of $(\mathbf{E}_u)$, or there exists $K\in\mathbb{R}$ such that

$$\left|w'\right|=r^{-(\eta+1)}\left(K-\frac{\varepsilon}{\gamma+N}r^{N-\eta}\right)_+^{1/(p-2)}; \tag{1.13}$$

and $w$ follows by integration, up to a constant, and then $u(x,t)=w(|x|/(\varepsilon pt)^{1/p})$. If $\varepsilon=1$, then $t>0$, $K>0$ and $w'$ has a compact support; up to a constant, $u$ has a compact support. If $\varepsilon=-1$, then $t<0$; for $K>0$, $w$ is strictly monotone; for $K<0$, $w$ is flat, constant near 0; for $K=0$, we find again (1.6). For $\varepsilon=\pm1,K>0$, observe that $\lim_{r\to0}w=\pm\infty$ if $p\leqq N$; and $\lim_{r\to0}w=a\in\mathbb{R}$, $\lim_{r\to0}w'=\pm\infty$ if $p>N>1$; and $\lim_{r\to0}w=a\in\mathbb{R}$, $\lim_{r\to0}w'=K$ if $p>N=1$. In particular we find solutions such that $w=cr^{|\eta|}(1+o(1))$ near 0, with $c>0$.

**(v) Case** $N=1$ **and** $\alpha=-(p-1)/(p-2)<0$. Here $\beta=1$, and we find the solutions

$$w(r)=\pm\left(Kr+\varepsilon\left|\alpha\right|^{p-1}\left|K\right|^{p}\right)_+^{(p-1)/(p-2)},\quad u(x,t)=\pm\left(K\left|x\right|+\left|\alpha\right|^{p-1}\left|K\right|^{p}t\right)_+^{(p-1)/(p-2)}, \tag{1.14}$$

If $\varepsilon=1,t>0$, then $w$ has a singularity at the level of the gradient, and either $K>0$, $w>0$, or $K<0$ and $w$ has a compact support. If $\varepsilon=-1,t<0$ then $K>0$, $w$ has a hole.

## 1.2 Main results

In the next sections we provide an exhaustive study of equation $(\mathbf{E}_w)$. Here we give the main results relative to the function $u$. Let us show how to return from $w$ to $u$. Suppose that the behaviour of $w$ is given by

$$\lim_{r\to0}r^{\lambda}w(r)=c\neq0,\qquad\lim_{r\to\infty}r^{\mu}w(r)=c'\neq0,\qquad\text{where }\lambda,\mu\in\mathbb{R}.$$

(i) Then for fixed $t \neq 0$, the function $u$ has a behaviour in $|x|^{-\lambda}$ near $x = 0$, and a behaviour in $|x|^{-\mu}$ for large $|x|$.

If $\lambda = 0$, then $u$ is defined on $\mathbb{R}^N \times (0, \pm\infty)$. Either $w$ is regular, then $u(.,t) \in C^1\left(\mathbb{R}^N \times (0,\infty)\right)$; we will say that $u$ is **regular**; nevertheless the regular solutions $u$ presents a singularity at time $t = 0$ if and only if $\alpha < -\gamma$ or $\alpha > 0$. Or a singularity can appear for $u$ at the level of the gradient.

If $\lambda < 0$, thus $u$ is defined on $\mathbb{R}^N \times (0, \pm\infty)$ and $u(0,t) = 0$; either $w$ is flat, we also say that $u$ is **flat**, or a singularity appears at the level of the gradient.

If $0 < \lambda < N$, then $u(.,t) \in L^1_{loc}\left(\mathbb{R}^N\right)$ for $t \neq 0$, we say that $x = 0$ is a **weak** *singularity.* We will show that there exist no stronger singularity.

If $\lambda < N < \mu$; then $u(.,t) \in L^1\left(\mathbb{R}^N\right)$.

(ii) For fixed $x \neq 0$, the behaviour of $u$ near $t = 0$, depends on the sign of $\beta$:

$$\lim_{t\to 0} |x|^{\mu} |t|^{(\alpha-\mu)/\beta} u(x,t) = C \neq 0 \quad \text{if} \quad \alpha > -\gamma,$$
$$\lim_{t\to 0} |x|^{\lambda} |t|^{(\alpha-\lambda)/\beta} u(x,t) = C \neq 0 \quad \text{if} \quad \alpha < -\gamma.$$

If $\mu < 0, \alpha > -\gamma$ or $\lambda < 0, \alpha < -\gamma$, then $\lim_{t\to 0} u(x,t) = 0$.

### 1.2.1 Solutions defined for $t > 0$

Here we look for solutions $u$ of $(\mathbf{E}_u)$ of the form (1.4) defined on $\mathbb{R}^N \backslash \{0\} \times (0,\infty)$. That means $\varepsilon\beta > 0$ or equivalently $\varepsilon = 1$, $-\gamma < \alpha$ (see Section 6) or $\varepsilon = -1, \alpha < -\gamma$ see (Section 7). We begin by the case $\varepsilon = 1$, treated at Theorem 6.1.

**Theorem 1.1** *Assume $\varepsilon = 1$, and $-\gamma < \alpha$.*

*(1) Let $\alpha < N$.*

*All regular solutions on $\mathbb{R}^N \backslash \{0\} \times (0,\infty)$ have a strict constant sign, in $|x|^{-\alpha}$ near $\infty$ for fixed $t$, with initial data $L|x|^{-\alpha}$ $(L \neq 0)$ in $\mathbb{R}^N$; thus $u(.,t) \notin L^1\left(\mathbb{R}^N\right)$, and $u$ is unbounded when $\alpha < 0$.*

*There exist nonnegative solutions such that near $x = 0$,*

$$\left.\begin{array}{lll} \text{for } p < N, & u \text{ has a weak singularity in } |x|^{-\eta}, & \\ \text{for } p = N, & u \text{ has a weak singularity in } \ln|x|, & \\ \text{for } p > N, & u \in C^0(\mathbb{R}^N \times (0,\infty)), & u(0,t) = a > 0, \text{ with a singular gradient,} \end{array}\right\} \tag{1.15}$$

*and $u$ has an **expanding compact support** for any $t > 0$, with initial data $L|x|^{-\alpha}$ in $\mathbb{R}^N \backslash \{0\}$.*

*There exist positive solutions with the same behaviour as $x \to 0$, in $|x|^{-\alpha}$ near $\infty$ for fixed $t$; and also solutions such that $u$ has one zero for fixed $t \neq 0$, and the same behaviour.*

*If $p > N$, there exist positive solutions satisfying (1.15), and also positive solutions such that*

$$u \in C^0(\mathbb{R}^N \times (0,\infty)), \quad u(0,t) = 0, \text{ in } |x|^{|\eta|} \text{ near } 0, \text{ with a singular gradient,} \tag{1.16}$$

*in $|x|^{-\alpha}$ near $\infty$ for fixed $t$, with and initial data $L\,|x|^{-\alpha}$ in $\mathbb{R}^N \backslash \{0\}$.*

*(2) Let $\alpha = N$.*

*All **regular (Barenblatt)** solutions are nonnegative, have a **compact support** for any $t > 0$. If $p \leqq N$, all the other solutions have one zero for fixed $t$, satisfy (1.15) or (1.16) and have the same behaviour at $\infty$.*

*(3) Let $N < \alpha$.*

*All regular solutions $u$ have a finite number $m \geqq 1$ of simple zeros for fixed $t$, and $u(.,t) \in L^1(\mathbb{R}^N)$. Either they are in $|x|^{-\alpha}$ near $\infty$ for fixed $t$, then there exist solutions with $m$ zeros, compact support, satisfying (1.15); or they have a compact support. All the solutions have $m$ or $m+1$ zeros. There exist solutions satisfying (1.15) with $m+1$ zeros, and in $|x|^{-\alpha}$ near $\infty$. If $p > N$, there exist solutions satisfying (1.15) with $m$ zeros; there exist also solutions with $m$ zeros, $u(0,t) = 0$, and a singular gradient, in $|x|^{-\alpha}$ near $\infty$.*

Next we come to the case $\varepsilon = -1$, which is the subject of Theorem 7.1.

**Theorem 1.2** *Assume $\varepsilon = -1$ and $\alpha < -\gamma$.*

*All the solutions $u$ on $\mathbb{R}^N \backslash \{0\} \times (0,\infty)$, in particular the regular ones, are **oscillating around** 0 for fixed $t > 0$ and large $|x|$, and $r^{-\gamma}w$ is asymptotically periodic in $\ln r$. Moreover there exist*

*solutions such that $r^{-\gamma}w$ is **periodic** in $\ln r$, in particular $C_1 t^{-|\alpha/\beta|} \leqq |u| \leqq C_2 t^{-|\alpha/\beta|}$ for some $C_1, C_2 > 0$;*

*solutions $u \in C^1(\mathbb{R}^N \times [0,\infty))$, $u(x,0) \equiv 0$, with a **shrinking hole**;*

***flat** solutions $u \in C^1(\mathbb{R}^N \times [0,\infty))$, in $|x|^{|\alpha|}$ near 0, with initial data $L\,|x|^{|\alpha|}$ $(L \neq 0)$;*

*solutions satisfying (1.15) near $x = 0$, and if $p > N$, solutions satisfying (1.16) near 0.*

### 1.2.2 Solutions defined for $t < 0$

We look for solutions $u$ of $(\mathbf{E}_u)$ of the form (1.4) defined on $\mathbb{R}^N \backslash \{0\} \times (-\infty, 0)$. That means $\varepsilon\beta < 0$ or equivalently $\varepsilon = 1$, $\alpha < -\gamma$ (see Section 8, Theorem 8.1) or $\varepsilon = -1, \alpha > -\gamma$ (see Section 9). In the case $\varepsilon = 1$, we get the following:

**Theorem 1.3** *Assume $\varepsilon = 1$, and $\alpha < -\gamma$.*

*The function $U(x,t) = C\left(\frac{|x|^p}{-t}\right)^{1/(p-2)}$ is a positive **flat** solution on $\mathbb{R}^N \backslash \{0\} \times (-\infty, 0)$.*

*All regular solutions have a constant sign, are unbounded in $|x|^\gamma$ near $\infty$ for fixed $t$, and blow up at $t = 0$ like $(-t)^{-|\alpha|/|\beta|}$ for fixed $x \neq 0$.*

*There exist **flat positive** solutions $u \in C^1(\mathbb{R}^N \times (-\infty, 0])$, in $|x|^{\gamma}$ near $\infty$ for fixed $t$, with **final data** $L\,|x|^{|\alpha|}$ $(L > 0)$.*

*There exist **nonnegative** solutions satisfying (1.15) near $0$, with a **spreading compact support**, blowing up near $t = 0$ (like $|t|^{-(\eta+|\alpha|)/|\beta|}$ for $p < N$, or $|t|^{-|\alpha|/|\beta|} \ln|t|$ for $p = N$, or $(-t)^{-|\alpha|/|\beta|}$for $p > N$).*

*There exist positive solutions with the same behaviour near $0$, in $|x|^{\gamma}$ near $\infty$, blowing up as above at $t = 0$, and solutions with one zero for fixed $t$, and the same behaviour. If $p > N$, there exist positive solutions satisfying (1.15) (resp. (1.16)) near 0, in $|x|^{\gamma}$ near $\infty$ for fixed $t$, blowing up at $t = 0$ like $|t|^{-|\alpha|/|\beta|}$ ( resp. $|t|^{(|\eta|-|\alpha|)/|\beta|}$) for fixed $x$.*

*Up to a symmetry, all the solutions are described.*

The most interesting case is $\varepsilon = -1, -\gamma < \alpha$. For simplicity we will assume that $p < N$. The case $p \geqq N$ is much more delicate, and the complete results can be read in terms of $w$ at Theorems 9.4, 9.6, 9.9, 9.10, 9.11 and 9.12. We discuss according to the position of $\alpha$ with respect to $-p'$ and $\alpha^*$ defined at (1.5). Notice that $\alpha^* < -p'$.

**Theorem 1.4** *Assume $\varepsilon = -1$, and $-p' \leqq \alpha \neq 0$. The function $U$ is still a flat solution on $\mathbb{R}^N \setminus \{0\} \times (-\infty, 0)$.*

*(1) Let $0 < \alpha$.*

*All regular solutions have a strict constant sign, in $|x|^{\gamma}$ near $\infty$ for fixed $t$, blowing up at $t = 0$ like $(-t)^{-1/(p-2)}$ for fixed $x \neq 0$.*

*There exist nonnegative solutions with a **focussing hole**: $u(x,t) \equiv 0$ for $|x| \leqq C\,|t|^{1/\beta}$, $t > 0$, in $|x|^{\gamma}$ near $\infty$ for fixed $t$, blowing up at $t = 0$ like $(-t)^{-1/(p-2)}$ for fixed $x \neq 0$.*

*There exist positive solutions $u$ with a **(weak) singularity** in $|x|^{-\eta}$at $x = 0$, in $|x|^{-\alpha}$ near $\infty$ for fixed $t$, with $u(.,t) \in L^1(\mathbb{R}^N)$ if $\alpha > N$, with final data $L\,|x|^{-\alpha}$ $(L > 0)$ in $\mathbb{R}^N \setminus \{0\}$.*

*There exist positive solutions $u$ in $|x|^{-\eta}$at $x = 0$, in $|x|^{\gamma}$ near $\infty$ for fixed $t$, blowing up at $t = 0$ like $(-t)^{-1/(p-2)}$ for fixed $x \neq 0$; solutions with one zero and the same behaviour.*

*(2) Let $-p' < \alpha < 0$.*

*All regular solutions have **one zero** for fixed $t$, and the same behaviour. There exist solutions with one zero, in $|x|^{-\eta}$at $x = 0$, in $|x|^{|\alpha|}$ near $\infty$ for fixed $t$, with final data $L\,|x|^{-\alpha}$ $(L > 0)$ in $\mathbb{R}^N \setminus \{0\}$. There exist solutions with one zero, $u$ in $|x|^{-\eta}$at $x = 0$, in $|x|^{\gamma}$ near $\infty$ for fixed $t$, blowing up at $t = 0$ like $(-t)^{-1/(p-2)}$ for fixed $x \neq 0$; solutions with two zeros and the same behaviour.*

*3) Let $\alpha = -p'$.*

*All regular solutions have **one zero** and are in $|x|^{|\alpha|}$ near $\infty$ for fixed $t$, and with **final data** $L\,|x|^{|\alpha|}$ $(L > 0)$. The other solutions have one or two zeros, are in $|x|^{-\eta}$at $x = 0$, in $|x|^{\gamma}$ near $\infty$ for fixed $t$.*

*In any case, up to a symmetry, all the solutions are described.*

**Theorem 1.5** *Assume* $\varepsilon = -1, -\gamma < \alpha < -p'$. *Then* $U$ *is still a flat solution on* $\mathbb{R}^N \backslash \{0\} \times (-\infty, 0)$.

*(1)Let* $\alpha \leqq \alpha^*$.

*Then there exist* ***positive flat solutions****, in* $|x|^\gamma$ *near 0, in* $|x|^{|\alpha|}$ *near* $\infty$ *for fixed* $t$, *with* ***final data*** $L\,|x|^{-\alpha}$ $(L > 0)$ *in* $\mathbb{R}^N$.

*All the other solutions, among them the* ***regular ones****, have an* ***infinity of zeros****:* $u(t,.)$ *is oscillating around* $0$ *for large* $|x|$. *There exist solutions with a focussing hole, and solutions with a singularity in* $|x|^{-\eta}$ *at* $x = 0$. *There exist solutions* ***oscillating also for small*** $|x|$, *such that* $r^{-\gamma}w$ *is periodic in* $\ln r$.

*(2) There exist a* ***critical unique value*** $\alpha_c \in (\max(\alpha^*, -p')$ *such that for* $\alpha = \alpha_c$, *there exists nonnegative solutions with a* ***focussing hole*** *near 0, in* $|x|^{|\alpha|}$ *near* $\infty$ *for fixed* $t$, *with* ***final data*** $L\,|x|^{-\alpha}$ $(L > 0)$ *in* $\mathbb{R}^N$. *And* $\alpha_c > -(p-1)/(p-2)$.
*There exist positive flat solutions, such that* $|x|^{-\gamma}\,u$ *is bounded on* $\mathbb{R}^N$ *for fixed* $t$, *blowing up at* $t = 0$ *like* $(-t)^{-1/(p-2)}$ *for fixed* $x \neq 0$. *The regular solutions are oscillating around* $0$ *as above. There exist solutions* ***oscillating around*** $0$, *such that* $r^{-\gamma}w$ *is* ***periodic*** *in* $\ln r$. *There are solutions with a weak singularity in* $|x|^{-\eta}$ *at* $x = 0$, *and oscillating around 0 for large* $|x|$.

*(3) Let* $\alpha^* < \alpha < \alpha_c$.

*The regular solutions are as above. There exist solutions of the same types as above. Moreover there exist* ***positive*** *solutions, such that* $r^{-\gamma}w$ *is* ***periodic*** *in* $\ln r$, *thus there exist* $C_1, C_2 > 0$ *such that*

$$C_1 \left(\frac{|x|^p}{|t|}\right)^{1/(p-2)} \leqq u \leqq C_2 \left(\frac{|x|^p}{|t|}\right)^{1/(p-2)}$$

*There exist* ***positive*** *solutions, such that* $r^{-\gamma}w$ *is asymptotically periodic in* $\ln r$ *near* $0$ *and in* $|x|^\gamma$ *near* $\infty$ *for fixed* $t$; *and also, solutions with a hole, and oscillating around 0 for large* $|x|$. *There exist solutions positive near* $0$, *oscillating near* $\infty$, *and* $r^{-\gamma}w$ *is* ***doubly asymptotically periodic*** *in* $\ln r$.

*4) Let* $\alpha_c < \alpha < -p'$.

*There exist nonnegative solutions with a focussing hole near 0, in* $|x|^\gamma$ *near* $\infty$ *for fixed* $t$, *blowing up at* $t = 0$ *like* $(-t)^{-1/(p-2)}$ *for fixed* $x \neq 0$. *Either the regular solutions have an* ***infinity*** *of zeros for fixed* $t$, *then the same is true for all the other solutions. Or they have a* ***finite*** *number* $m \geqq 2$ *of zeros, and can be in* $|x|^\gamma$ *or* $|x|^{|\alpha|}$ *near* $\infty$ *(in that case they have a final data* $L\,|x|^{|\alpha|}$*); all the other solutions have* $m$ *or* $m+1$ *zeros.*

In the case $\alpha = \alpha_c$, we find again the existence and uniqueness of the focussing solutions introduced in [8].

## 2 Different formulations of the problem

In all the sequel we assume

$$\alpha \neq 0,$$

recalling that the solutions $w$ are given explicitely by (1.13) when $\alpha = 0$. Defining

$$J_N(r) = r^N \left( w + \varepsilon r^{-1} \left| w' \right|^{p-2} w' \right), \qquad J_\alpha(r) = r^{\alpha - N} J_N(r), \tag{2.1}$$

equation $(\mathbf{E}_w)$ can be written in an equivalent way under the forms

$$J'_N(r) = r^{N-1}(N-\alpha)w, \qquad J'_\alpha(r) = -\varepsilon(N-\alpha)r^{\alpha-2} \left| w' \right|^{p-2} w'. \tag{2.2}$$

If $\alpha = N$, then $J_N$ is constant, so we find again (1.9).

We mainly use logarithmic substitutions; given $d \in \mathbb{R}$, setting

$$w(r) = r^{-d} y_d(\tau), \qquad Y_d = -r^{(d+1)(p-1)} \left| w' \right|^{p-2} w', \qquad \tau = \ln r, \tag{2.3}$$

we obtain the equivalent system:

$$\left.\begin{aligned} y'_d &= d y_d - |Y_d|^{(2-p)/(p-1)} Y_d, \\ Y'_d &= (p-1)(d-\eta)Y_d + \varepsilon e^{(p+(p-2)d)\tau}(\alpha y_d - |Y_d|^{(2-p)/(p-1)} Y_d). \end{aligned}\right\} \tag{2.4}$$

At any point $\tau$ where $w'(\tau) \neq 0$, the functions $y_d, Y_d$ satisfy the equations

$$y''_d + (\eta - 2d)y'_d - d(\eta - d)y_d + \frac{\varepsilon}{p-1} e^{((p-2)d+p)\tau} \left| d y_d - y'_d \right|^{2-p} \left( y'_d + (\alpha - d) y_d \right) = 0, \tag{2.5}$$

$$\begin{aligned} Y''_d + (p-1)(\eta - 2d - p')Y'_d + \varepsilon e^{((p-2)d+p)\tau} |Y_d|^{(2-p)/(p-1)} \left( Y'_d/(p-1) + (\alpha - d)Y_d \right) \\ -(p-1)^2(\eta - d)(p' + d)Y_d = 0, \end{aligned} \tag{2.6}$$

The main case is $d = -\gamma$: setting $y = y_{-\gamma}$,

$$w(r) = r^\gamma y(\tau), \qquad Y = -r^{(-\gamma+1)(p-1)} \left| w' \right|^{p-2} w', \qquad \tau = \ln r, \tag{2.7}$$

we are lead to the *autonomous* system

$$\left.\begin{aligned} y' &= -\gamma y - |Y|^{(2-p)/(p-1)} Y, \\ Y' &= -(\gamma + N)Y + \varepsilon(\alpha y - |Y|^{(2-p)/(p-1)} Y). \end{aligned}\right\} \tag{$\mathbf{S}$}$$

Its study is fundamental: its phase portrait allows to study all the *signed* solutions of equation $(\mathbf{E}_w)$. Equation (2.5) takes the form

$$(p-1)y''+(N+\gamma p)y'+\gamma(\gamma+N)y+\varepsilon\left|\gamma y+y'\right|^{2-p}\left(y'+(\alpha+\gamma)y\right)=0, \qquad (\mathbf{E}_y)$$

Notice that $J_N(r)=r^{N+\gamma}(y(\tau)-\varepsilon Y(\tau))$.

**Remark 2.1** *Since (**S**) is autonomous, for any solution $w$ of $(\mathbf{E}_w)$ of the problem, all the functions $w_\xi(r)=\xi^{-\gamma}w(\xi r),\xi>0$, are also solutions.*

**Notation 2.2** *In the sequel we set $\varepsilon\infty:=+\infty$ if $\varepsilon=1$, $\varepsilon\infty:=-\infty$ if $\varepsilon=-1$.*

## 2.1 The phase plane of system (S)

In the phase plane $(y,Y)$ we denote the four quadrants by

$$\mathcal{Q}_1=(0,\infty)\times(0,\infty), \quad \mathcal{Q}_2=(-\infty,0)\times(0,\infty), \quad \mathcal{Q}_3=-\mathcal{Q}_1, \quad \mathcal{Q}_4=-\mathcal{Q}_2.$$

**Remark 2.3** *The vector field at any point $(0,\xi),\xi>0$ satisfies $y'=-\xi^{1/(p-1)}<0$, thus points to $\mathcal{Q}_2$; moreover $Y'<0$ if $\varepsilon=1$. The field at any point $(\varphi,0),\varphi>0$ satisfies $Y'=\varepsilon\alpha\varphi$, thus points to $\mathcal{Q}_1$ if $\varepsilon\alpha>0$ and to $\mathcal{Q}_4$ if $\varepsilon\alpha<0$; moreover $y'=-\gamma\varphi<0$.*

If $\varepsilon(\gamma+\alpha)\geqq 0$, system (**S**) has a unique stationary point $(0,0)$. If $\varepsilon(\gamma+\alpha)<0$, it admits three stationary points:

$$(0,0), \qquad M_\ell=(\ell,-(\gamma\ell)^{p-1})\in\mathcal{Q}_4, \qquad M'_\ell=-M_\ell\in\mathcal{Q}_2, \tag{2.8}$$

where $\ell$ is defined at (1.7). The point $(0,0)$ is singular because $p>2$; its study concern in particular the solutions $w$ with a *double zero.* When $\varepsilon(\gamma+\alpha)<0$, the point $M_\ell$ is associated to the solution $w\equiv\ell r^\gamma$ of equation $(\mathbf{E}_w)$ given at (1.1).

**Linearization around** $M_\ell$. Near the point $M_\ell$, setting

$$y=\ell+\overline{y}, \qquad Y=-(\gamma\ell)^{p-1}+\overline{Y}, \tag{2.9}$$

system (**S**) is equivalent in $\mathcal{Q}_4$ to

$$\overline{y}'=-\gamma\overline{y}-\varepsilon\nu(\alpha)\overline{Y}+\Psi(\overline{Y}), \qquad \overline{Y}'=\varepsilon\alpha\overline{y}-(\gamma+N+\nu(\alpha))\overline{Y}+\varepsilon\Psi(\overline{Y}), \tag{2.10}$$

where

$$\nu(\alpha)=-\frac{\gamma(N+\gamma)}{(p-1)(\gamma+\alpha)}, \text{ and } \Psi(\vartheta)=((\gamma\ell)^{p-1}-\vartheta)^{1/(p-1)}-\gamma\ell+\frac{(\gamma\ell)^{2-p}}{p-1}\vartheta, \qquad \vartheta<(\gamma\ell)^{p-1}, \tag{2.11}$$

thus $\varepsilon\nu(\alpha) > 0$. The linearized problem is given by

$$\overline{y}' = -\gamma\overline{y} - \varepsilon\nu(\alpha)\overline{Y}, \qquad \overline{Y}' = \varepsilon\alpha\overline{y} - (\gamma + N + \nu(\alpha))\overline{Y}.$$

Its eigenvalues $\lambda_1 \leqq \lambda_2$ are the solutions of equation

$$\lambda^2 + (2\gamma + N + \nu(\alpha))\lambda + p'(N + \gamma) = 0 \tag{2.12}$$

The discriminant $\Delta$ of the equation (2.12) is given by

$$\Delta = (2\gamma + N + \nu(\alpha))^2 - 4p'(N + \gamma) = (N + \nu(\alpha))^2 - 4\nu(\alpha)\alpha. \tag{2.13}$$

For $\varepsilon = 1$, $M_\ell$ is a *sink*, and a node point, since $\nu(\alpha) > 0$, and $\alpha < 0$, thus $\Delta > 0$. For $\varepsilon = -1$, we have $\nu(\alpha) < 0$; the nature of $M_\ell$ depends on the critical value $\alpha^*$ defined at (1.5); indeed

$$\alpha = \alpha^* \Longleftrightarrow \lambda_1 + \lambda_2 = 0.$$

Then $M_\ell$ is a *sink* when $\alpha > \alpha^*$ and a *source* when $\alpha < \alpha^*$. Moreover $\alpha^*$ corresponds to a spiral point, and $M_\ell$ is a node point when $\Delta \geqq 0$, that means $\alpha \leqq \alpha_1$, or $\gamma > N/2 + \sqrt{p'(N+\gamma)}$ and $\alpha_2 \leqq \alpha$, where

$$\alpha_1 = -\gamma + \frac{\gamma(N+\gamma)}{(p-1)(2\gamma + N + 2(p'(N+\gamma))^{1/2})}, \qquad \alpha_2 = -\gamma + \frac{\gamma(N+\gamma)}{(p-1)(2\gamma + N - 2(p'(N+\gamma))^{1/2})}. \tag{2.14}$$

When $\Delta > 0$, and $\lambda_1 < \lambda_2$, one can choose a basis of eigenvectors

$$e_1 = (-\varepsilon\nu(\alpha), \lambda_1 + \gamma) \quad \text{and} \quad e_2 = (\varepsilon\nu(\alpha), -\gamma - \lambda_2). \tag{2.15}$$

**Remark 2.4** *One verifies that* $\alpha^* < -1$; *and* $\alpha^* < -(p-1)/(p-2)$ *if and only if* $p > N$. *Also* $\alpha_2 \leqq 0$, *and* $\alpha_2 = 0 \Longleftrightarrow N = p/((p-2)^2$; *and* $\alpha_2 > -p' \Longleftrightarrow \gamma^2 - 7\gamma - 8N < 0$, *which is not always true.*

As in [4, Theorem 2.16] we prove that the Hopf bifurcation point is not degenerate, which implies the existence of small cycles near $\alpha^*$.

**Proposition 2.5** *Let* $\varepsilon = -1$, *and* $\alpha = \alpha^* > -\gamma$. *Then* $M_\ell$ *is a weak source. If* $\alpha > \alpha^*$ *and* $\alpha - \alpha^*$ *is small enough, there exists a unique limit cycle in* $\mathcal{Q}_4$, *attracting at* $-\infty$.

## 2.2 Other systems for positive solutions

When $w$ has a constant sign, we define two functions associated to $(y, Y)$ :

$$\zeta(\tau) = \frac{|Y|^{(2-p)/(p-1)}\, Y}{y}(\tau) = -\frac{rw'(r)}{w(r)}, \qquad \sigma(\tau) = \frac{Y}{y}(\tau) = -\frac{|w'(r)|^{p-2}\, w'(r)}{rw(r)}. \tag{2.16}$$

Thus $\zeta$ describes the behaviour of $w'/w$ and $\sigma$ is the slope in the phase plane $(y, Y)$. They satisfy the system

$$\left.\begin{aligned} \zeta' &= \zeta(\zeta-\eta) + \varepsilon\, |\zeta y|^{2-p}\, (\alpha-\zeta)/(p-1) = \zeta(\zeta-\eta+\varepsilon(\alpha-\zeta)/(p-1)\sigma), \\ \sigma' &= \varepsilon(\alpha-N) + \left(|\sigma y|^{(2-p)/(p-1)}\, \sigma - N\right)(\sigma-\varepsilon) = \varepsilon(\alpha-\zeta) + (\zeta-N)\,\sigma. \end{aligned}\right\} \tag{$\mathbf{Q}$}$$

In particular, System ($\mathbf{Q}$) provides a short proof of the local existence and uniqueness of the *regular* solutions: they correspond to its stationary point $(0, \varepsilon\alpha/N)$, see Section 3.1.

Moreover, if $w$ and $w'$ have a strict constant sign, that means in any quadrant $\mathcal{Q}_i$, we can define

$$\psi = \frac{1}{\sigma} = \frac{y}{Y} \tag{2.17}$$

We obtain a new system relative to $(\zeta, \psi)$ :

$$\left.\begin{aligned} \zeta' &= \zeta(\zeta-\eta+\varepsilon(\alpha-\zeta)\psi/(p-1)), \\ \psi' &= \psi(N-\zeta+\varepsilon\,(\zeta-\alpha)\psi)\,. \end{aligned}\right\} \tag{$\mathbf{P}$}$$

We are reduced to a polynomial system, thus with no singularity. System ($\mathbf{P}$) gives the existence of singular solutions when $p > N$, corresponding to its stationary point $(\eta, 0)$, see Section 5.

We will also consider another system in any $\mathcal{Q}_i$ : setting

$$\zeta = -1/g, \qquad \sigma = -s, \qquad d\tau = gsd\nu = |Y|^{(p-2)/(p-1)}\, d\nu, \tag{2.18}$$

we find

$$\left.\begin{aligned} dg/d\nu &= g(s(1+\eta g) + \varepsilon(1+\alpha g)/(p-1)), \\ ds/d\nu &= -s(\varepsilon(1+\alpha g) + (1+Ng)s). \end{aligned}\right\} \tag{$\mathbf{R}$}$$

System ($\mathbf{R}$) allows to get the existence of solutions $w$ with a hole or a compact support, and other solutions, corresponding to its stationary points $(0, -\varepsilon)$ and $(-1/\alpha, 0)$; it provides a complete study of the singular point $(0, 0)$ of system ($\mathbf{S}$), see Sections 3.3, 5; and of the focussing solutions, see Section 9.

**Remark 2.6** *The particular solutions can be found again in the different phase planes, where their trajectories are lines:*

*For $\alpha = N$, the solutions (1.10) correspond to $Y \equiv \varepsilon y$, that means $\sigma \equiv \varepsilon$.*

*For $\alpha = \eta \neq 0$ the solutions (1.11) correspond to $\zeta \equiv \eta$.*

*For $\alpha = -p'$, the solutions (1.12) are given by $\zeta + \varepsilon N \sigma \equiv \alpha$.*

*For $N = 1$, $\alpha = -(p-2)/(p-1)$, the solutions (1.14) satisfy $\alpha g + \varepsilon s \equiv -1$.*

# 3 Global existence

## 3.1 Local existence and uniqueness

**Proposition 3.1** *Let $r_1 > 0$ and $a, b \in \mathbb{R}$. If $(a, b) \neq (0, 0)$, there exists a unique solution $w$ of equation ($\boldsymbol{E}_w$) in a neighborhood $\mathcal{V}$ of $r_1$, such that $w$ and $|w'|^{p-2} w' \in C^1(\mathcal{V})$ and $w(r_1) = a$, $w'(r_1) = b$. It extends on a maximal interval $I$ where $(w(r), w'(r)) \neq (0, 0)$.*

**Proof.** If $b \neq 0$, the Cauchy theorem directly applies to system (**S**). If $b = 0$ the system is a priori singular on the line $\{Y = 0\}$ since $p > 2$. In fact it is only singular at $(0, 0)$. Indeed near any point $(\xi, 0)$ with $\xi \neq 0$, one can take $Y$ as a variable, and

$$\frac{dy}{dY} = F(Y, y), \qquad F(Y, y) := \frac{\gamma y + |Y|^{(2-p)/(p-1)} Y}{(\gamma + N)Y + \varepsilon(|Y|^{(2-p)/(p-1)} Y - \alpha y)},$$

where $F$ is continuous in $Y$ and $C^1$ in $y$, hence local existence and uniqueness hold. ■

**Notation 3.2** *For any point $P_0 = (y_0, Y_0) \in \mathbb{R}^2 \setminus \{(0, 0)\}$, the unique trajectory in the phase plane $(y, Y)$ of system ($\boldsymbol{S}$) going through $P_0$ is denoted by $\mathcal{T}_{[P_0]}$. By symmetry, $\mathcal{T}_{[-P_0]} = -\mathcal{T}_{[P_0]}$.*

Next we show the existence of regular solutions. Our proof is short, based on phase plane portrait, and not on a fixed point method, rather delicate because $p > 2$, see [3].

**Theorem 3.3** *For any $a \in \mathbb{R}$, $a \neq 0$, there exists a unique solution $w = w(., a)$ of equation ($\boldsymbol{E}_w$) in an interval $[0, r_0)$, such that $w$ and $|w'|^{p-2} w' \in C^1([0, r_0))$ and*

$$w(0) = a, \qquad w'(0) = 0; \tag{3.1}$$

*and then $\lim_{r \to 0} |w'|^{p-2} w'/rw = -\varepsilon\alpha/N$. In other words in the phase plane $(y, Y)$ there exists a unique trajectory $\mathcal{T}_r$ such that $\lim_{\tau \to -\infty} y = \infty$, and $\lim_{\tau \to -\infty} Y/y = \varepsilon\alpha/N$.*

**Proof.** We have assumed $\alpha \neq 0$ (when $\alpha = 0, w \equiv a$ from (1.13)). If such a solution $w$ exists, then from (2.1) and (2.2), $J_N'(r) = r^{N-1}(N - \alpha)a(1 + o(1))$ near 0. Thus $J_N(r) = r^{N-1}(1 - \alpha/N)a(1 + o(1))$, hence $\lim_{r \to 0} |w'|^{p-2} w'/rw = -\varepsilon\alpha/N$; in other words, $\lim_{\tau \to -\infty} \sigma = \varepsilon\alpha/N$. And

$\lim_{\tau\to-\infty} y = \infty$, thus $\lim_{\tau\to-\infty} \zeta = 0$, and $\varepsilon\alpha\zeta > 0$ near $-\infty$. Reciprocally consider system (**Q**). The point $(0, \varepsilon\alpha/N)$ is stationary. Setting $\sigma = \varepsilon\alpha/N + \bar{\sigma}$, the linearized system near this point is given by

$$\zeta' = p'\zeta, \qquad \bar{\sigma}' = \varepsilon\zeta(\alpha - N)/N - N\bar{\sigma}.$$

One finds is a saddle point, with eigenvalues $-N$ and $p'$. Then there exists a unique trajectory $\mathcal{T}'_r$ in the phase-plane $(\zeta, \sigma)$ starting at $-\infty$ from $(0, \varepsilon\alpha/N)$ with the slope $\varepsilon(\alpha - N)/N(N + p') \neq 0$ and $\varepsilon\alpha\zeta > 0$. It corresponds to a unique trajectory $\mathcal{T}_r$ in the phase plane $(y, Y)$, and $\lim_{\tau\to-\infty} y = \infty$, since $y = |\sigma| \, |\zeta|^{1-p})^{1/(p-2)}$. For any solution $(\zeta, \sigma)$ describing $\mathcal{T}'_r$, the function $w(r) = r^{\gamma}(|\sigma| \, |\zeta|^{1-p} (\tau))^{1/(p-2)}$ satisfies $\lim_{r\to0} |w'|^{p-2} w'/rw = -\varepsilon\alpha/N$. As a consequence, $w^{(p-2)/(p-1)}$ has a finite nonzero limit, and $\lim_{r\to0} w' = 0$; thus $w$ is regular. Local existence and uniqueness follows for any $a \neq 0$, by Remark 2.1. ■

**Definition 3.4** *The trajectory $\mathcal{T}_r$ in the plane $(y, Y)$ and its opposite $-\mathcal{T}_r$ will be called regular trajectories. We shall say that $y$ is regular. Observe that $\mathcal{T}_r$ starts in $\mathcal{Q}_1$ if $\varepsilon\alpha > 0$, and in $\mathcal{Q}_4$ if $\varepsilon\alpha < 0$.*

**Remark 3.5** *From Theorem 3.3 and Remark 2.1, all regular solutions are obtained from one one of them: $w(r, a) = aw(a^{-1/\gamma}r, 1)$. Thus they have the same behaviour near $\infty$.*

### 3.2 Sign properties

Next we give informations on the zeros of $w$ or $w'$, by using the monotonicity properties of the functions $y_d, Y_d$, in particular $y, Y$, and $\zeta$ and $\sigma$. At any extremal point $\tau$, they satisfy respectively

$$y''_d(\tau) = y_d(\tau) \left( d(\eta - d) + \frac{\varepsilon(d-\alpha)}{p-1} e^{((p-2)d+p)\tau} \left| d y_d(\tau) \right|^{2-p} \right), \tag{3.2}$$

$$Y''_d(\tau) = Y_d(\tau) \left( (p-1)^2(\eta - d)(p' + d) + \varepsilon(d-\alpha) e^{((p-2)d+p)\tau} \left| Y_d(\tau) \right|^{(2-p)/(p-1)} \right), \tag{3.3}$$

$$(p-1)y''(\tau) = \gamma^{2-p} y(\tau) \left( -\gamma^{p-1}(N+\gamma) - \varepsilon(\gamma+\alpha) \left| y(\tau) \right|^{2-p} \right) = - \left| Y(\tau) \right|^{(2-p)/(p-1)} Y'(\tau), \tag{3.4}$$

$$Y''(\tau) = Y(\tau) \left( -\gamma(N+\gamma) - \varepsilon(\gamma+\alpha) \left| Y(\tau) \right|^{(2-p)/(p-1)} \right) = \varepsilon\alpha y'(\tau), \tag{3.5}$$

$$(p-1)\zeta''(\tau) = -\varepsilon(p-2)((\alpha-\zeta) \left|\zeta\right|^{2-p} \left|y\right|^{-p} yy')(\tau) = \varepsilon(p-2)((\alpha-\zeta)(\gamma+\zeta) \left|\zeta y\right|^{2-p})(\tau), \tag{3.6}$$

$$(p-1)\sigma''(\tau) = -(p-2)((\sigma-\varepsilon) \left|\sigma\right|^{(2-p)/(p-1)} Y \left|y\right|^{(4-3p)/(p-1)} y')(\tau) = \zeta'(\tau)(\sigma(\tau)-\varepsilon). \tag{3.7}$$

**Proposition 3.6** *Let $w \not\equiv 0$ be any solution of ($\boldsymbol{E}_w$) on an interval $I$.*

*(i) If $\varepsilon = 1$ and $\alpha \leqq N$, then $w$ has at most one simple zero; if $\alpha < N$ and $w$ is regular, it has no zero. If $\alpha = N$ it has no simple zero and a compact support. If $\alpha > N$ and $w$ is regular, it has at least one simple zero.*

*(ii) If $\varepsilon = -1$ and $\alpha \geqq \min(0, \eta)$, then $w$ has at most one simple zero. If $w \not\equiv 0$ has a double zero, then it has no simple zero. If $\alpha > 0$ and $w$ is regular, it has no zero.*

*(iii) If $\varepsilon = -1$ and $-p' \leqq \alpha < \min(0, \eta)$, then $w'$ has at most one simple zero, consequently $w$ has at most two simple zeros, and at most one if $w$ is regular. If $\alpha < -p'$, the regular solutions have at least two zeros.*

**Proof.** (i) Let $\varepsilon = 1$. Consider two consecutive simple zeros $\rho_0 < \rho_1$ of $w$, with $w > 0$ on $(\rho_0, \rho_1)$ ; hence $w'(\rho_1) < 0 < w'(\rho_0)$. If $\alpha \leqq N$, we find from (2.1),

$$J_N(\rho_1) - J_N(\rho_0) = -\rho_1^{N-1} \left|w'(\rho_1)\right|^{p-2} - \rho_0^{N-1} w'(\rho_0)^{p-1} = (N - \alpha) \int_{\rho_0}^{\rho_1} s^{N-1} w ds,$$

which is contradictory; thus $w$ has at most one simple zero. The contradiction holds as soon as $\rho_0$ is simple, even if $\rho_1$ is not. If $w$ is regular with $w(0) > 0$, and $\rho_1$ is a first zero, and $\alpha < N$,

$$J_N(\rho_1) = -\rho_1^{N-1} \left|w'(\rho_1)\right|^{p-1} = (N - \alpha) \int_0^{\rho_1} s^{N-1} w ds > 0,$$

which is still impossible. If $\alpha = N$, the (Barenblatt) solutions are given by (1.10). Next suppose $\alpha > N$ and $w$ regular. If $w > 0$, then $J_N < 0$, thus $w^{-1/(p-1)} w' + r^{1/(p-1)} < 0$. Then the function $r \mapsto r^{p'} + \gamma w^{(p-2)/(p-1)}$ is non increasing and we reach a contradiction for large $r$. Thus $w$ has a first zero $\rho_1$, and $J_N(\rho_1) < 0$, thus $w'(\rho_1) \neq 0$.

(ii) Let $\varepsilon = -1$ and $\alpha \geqq \min(\eta, 0)$. Here we use the substitution (2.3) from some $d \neq 0$. If $y_d$ has a maximal point, where it is positive, and is not constant, then (3.2) holds. Taking $d \in (0, \min(\alpha, \eta))$ if $\eta > 0$, $d = \eta$ if $\eta \leqq 0$, we reach a contradiction. Hence $y_d$ has at most a simple zero, and no simple zero if it has a double one. Suppose $w$ regular and $\alpha > 0$. Then $w' > 0$ near 0, from Theorem 3.3. As long as $w$ stays positive, any extremal point $r$ is a strict minimum, from ($\mathbf{E}_w$), thus in fact $w'$ stays positive.

(iii) Let $\varepsilon = -1$ and $-p' \leqq \alpha < \min(0, \eta)$. Suppose that $w'$ and has two consecutive zeros $\rho_0 < \rho_1$, and one of them is simple, and use again (2.3) with $d = \alpha$. Then the function $Y_\alpha$ has an extremal point $\tau$, where it is positive and is not constant; from (3.3),

$$Y_\alpha''(\tau) = (p - 1)^2 (\eta - \alpha)(p' + \alpha) Y_\alpha(\tau), \tag{3.8}$$

thus $Y_\alpha''(\tau) \geqq 0$, which is contradictory. Next consider the regular solutions. They satisfy $Y_\alpha(\tau) = e^{(\alpha(p-1)+p)\tau}(|\alpha|\, a/N)(1 + o(1)$ near $-\infty$, from Theorem 3.3 and (2.3), thus $\lim_{\tau \to -\infty} Y_\alpha = 0$. As above $Y_\alpha$ cannot have any extremal point, then $Y_\alpha$ is positive and increasing. In turn $w' < 0$ from (2.3), hence $w$ has at most one zero. ■

**Proposition 3.7** *Let $w \not\equiv 0$ be any solution of ($\boldsymbol{E}_w$) on an interval $I$. If $\varepsilon = 1$, then $w$ has a finite number of isolated zeros. If $\varepsilon = -1$, it has a finite number of isolated zeros in any interval $[m, M] \cap I$ with $0 < m < M < \infty$.*

**Proof.** Let $Z$ be the set of isolated zeros on $I$. If $w$ has two consecutive isolated zeros $\rho_1 < \rho_2$, and $\tau \in (e^{\rho_1}, e^{\rho_2})$ is a maximal point of $|y_d|$, from (3.2), it follows that

$$\varepsilon e^{((p-2)d+p)\tau} \left|dy_d(\tau)\right|^{2-p}(d-\alpha) \leqq (p-1)d(d-\eta). \tag{3.9}$$

That means with $\rho = e^{\tau} \in (\rho_1, \rho_2)$,

$$\varepsilon \rho^p \left|w(\rho)\right|^{2-p}(d-\alpha) \leqq (p-1)d^{p-1}(d-\eta). \tag{3.10}$$

First suppose $\varepsilon = 1$ and fix $d > \alpha$. Consider the energy function

$$E(r) = \frac{1}{p'}\left|w'\right|^p + \frac{\alpha}{2}w^2.$$

It is nonincreasing since $E'(r) = -(N-1)r^{-1}\left|w'\right|^p - rw'^2$, thus bounded on $I \cap [\rho_1, \infty)$. Then $w$ is bounded, $\rho_2$ is bounded, $Z$ is a bounded set. If $Z$ is infinite, there exists a sequence of zeros $(r_n)$ converging to some point $\overline{r} \in [0, \infty)$, and a sequence $(\tau_n)$ of maximal points of $|y_d|$ converging to $\overline{\tau} = \ln \overline{r}$. If $\overline{r} > 0$, then $w(\overline{r}) = w'(\overline{r}) = 0$; we get a contradiction by taking $\rho = \rho_n = e^{\tau_n}$ in (3.10), because the left-hand side tends to $\infty$. If $\overline{r} = 0$, fixing now $d < \eta$, there exists a sequence $(\tau_n)$ of maximal points of $|y_d|$ converging to $-\infty$. Then $w(\rho_n) = O(\rho_n^{p/(p-2)})$, and $w'(\rho_n) = -d\rho_n^{-1}w(\rho_n) = O(\rho_n^{2/(p-2)})$, thus $E(\rho_n) = o(1)$. Since $E$ is monotone, it implies $\lim_{r\to 0} E(r) = 0$, hence $E \equiv 0$, and $w \equiv 0$, which is contradictory. Next suppose $\varepsilon = -1$ and fix $d < \alpha$. If $Z \cap [m, M]$ is infinite, we construct a sequence converging vers some $\overline{r} > 0$ and reach a contradiction as above. ■

**Proposition 3.8** *Let $y$ be any non constant solution of ($\boldsymbol{E}_y$), on a maximal interval $I$ where $(y, Y) \neq (0, 0)$, and $s$ be an extremity of $I$.*

*(i) If $y$ has a constant sign near $s$, then the same is true for $Y$.*

*(ii) If $y > 0$ is strictly monotone near $s$, then $Y, \zeta, \sigma$ are monotone near $s$.*

*(iii) If $y > 0$ is not strictly monotone near $s$, then $s = \pm\infty$, $\varepsilon(\gamma + \alpha) < 0$ and $y$ oscillates around $\ell$.*

*(iv) If $y$ is oscillating around 0 near $s$, then $\varepsilon = -1, s = \pm\infty, \alpha < -p'$; if $\alpha > -\gamma$, then $|y| > \ell$ at the extremal points.*

**Proof.** (i) The function $w$ has at most one extremal point on $I$ : at such a point, it satisfies $(\left|w'\right|^{p-2}w')' = -\varepsilon\alpha w$ with $\alpha \neq 0$. From (2.7), $Y$ has a constant sign near $s$.

(ii) Suppose $y$ strictly monotone near $s$. At any extremal point $\tau$ of $Y$, we find $Y''(\tau) = \varepsilon\alpha y'(\tau)$ from (3.5). Then $y'(\tau) \neq 0$, $Y''(\tau)$ has a constant sign. Thus $\tau$ is unique, and $Y$ is strictly monotone near $s$. Next consider $\zeta$. If there exists $\tau_0$ such that $\zeta(\tau_0) = \alpha$, then $\zeta'(\tau_0) = \alpha(\alpha - \eta)$, from system **(Q)**. If $\alpha \neq \eta$, then $\tau_0$ is unique, thus $\alpha - \zeta$ has a constant sign near $s$. Then $\zeta''(\tau)$ has a constant sign at any extremal point $\tau$ of $\zeta$, from (3.6), thus $\zeta$ is strictly monotone near $s$. If $\alpha = \eta$, then $\zeta \equiv \alpha$. At last consider $\sigma$. If there exists $\tau_0$ such that $\sigma(\tau_0) = \varepsilon$, then $\sigma'(\tau_0) = \varepsilon(\alpha - N)$ from System **(Q)**. If $\alpha \neq N$, then $\tau_0$ is unique, and $\sigma - \varepsilon$ has a constant sign near $s$. Thus $\sigma''(\tau)$ has a constant sign at any extremal point $\tau$ of $\sigma$, from (3.7) and assertion (i). If $\alpha = N$, then $\sigma \equiv \varepsilon$.

(iii) Let $y$ be positive and not strictly monotone near $s$. There exists a sequence $(\tau_n)$ strictly monotone, converging to $\pm\infty$, such that $y'(\tau_n) = 0$, $y''(\tau_{2n}) > 0 > y''(\tau_{2n+1})$. Since $y(\tau_n) = \gamma^{-1} |Y|^{(2-p)/(p-1)} Y(\tau_n)$, we deduce $Y < 0$ near $s$, from (i). From (3.5),

$$-\varepsilon(\gamma+\alpha)y(\tau_{2n+1})^{2-p} \leqq \gamma^{p-1}(N+\gamma) \leqq -\varepsilon(\gamma+\alpha))y(\tau_{2n})^{2-p}, \tag{3.11}$$

thus $\varepsilon(\gamma+\alpha) < 0$ and $y(\tau_{2n}) < \ell < y(\tau_{2n+1})$, and $Y(\tau_{2n+1}) < -(\gamma\ell)^{p-1} < Y(\tau_{2n})$. If $s$ is finite, then $y(s) = y'(s) = 0$, which is impossible; thus $s = \pm\infty$.

(iv) If $y$ is changing sign, then $\varepsilon = -1$ and $\alpha < -p'$, from Propositions 3.6 and 3.7. At any extremal point $\tau$,

$$(\alpha+\gamma)\, |y(\tau)|^{2-p} \leqq \gamma^{p-1}(N+\gamma)$$

from (3.4); if $\alpha > -\gamma$ it means $|y| > \ell$. ■

### 3.3 Double zeros and global existence

**Theorem 3.9** *For any $\overline{r} > 0$, there exists a unique solution $w$ of $(\boldsymbol{E}_w)$ defined in a interval $[\overline{r}, \overline{r} \pm h)$ such that*

$$w > 0 \quad \text{on } (\overline{r}, \overline{r} \pm h) \quad \text{and} \quad w(\overline{r}) = w'(\overline{r}) = 0.$$

*Moreover $\varepsilon h < 0$ and*

$$\lim_{r \to \overline{r}} |(\overline{r} - r)|^{(p-1)/(2-p)}\, \overline{r}^{1/(2-p)} w(r) = \pm((p-2)/(p-1))^{(p-1)/(p-2)}. \tag{3.12}$$

*In other words in the phase plane $(y, Y)$ there exists a unique trajectory $\mathcal{T}_\varepsilon$ converging to $(0,0)$ at $\varepsilon\infty$. It has the slope $\varepsilon$ and converges in finite time; it depends locally continuously of $\alpha$.*

**Proof.** Suppose that a solution $w \not\equiv 0$ exists on $[\overline{r}, \overline{r} \pm h)$ with $w(\overline{r}) = w'(\overline{r}) = 0$. From Propositions 3.7 and 3.8, up to a symmetry, $y > 0$, $|Y| > 0$ near $\overline{\tau} = \ln \overline{r}$, and $\lim_{\tau \to \ln \overline{r}} y = 0$, and $\sigma, \zeta$ are monotone near $\ln r$. Let $\mu$ and $\lambda$ be their limits. If $|\mu| = \infty$, then $|\lambda| = \infty$, because $\zeta = |Y|^{(2-p)/(p-1)} \sigma$, $|\zeta|^{p-2} \zeta = \sigma y^{2-p}$; then $f = 1/\zeta$ tends to 0; but

$$f' = -1 + \eta f + \varepsilon \frac{1 - \alpha f}{(p-1)\sigma}, \tag{3.13}$$

thus $f'$ tends to $-1$, which is impossible. Thus $\mu$ is finite. If $\lambda$ is finite, then $\mu = 0$, thus $\lambda = \alpha$, from system $(\mathbf{Q})$, $\ln w$ is integrable at $\overline{r}$, which is not true. Then $\lambda = \varepsilon\infty$, hence

$$\mu = \lim_{\tau \to \ln \overline{r}} \sigma = \varepsilon,$$

from system (**Q**). Then $\varepsilon Y > 0$ near $\bar{\tau}$, then $\varepsilon w' < 0$ near $\bar{r}$, thus $\varepsilon h < 0$. Consider system (**R**): as $\tau$ tends to $\bar{\tau}$, $\nu$ tends to $\pm\infty$, and $(g, s)$ converges to the stationary point $(0, -\varepsilon)$.

Reciprocally, setting $s = -\varepsilon/\beta + h$, the linearized system of system (**R**) at this point is given by

$$\frac{dg}{d\nu} = -\varepsilon\frac{p-2}{p-1}g, \qquad \frac{dh}{d\nu} = (\alpha - N)g + \varepsilon h.$$

The eigenvalues are $-\varepsilon(p-2)/(p-1)$ and $\varepsilon$, thus we find a saddle point. There are two trajectories converging to $(0, -\varepsilon)$. The first one satisfies $g \equiv 0$, it does not correspond to a solution of the initial problem. Then there exists a unique trajectory converging to $(0, -\varepsilon)$, as $\nu$ tends to $\varepsilon\infty$, with $g > 0$ near $\varepsilon\infty$. It is associated to the eigenvalue $-\varepsilon(p-2)/(p-1)$ and the eigenvector $((2p-3)/(p-1), \varepsilon(N-\alpha))$. It satisfies $dg/d\nu = -\varepsilon((p-2)/(p-1))g(1+o(1))$, thus $dg/d\tau = ((p-2)/(p-1))(1+o(1))$. Then $\tau$ has a finite limit $\bar{\tau}$, and $\tau$ increases to $\bar{\tau}$ if $\varepsilon = 1$ and decreases to $\bar{\tau}$ if $\varepsilon = -1$. In turn $|Y|^{(p-2)/(p-1)} = gs$ tends to 0, and $s$ tends to $\varepsilon$, thus $(y, Y)$ tends to $(0, 0)$ as $\tau$ tends to $\bar{\tau}$. Then $w$ and $w'$ converges to 0 at $\bar{r} = e^{\bar{\tau}}$. And $w'w^{-1/(p-1)} + (\varepsilon + o(1))r^{1/(p-1)} = 0$, which implies (3.12).

**Corollary 3.10** *Let $r_1 > 0$, and $a, b \in \mathbb{R}$ and $w$ be any local solution such that $w(r_1) = a$, $w'(r_1) = b$.*

*(i) If $(a, b) = (0, 0)$, then $w$ has a unique extension by 0 on $(r_1, \infty)$ if $\varepsilon = 1$, on $(0, r_1)$ if $\varepsilon = -1$.*

*(ii) If $(a, b) \neq (0, 0)$, $w$ has a unique extension to $(0, \infty)$.*

■

**Proof.** (i) Assume $a = b = 0$, the function $w \equiv 0$ is a solution. Let $w$ be any local solution near $r_1$, defined in an interval $(r_1 - h_1, r_1 + h_1)$ with $w(r_1) = w'(r_1) = 0$. Suppose that there exists $h_2 \in (0, h_1)$ such that $w(r_1 + \varepsilon h_1) \neq 0$. Let $\bar{h} = \inf\{h \in (0, h_1) : w(r_1 + \varepsilon h) \neq 0\}$, and $\bar{r} = r_1 + \varepsilon\bar{h}$, thus $w(\bar{r}) = w'(\bar{r}) = 0$, and for example $w > 0$ on some interval $(\bar{r}, \bar{r} + \varepsilon k))$ with $k > 0$. This contradicts theorem 3.9. Thus $w \equiv 0$ on $(r_1, r_1 + \varepsilon h_1)$.

(ii) From Theorems 3.9 and 3.3, $w$ has no double zero for $\varepsilon(r - r_1) < 0$, and has a unique extension to a maximal interval with no double zero. From (i) it has a unique extension to $(0, \infty)$. In particular any local regular solution is defined on $[0, \infty)$. ■

# 4 Asymptotic behaviour

Next the function $y$ is supposed to be monotone, thus $w$ has a constant sign near 0 or $\infty$, we can assume that $w > 0$.

**Proposition 4.1** *Let $y$ be any solution of ($\boldsymbol{E}_y$) strictly monotone and positive near $s = \pm\infty$.*

*(1) Then $(\zeta,\sigma)$ has a limit $(\lambda,\mu)$ near $s$, given by is some of the values*

$$
\begin{aligned}
&A_\gamma=\left(-\gamma,\varepsilon\frac{\alpha+\gamma}{N+\gamma}\right),\quad A_r=(0,\varepsilon\alpha/N)\,,\quad A_\alpha=(\alpha,0)\,,\\
&L_\eta=\eta\,(1,\infty)\,(\text{if } p\neq N),\quad L_+=(0,\infty)\,(\text{if } p\geqq N),\quad L_-=(0,-\infty)\,(\text{if } p>N).
\end{aligned}\tag{4.1}
$$

*(2) More precisely,*

*(i) Either $\varepsilon(\gamma+\alpha)<0$ and $(y,Y)$ converges to $\pm M_\ell$. Then $(\lambda,\mu)=A_\gamma$ and $(\varepsilon=1,s=\infty)$ or $(\varepsilon=-1,s{=}{-}\infty$ for $\alpha\leqq\alpha^*$, $s=\infty$ for $\alpha>\alpha^*)$.*

*(ii) Or $(y,Y)$ converges to $(0,0)$. Then ($s=\infty$ and $-\gamma<\alpha$) or ($s=-\infty$ and $\alpha<-\gamma$), or ($s=\varepsilon\infty$ and $\alpha=-\gamma$) and $(\lambda,\mu)=A_\alpha$.*

*(iii) Or $\lim_{\tau\to s}y=\infty$. Then $s=-\infty$. If $p<N$, then $(\lambda,\mu)=A_r$ or $L_\eta$. If $p=N$, then $(\lambda,\mu)=A_r$ or $L_+$. If $p>N$, then $(\lambda,\mu)=A_r,L_\eta,L_+$ or $L_-$.*

**Proof.** (1) The functions $Y,\sigma,\zeta$ are also monotone, and by definition $\zeta\sigma>0$. Thus $\zeta$ has a limit $\lambda\in[-\infty,\infty]$ and $\sigma$ has a limit $\mu\in[-\infty,\infty]$, and $\lambda\mu\geqq0$.

(i) $\lambda$ is finite. Indeed if $\lambda=\pm\infty$, then $f=1/\zeta$ tends to 0. From (3.13), either $\mu=\pm\infty$, then $f'$ tends to $-1$, which is imposible; or $\mu$ is finite, thus $\mu=\varepsilon$ from system (**Q**), then $f'$ tends to $(2-p)/(p-1)$, which is still contradictory.

(ii) Either $\mu$ is finite, thus $(\lambda,\mu)$ is a stationary point of system **(Q)**, equal to $A_\gamma,A_r$ or $A_\alpha$.

(iii) 0r $\mu=\pm\infty$ and $(\lambda,0)$ is a stationary point of system (**P**).

• If $p\neq N$, either $\lambda=\eta\neq0$ and $(\lambda,\mu)=L_\eta$; or $\lambda=0$ and $(\lambda,\mu)=L_+$ or $L_-$. In the last case $(\zeta,\psi)$ converges to $(0,0)$, and $\zeta'/\psi'=-(\eta\zeta/N\psi)(1+o(1))$, thus $\eta<0$, that means $p>N$.

• If $p=N$, then again $(\zeta,\psi)$ converges to $(0,0)$, thus $\mu=\pm\infty$, and $\psi'=N\psi(1+o(1))$, and necessarily $s=-\infty$. We make the substitution (2.4) with $d=0$. Then $y_0(\tau)=w(r)$, and $y_0$ satisfies

$$
y_0'=-\left|Y_0\right|^{(2-p)/(p-1)}Y_0=-\zeta y_0=o(y_0),\qquad Y_0'=\varepsilon e^{p\tau}y_0(\alpha-\zeta)=\varepsilon e^{p\tau}y_0\alpha(1+o(1).
$$

Thus for any $\upsilon>0$, we get $y_0=O(e^{-\upsilon\tau})$ and $1/y_0=O(e^{\upsilon\tau})$. Then $Y_0'$ is integrable, and $Y_0$ has a finite limit $|k|^{p-2}\,k$. Suppose that $k=0$. Then $Y_0=O(e^{(p-\upsilon)\tau})$, and $y_0$ has a finite limit $a\geqq0$. If $a\neq0$, then $Y_0'=\varepsilon\alpha ae^{p\tau}(1+o(1))$; in turn $Y_0=p^{-1}\varepsilon\alpha ae^{p\tau}(1+o(1))$, and $\psi=e^{p\tau}y_0/Y_0$ does not tend to 0. If $a=0$, then $y_0=O(e^{p'\tau})$, which contradicts the estimate of $1/y_0$. Thus $k>0$ and

$$
y_0=-k\tau(1+o(1),\qquad Y_0=k^{p-1}(1+o(1));\tag{4.2}
$$

hence $(\lambda,\mu)=L_+$.

(2) Since $y$ is monotone, we encounter one of the three following cases:

(i) $(y,Y)$ converges to $\pm M_\ell$. Then $(\lambda,\mu)=A_\gamma$ and $M_\ell$ is a source (or a weak source) for $\alpha \leqq \alpha^*$, a sink for $\alpha > \alpha^*$.

(ii) $y$ tends to 0. Since $\lambda$ is finite, $(y,Y)$ converges to $(0,0)$. And $|\sigma| = |\zeta|^{p-1} y^{p-2}$ tends to 0, thus $(\lambda,\mu)=A_\alpha$. If $-\gamma<\alpha$, seeing that $y' = -y(\gamma+\zeta)<0$ we find $s=\infty$. If $\alpha<-\gamma$, then $s=-\infty$. If $\alpha=-\gamma<0$, then $\varepsilon(\gamma+\zeta)>0$, from the first equation of **(Q),** thus $\varepsilon y'<0$, hence $s=\varepsilon\infty$.

(iii) $y$ tends to $\infty$. Either $\lambda \neq 0$, thus $|\sigma| = |\zeta|^{p-1} y^{p-2}$ tends to $\infty$, and $\lambda=\eta$ from system (**Q**), thus $p\neq N$, $(\lambda,\mu)=L_\eta$. Or $\lambda=0$ and $\mu$ is finite, thus $\mu=\varepsilon\alpha/N$, $(\lambda,\mu)=A_r$. Or $(\lambda,\mu)=L_0$; then either $p=N$, $L_0=L_\eta$, or $p>N$. In any case, $y'=-y(\gamma+\zeta)<0$, from (1.2), hence $s=-\infty$.
■

Next we apply these results to the functions $w$ :

**Proposition 4.2** *We keep the assumptions of Proposition 4.1. Let $w$ be the solution of $(\boldsymbol{E}_w)$ associated to $y$ by (2.7).*

*(i) If $(\lambda,\mu)=A_\gamma$ (near 0 or $\infty$), then*

$$\lim r^{-\gamma} w = \ell. \tag{4.3}$$

*(ii) If $(\lambda,\mu)=A_\alpha$ (near 0 or $\infty$), then*

$$\lim r^{\alpha} w = L > 0 \qquad \textit{if } \alpha \neq -\gamma, \tag{4.4}$$

$$\lim r^{-\gamma} (\ln r)^{1/(p-2)} w = ((p-2)\gamma^{p-1}(N+\gamma))^{-1/(p-2)} \qquad \textit{if } \alpha=-\gamma. \tag{4.5}$$

*(iii) If $p<N$ and $(\lambda,\mu)=L_\eta$, then*

$$\lim_{r\to 0} r^{\eta} w = c > 0. \tag{4.6}$$

*(iv) If $p>N$ and $(\lambda,\mu)=L_\eta$, then*

$$\lim_{r\to 0} r^{-|\eta|} w = c > 0. \tag{4.7}$$

*(v) If $p=N$ and $(\lambda,\mu)=L_+$, then*

$$\lim_{r\to 0} |\ln r|^{-1} w = k > 0, \qquad \lim_{r\to 0} r w' = -k \qquad \textit{if } p=N. \tag{4.8}$$

*(vi) If $p>N$ and $(\lambda,\mu)=L_+$, or $L_-$, then*

$$\lim_{r\to 0} w = a > 0, \qquad \lim_{r\to 0}(-r^{(N-1)/(p-1)} w') = c > 0, \tag{4.9}$$

*or*

$$\lim_{r\to 0} w = a > 0, \qquad \lim_{r\to 0}(-r^{(N-1)/(p-1)} w') = c < 0. \tag{4.10}$$

**Proof.** (i) This follows directly from (2.7).

(ii) From (2.16), $rw'(r) = -\alpha w(r)(1 + o(1)$. We are lead to three cases.

• Either $-\gamma < \alpha$, and $s = \infty$. For any $\upsilon > 0$, we find $w = O(r^{-\alpha+\upsilon})$ and $1/w = O(r^{\alpha+\upsilon})$ near $\infty$ and $w' = O(r^{-\alpha-1+\upsilon})$. Then $J'_\alpha(r) = O(r^{\alpha(2-p)-p-1+\upsilon})$, hence $J'_\alpha$ is integrable, $J_\alpha$ has a limit $L$. And $\lim r^\alpha w = L$, seeing that $J_\alpha(r) = r^\alpha w(1 + o(1))$. If $L = 0$, then $r^\alpha w = O(r^{\alpha(2-p)-p+\upsilon})$, which contradicts the estimate of $1/w = O(r^{\alpha+\upsilon})$ for $\upsilon$ small enough. Thus $L > 0$.

• Or $\alpha < -\gamma$ and $s = -\infty$. For any $\upsilon > 0$, we find $w = O(r^{-\alpha-\upsilon})$ and $1/w = O(r^{\alpha+\upsilon})$ near 0 and $w' = O(r^{-\alpha-1-\upsilon})$. Then $J'_\alpha(r) = O(r^{\alpha(2-p)-p-1-\upsilon})$, and $J'_\alpha$ is still integrable, $J_\alpha$ has a limit $L$, and $\lim r^\alpha w = L$. If $L = 0$, then $r^\alpha w = O(r^{\alpha(2-p)-p-\upsilon})$, which contradicts the estimate of $1/w$. Thus again $L > 0$.

• Or $\alpha = -\gamma$ and $s = \varepsilon\infty$. Then $Y = -\gamma^{p-1}y^{p-1}(1 + o(1))$, and $\mu = 0$, thus $y - \varepsilon Y = y(1 + o(1))$. From System (**S**),

$$(y - \varepsilon Y)' = \varepsilon(N + \gamma)Y = -\varepsilon(N + \gamma)\gamma^{p-1} (y - \varepsilon Y)^{p-1} (1 + o(1)).$$

Then $y = (N + \gamma)\gamma^{p-1}(p - 2) |\tau|)^{-1/(p-2)}(1 + o(1))$, which is equivalent to (4.5).

(iii) From (2.16), we get $rw'(r) = -\eta w(r)(1+o(1)$. We use (2.3) with $d = \eta$, thus $y_\eta = r^\eta w$. We find $y_\eta = O(e^{-\upsilon\tau})$, $1/y_\eta = O(e^{-\upsilon\tau})$, in turn $Y_\eta = O(e^{-\upsilon\tau})$. From (2.4), $Y'_\eta = O(e^{(p+(p-2)\eta-\upsilon)\tau})$, thus $Y'_\eta$ is integrable, hence $Y_\eta$ has a finite limit. Now $(e^{-\eta\tau}y_\eta)' = -e^{-\eta\tau}Y_\eta^{1/(p-1)}$, and $\eta > 0$, thus $y_\eta$ has a limit $c$. If $c = 0$, then $Y_\eta = O(e^{(p+(p-2)\eta-\upsilon)\tau})$, $y_\eta = O(e^{((p+(p-2)\eta)/(p-1)-\upsilon)\tau})$, which contradicts $1/y_\eta = O(e^{-\upsilon\tau})$ for $\upsilon$ small enough. Then (4.6) holds.

(iv) As above, $Y_\eta$ has a finite limit. In turn $r^{-|\eta|+1}w' = |Y_\eta|^{(2-p)/(p-1)} Y_\eta$ has a limit $c|\eta|$ and $w$ has a limit $a \geqq 0$. From (2.16), $rw' = |\eta| w(1 + o(1)$, hence $a = 0$. Then $c \geqq 0$; if $b = 0$, then $Y < 0$, the function $v = -e^{(\gamma+N)\tau}Y > 0$ tends to 0 and

$$v' = -e^{(\gamma+N)\tau}\varepsilon(\alpha - \eta)y(1 + o(1)) = -\varepsilon(\alpha - \eta) |\eta| e^{-(\gamma+N)(p-2)/(p-1)\tau}v^{1/(p-1)};$$

we reach again a contradiction.Thus $a = 0$ and $c > 0$, and (4.7) holds.

(v) Assertion (4.8) follows from (4.2).

(vi) Here $rw' = o(w)$, thus $w + |w'| = O(r^{-k})$ for any $k > 0$. Then $J'_N$ is integrable, $J_N$ has a limit at 0, and $\lim_{r\to 0} r^N w = 0$. Thus $\lim_{r\to 0} r^{(N-1)/(p-1)}w' = -c \in \mathbb{R}$, $\lim_{r\to 0} J_N = -\varepsilon |c|^{p-2} c$, $\lim_{r\to 0} w = a \geq 0$. If $c = 0$, then $J_N(r) = \int_0^r J'_N(s)ds$, implying that $\lim_{r\to 0} w' = 0$. Either $a > 0$ and then $w$ is regular, then $\lim_{\tau\to-\infty} \sigma = \varepsilon$; or $a = 0$, then $w' > 0$ and $(w')^{p-1} = O(rw)$; in both cases we get a contradiction. Thus $c \neq 0$. If $a = 0$, we find $\lim_{\tau\to-\infty} \zeta = \eta$, which is not true, hence $a > 0$. In any case (4.9) or (4.10) holds. ■

Now we study the cases where $y$ is not monotone, and eventually changing sign.

**Proposition 4.3** *Suppose $\varepsilon = -1$. Let $w \not\equiv 0$ be any solution of ($\boldsymbol{E}_w$).*

*(i) If $\alpha \leqq -\gamma$, then $w$ is oscillating near $0$ at $\infty$.*

*(ii) If $\alpha < 0$, then $y$ and $Y$ are bounded at $\infty$.*

**Proof.** (i) Suppose by contradiction that $w \geqq 0$ for large $r$, then $y \geqq 0$ for large $\tau$. If $y > 0$ near $\infty$, then from Proposition 3.8, either $y$ is constant, which is impossible since $(0,0)$ is the unique stationary point; or $y$ is strictly monotone, which contradicts Proposition 4.1. Then there exists a sequence $(\tau_n)$ tending to $\infty$ such that $y(\tau_n) = y'(\tau_n) = 0$; from Theorem 3.10, $y \equiv 0$ on $(-\infty, \tau_n)$, thus $y \equiv 0$.

(ii) Consider the function

$$\tau \mapsto R(\tau) = \frac{y^2}{2} + \frac{|Y|^{p'}}{p' \, |\alpha|};$$

it satisfies

$$R'(\tau) = -\gamma y^2 + \frac{1}{|\alpha|} |Y|^{2/(p-1)} - \frac{N+\gamma}{|\alpha|} |Y|^{p'}.$$

From the Young inequality,

$$|\alpha| \left(R'(\tau) + \gamma R(\tau)\right) = |Y|^{2/(p-1)} - (N + \frac{1}{p-2}) |Y|^{p'} \leqq (\frac{2}{Np+\gamma})^{(p-2)/2} \leqq 1$$

thus $R(\tau)$ is bounded for large $\tau$, at least by $1/|\alpha|\gamma$. ∎

**Proof.**

**Proposition 4.4** *(i) Assume $\varepsilon = 1$, or $\varepsilon = -1$, $\alpha \notin (\alpha_2, \alpha_1)$. Then for any trajectory of system ($\boldsymbol{S}$) in $\mathcal{Q}_4$ near $\pm\infty$, $y$ is strictly monotone near $\pm\infty$.*

*(ii) Assume $\varepsilon = 1$, and $\alpha \leqq \alpha^*$ or $-p' \leqq \alpha$. Then system ($\boldsymbol{S}$) admits no cycle in $\mathcal{Q}_4$ (or $\mathcal{Q}_2$).*

∎

**Proof.** (i) In any case $M_\ell$ is a node point. Following [4, Theorem 2.24], we use the linearization defined by (2.9). Consider the line $L$ given by the equation $A\overline{y} + \overline{Y} = 0$, where $A$ is a real parameter. The points of $L$ are in $\mathcal{Q}_4$ whenever $\overline{Y} < (\gamma\ell)^{p-1}$ and $-\ell < \overline{y}$. We get

$$A\overline{y}' + \overline{Y}' = \left(\varepsilon\nu(\alpha)A^2 + (N + \nu(\alpha))A + \varepsilon\alpha\right)\overline{y} + (A + \varepsilon)\Psi(\overline{Y}).$$

From (2.13), apart from the case $\varepsilon = 1, \alpha = N$, we can find an $A$ such that

$$\varepsilon\nu(\alpha)A^2 + (N + \nu(\alpha))A + \varepsilon\alpha = 0,$$

and $A + \varepsilon \neq 0$. Moreover $\Psi(\overline{Y}) \leqq 0$ on $L \cap \mathcal{Q}_4$. Indeed $(p-1)\Psi'(t) = -((\gamma\ell)^{p-1} - t)^{(2-p)/(p-1)} + (\gamma\ell)^{2-p}$, thus $\Psi$ has a maximum $0$ on $\left(-\infty, (\delta\ell)^{p-1}\right)$ at point $0$. Then the orientation of the vector field does not change along $L \cap \mathcal{Q}_4$. In particular $y$ cannot oscillate around $\ell$, thus $y$ is monotone, from Proposition 3.8. If $\varepsilon = 1, \alpha = N$, then $Y \equiv y \in (\ell, \infty)$ defines the trajectory $\mathcal{T}_r$, corresponding to the solutions given by (1.10) with $K > 0$. No solution $y$ can oscillate around $\ell$, since the trajectory cannot meet $\mathcal{T}_r$.

(ii) Suppose that there exists a cycle in $\mathcal{Q}_4$.

• Assume $\alpha \leqq \alpha^*$. Here $M_\ell$ is a source, or a weak source, from Proposition 2.5. Any trajectory starting from $M_\ell$ at $-\infty$ has a limit cycle in $\mathcal{Q}_1$, which is attracting at $\infty$. Writing System (**S**) under the form $y' = f_1(y,Y), Y' = f_2(y,Y)$, the mean value of the Floquet integral on the period $[0, \mathcal{P}]$ is given by

$$I = \oint (\frac{\partial f_1}{\partial y}(y,Y) + \frac{\partial f_2}{\partial Y}(y,Y))d\tau = \oint (\frac{|Y|^{(2-p)/(p-1)}}{p-1} - 2\gamma - N)d\tau. \tag{4.11}$$

Such a cycle is not unstable, thus $I \leqq 0$. Now

$$\oint (\alpha y' - \gamma Y')d\tau = 0 = (\alpha+\gamma)\oint |Y|^{1/(p-1)}\, d\tau - \gamma(\gamma+N)\oint |Y|\, d\tau.$$

From the Jensen and Hölder inequalities, since $1/(p-1) < 1$,

$$\gamma(\gamma+N)(\oint |Y|^{1/(p-1)}\, d\tau)^{p-2} \leqq \alpha + \gamma,$$

$$1 \leqq \left(\oint |Y|^{(2-p)/(p-1)})d\tau\right)\left(\oint |Y|^{1/(p-1)}\, d\tau\right)^{p-2} \leqq \frac{(p-1)(2\gamma+N)}{\gamma(\gamma+N)}(\alpha+\gamma),$$

then $\alpha^* < \alpha$, which is contradictory.

• Assume $-p' \leqq \alpha < 0$. Consider the functions $y_\alpha = e^{(\alpha+\gamma)\tau}y$ and $Y_\alpha = e^{(\alpha+\gamma)(p-1)\tau}Y$ defined by (2.3) with $d = \alpha$. They vary respectively from 0 to $\infty$ and from 0 to $-\infty$. They have no extremal point. Indeed at such a point, from (3.2) and (3.3) $y''_\alpha$ or $Y''_\alpha$ have a strict constant sign for $\alpha \neq \eta, p'$, which is contradictory. If $\alpha = \eta$ or $p'$, from uniqueness $y_\alpha$ or $Y_\alpha$ is constant, thus $y$ or $Y$ is monotone, which is impossible. In any case $y'_\alpha > 0 > Y'_\alpha$ on $(-\infty, \infty)$. Next, from (2.5) and (2.6),

$$\frac{y''_\alpha}{y'_\alpha} + \eta - 2\alpha - \frac{1}{p-1}Y^{(2-p)/(p-1)} = \alpha(\eta-\alpha)\frac{y_\alpha}{y'_\alpha}, \tag{4.12}$$

$$\frac{Y''_\alpha}{Y'_\alpha} + (p-1)(\eta - 2\alpha - p') - \frac{1}{p-1}Y^{(2-p)/(p-1)} = (p-1)^2(\eta-\alpha)(p'+\alpha)\frac{Y_\alpha}{Y'_\alpha}. \tag{4.13}$$

Let us integrate on the period $\mathcal{P}$. If $\eta \leqq \alpha < 0$, then $\eta - N - 2(\alpha+\gamma) \geqq 0$ from (4.12), which is contradictory. If $-p' \leqq \alpha < \eta$, then $-2(\alpha + p' + \gamma) > 0$ from (4.13), still contradictory. ∎

## 5 New local existence results

At Proposition 4.1 we gave all the *possible* behaviours of the positive solutions near $\pm\infty$. Next we prove their existence, and uniqueness or multiplicity. The case $p > N$ is very delicate.

**Theorem 5.1** *(i) Suppose $p < N$. In the phase plane $(y, Y)$ of system ($\boldsymbol{S}$) there exist an infinity of trajectories $\mathcal{T}_\eta$ such that $\lim_{\tau\to-\infty}(\zeta, \sigma) = L_\eta$; the corresponding $w$ satisfy (4.6).*

*(ii) Suppose $p > N$. There exist a unique trajectory $\mathcal{T}_u$ such that $\lim_{\tau\to-\infty}(\zeta, \sigma) = L_\eta$; in other words for any $c \neq 0$, there exists a unique solution $w$ of equation ($\boldsymbol{E}_w$) such that (4.7) holds.*

**Proof.** Suppose that such a trajectory exists in the plane $(y, Y)$. In the phase plane $(\zeta, \psi)$ of System (**P**), $\zeta$ and $\psi$ keep a strict constant sign, because the two axes $\zeta = 0$ and $\psi = 0$ contain particular trajectories, and $(\zeta, \psi)$ converges to $(\eta, 0)$ at $-\infty$. Reciprocally, setting $\zeta = \eta + \bar{\zeta}$, the linearized problem at point $(\eta, 0)$

$$\bar{\zeta}' = \eta\bar{\zeta} + \eta(\alpha - \eta)\varepsilon\psi/(p-1), \qquad \psi' = (N - \eta)\psi,$$

admits the eigenvalues $\eta$ and $N - \eta$. The trajectories linked to the eigenvalue $\eta$ are tangent to the line $\psi = 0$.

(i) Case $p < N$. Then $\eta > 0$, and $(\eta, 0)$ is a source. In the plane $(\zeta, \psi)$ there exist an infinity of trajectories, starting from this point at $-\infty$, such that $\psi > 0$, and $\lim_{\tau\to-\infty}\zeta = \eta$, thus $\zeta > 0$. In the phase plane $(y, Y)$, setting $y = (\psi\,|\zeta|^{p-2}\,\zeta)^{2-p}$ and $Y = y/\psi$, they correspond to an infinity of trajectories in the plane $(y, Y)$ such that $\lim_{\tau\to-\infty}(\zeta, \sigma) = L_\eta$, and (4.6) holds from Proposition (4.2).

(ii) Case $p > N$. Then $\eta < 0$, and $(\eta, 0)$ is a saddle point. In the plane $(\zeta, \psi)$, there exists a unique trajectory starting from $(\eta, 0)$, tangentially to the vector $(\eta(\alpha - \eta)\varepsilon/(p-1), N - \eta)$, with $\psi < 0$; it defines a unique trajectory $\mathcal{T}_u$ in the plane $(y, Y)$, and (4.7) holds. From Remark 2.1, we get a solution for any $c \neq 0$. ■

**Theorem 5.2** *(i) Suppose $p = N$. In the phase plane $(y, Y)$, there exists an infinity of trajectories $\mathcal{T}_+$ such that $\lim_{\tau\to-\infty}(\zeta, \sigma) = L_+$; then $w$ satisfies (4.8).*

*(ii) Suppose $p > N$. Then there exist an infinity of trajectories $\mathcal{T}_+$ (resp. $\mathcal{T}_-$) such $\lim_{\tau\to-\infty}(\zeta, \sigma) = L_+$ (resp. $L_-$); then the corresponding solutions $w$ of ($\boldsymbol{E}_w$) satisfy (4.9) (resp. (4.10).*

*More precisely for any $k > 0$ (for $p = N$) or any $a > 0$ and $c \neq 0$ (for $p > N$) there exists a unique function $w$ satisfying those conditions.*

**Proof.** If $\lim_{\tau\to-\infty}(\zeta, \sigma) = L_\pm$, then $\lim_{\tau\to-\infty}(\zeta, \psi) = (0, 0)$, with $\zeta\psi > 0$ in case of $L_+$, $\zeta\psi < 0$ in case of $L_-$. The linearization of System (**P**) near $(0, 0)$ is given by

$$\zeta' = |\eta|\,\zeta, \qquad \psi' = N\psi.$$

(i) Case $p=N$. The phase plane study is delicate because 0 is a center, thus we use a fixed method. Suppose that such a trajectory exists, and consider the substitution (2.3) with $d=0$. From (4.2), there exists $k>0$ such that $\zeta=|Y_0|^{(2-p)/(p-1)}/y_0=-\tau^{-1}(1+o(1))>0$, and $\psi=-k^{2-p}\tau e^{N\tau}(1+o(1))>0$. Then $\zeta'=\tau^{-2}(1+o(1))$ from System (**P**). The function

$$V=\psi e^{-N/\zeta}\zeta$$

satisfies $\lim_{\tau\to-\infty}V=k^{2-p}$, and

$$V'=\frac{Ve^{N/\zeta}}{(N-1)\zeta^2}(\varepsilon\left(\alpha-\zeta\right)\left(N-(N-2)\zeta\right)V+2N(N-1)\zeta^2e^{-N/\zeta}).$$

Thus $\varepsilon\alpha(V-k^{2-p})>0$ near $-\infty$. Moreover $\lim_{\tau\to-\infty}\zeta'/V'=0$, so that $\zeta$ can be considered as a function of $V$ near $k^{2-p}$, with $\lim_{V\to k^{2-p}}\zeta=0$ and

$$\frac{d\zeta}{dV}=K(V,\zeta),\qquad K(V,\zeta):=\frac{\zeta^2}{V}\frac{\varepsilon\left(\alpha-\zeta\right)V+(N-1)\zeta^2e^{-N/\zeta}}{\varepsilon\left(\alpha-\zeta\right)\left(N-(N-2)\zeta\right)V+2N(N-1)\zeta^2e^{-N/\zeta}}.$$

Reciprocally, extending the function $\zeta^2e^{-N/\zeta}$ by 0 for $\zeta\leqq 0$, the function $K$ is of class $C^1$ near $(k^{2-p},0)$. For any $k>0$, there exists a unique local solution $V\mapsto\zeta(V)$ on a interval $\mathcal{V}$ where $\varepsilon\alpha(V-k^{2-p})>0$, such that $\zeta(k^{2-p})=0$. And $d\zeta/dV=(\zeta^2/Nk^{2-p})(1+o(1))$ near 0, thus $\zeta>0$. In the plane $(\zeta,\psi)$, taking one point $P$ on the curve $\mathcal{C}=\left\{(\zeta(V),V\zeta(V)e^{N/\zeta(V)}):v\in\mathcal{V}\right\}$, there exists a unique solution of System (**P**) issued from $P$ at time 0. Its trajectory is on $\mathcal{C}$, thus it converges to $(0,0)$, with $\zeta,\psi>0$. It corresponds to a unique trajectory $\mathcal{T}_+$ in the plane $(y,Y)$, and $(\zeta,\sigma)$ converges to $L_+$, as $\tau$ tends to $-\infty$, from Proposition 4.1. The corresponding functions $w$ satisfy (4.8) from Proposition (4.2).

(ii) Case $p>N$. Here $(0,0)$ is a source for System (**P**). The lines $\zeta=0$ and $\psi=0$ contain trajectories. There exists an infinity of trajectories converging to $(0,0)$, with $\zeta\psi\neq 0$; moreover, if $N\geqq 2$, then $|\eta|<N$, thus $\lim_{\tau\to-\infty}(\psi/\zeta)=0$. Our claim is more precise. Given $a>0$ and $c\neq 0$, we look for a solution $w$ of ($\mathbf{E}_w$) such that $\lim_{r\to 0}w=a$, $\lim_{r\to 0}r^{\eta+1}w'=-c$. By scaling we can assume $a=1$. If $w_1$ is a such a solution, then $\zeta$ and $\psi$ have the sign of $c$ near 0, and $\zeta\left(\tau\right)=ce^{|\eta|\tau}(1+o(1))$ and $|c|^{p-2}c\psi\left(\tau\right)=e^{N\tau}(1+o(1))$. The function

$$v=c(|c|^{p-2}c\psi)^{1/\kappa}/\zeta,\qquad\text{with }\kappa=N/\left|\eta\right|>1,$$

satisfies $\lim_{\tau\to-\infty}v=1$, and can be expressed locally as a function of $\zeta$, and

$$\frac{dv}{d\zeta}=H(\zeta,v),\qquad H(\zeta,v):=-\frac{v}{\kappa}\frac{(p-1)(\kappa+1)+\varepsilon(\kappa-p+1)\left|c\right|^{1-p-\kappa}(\zeta-\alpha)\left|\zeta\right|^{\kappa-1}v^\kappa}{(p-1)(\zeta-\eta)+\varepsilon\left|c\right|^{1-p-\kappa}(\alpha-\zeta)\left|\zeta\right|^{\kappa-1}\zeta v^\kappa}.$$

Reciprocally, there exists a unique solution $\zeta\mapsto v(\zeta)$ of this equation on a small interval $[0,hc)$, with $h>0$, such that $v(0)=1$. Indeed $H$ is locally continuous in $\xi$ and $C^1$ in $v$. Taking one point $P$ on the curve $\mathcal{C}' = \left\{ (\zeta, |c|^{1-p-\kappa} |\zeta|^{\kappa-1} \zeta v(\zeta)) : \zeta \in [0, hc) \right\}$, there exists a unique solution of System (**P**) issued from $P$ at time 0. Its trajectory is on $\mathcal{C}'$, thus converges to $(0,0)$ with $\zeta\psi > 0$. It corresponds to a solution $(y, Y)$ of System (**S**), such that $(\zeta, \sigma)$ converges to $L_+$, as $\tau$ tends to $-\infty$, from Proposition 4.1. The corresponding function, called $w_2$, satisfies $\lim_{r\to 0} r^{\eta+1} w_2^{\gamma^{-1}|\eta|-1} w_2' = -c$; thus $w_2$ has a limit $a_2$, and $\lim_{r\to 0} r^{\eta-1} w_2' = a_2^{1-s} b$. Moreover $a_2 \neq 0$, because $a_2 = 0$ implies that $r^{-\gamma} w_2$ has a nonzero limit, thus $(\zeta, \sigma)$ converges to $A_\gamma$. The function $w(r) = a_2^{-1} w_2(a_2^{1/\gamma} r)$ satisfies $\lim_{r\to 0} w = 1$, and $\lim_{r\to 0} r^{\eta-1} w' = -c$, and the proof is done.

**Theorem 5.3** *(i) In the phase plane $(y, Y)$, for any $\alpha \neq 0$ there exists at least a trajectory $\mathcal{T}_\alpha$ converging to $(0,0)$ with $y > 0$, and $\lim(\zeta, \sigma) = A_\alpha$. The convergence holds at $\infty$ if $-\gamma < \alpha$, or $-\infty$ if $\alpha < -\gamma$, or $\varepsilon\infty$ if $\alpha = -\gamma$.*

*(ii) If $\varepsilon(\gamma + \alpha) < 0$, $\mathcal{T}_\alpha$ is unique, it is the unique trajectory converging to $(0,0)$ at $-\varepsilon\infty$ with $y > 0$, and it depends locally continuously of $\alpha$.*

■

**Proof.** (i) Suppose that such a trajectory exists. Then $\tau$ tends to $\infty$ if $-\gamma < \alpha$, or $-\infty$ if $\alpha < -\gamma$, or $\varepsilon\infty$ if $\alpha = -\gamma$, from Proposition 4.1. Consider System (**R**), where $g$, $s$ and $\nu$ are defined by (2.18). Then $(g, s)$ converges to $(-1/\alpha, 0)$, with $gs > 0$, and $\nu$ tends to the same limits as $\tau$, since $Y$ converges to 0. Reciprocally, in the plane $(g, s)$, let us show the existence of a trajectory converging to $(-1/\alpha, 0)$, different from the line $s = 0$. Setting $g = -1/\alpha + \bar{g}$, the linearized system at this point is

$$\frac{d\bar{g}}{d\nu} = -\frac{\varepsilon}{p-1}\bar{g} + \frac{\eta - \alpha}{\alpha^2} s, \qquad \frac{ds}{d\nu} = 0,$$

thus we find a center: the eigenvalues are 0 and $\lambda = \varepsilon/(p-1)$. Since the system is polynomial, it is known that System (**R**) admits a trajectory, depending locally continuously of $\alpha$, such that $sg > 0$, and tangent to the eigenvector $((p-1)(\eta - \alpha), \varepsilon\alpha^2)$. It satisfies $ds/d\nu = (p-2)(\alpha + \gamma)s^2(1 + o(1))$. Then $ds/d\tau = -(p-2)\alpha(\alpha + \gamma)s(1 + o(1))$, thus $\tau$ tends to $\pm\infty$. And $|y|^{p-2} = |s| \, |g|^{1/(p-1)}$, then $y$ tends to 0, $(y, Y)$ converges to $(0,0)$, and $\lim(\zeta, \sigma) = A_\alpha$.

(ii) Suppose $\varepsilon(\gamma + \alpha) < 0$. Consider two trajectories $\mathcal{T}_1, \mathcal{T}_2$ in the plane $(y, Y)$, converging to $(0,0)$ at $-\varepsilon\infty$, with $y > 0$. They are different from $\mathcal{T}_\varepsilon$ which converges at $\varepsilon\infty$, thus $\lim(\zeta_i, \sigma_i) = (\alpha, 0)$ from Proposition 4.1. Then $\zeta_1, \zeta_2$ can locally be expressed as a function of $y$, and

$$y\frac{d(\zeta_1 - \zeta_2)^2}{dy} = 2(F(\zeta_1, y) - F(\zeta_2, y))\,(\zeta_1 - \zeta_2)$$

near 0, where

$$F(\zeta, y) = \frac{1}{\gamma + \zeta}(-\zeta(\zeta - \eta) + \frac{\varepsilon}{p-1} |\zeta y|^{2-p} (\zeta - \alpha)).$$

Then $(\zeta_1 - \zeta_2)^2$ is nonincreasing, seeing that $\partial F/\partial\zeta(\zeta, y) = -((p-1)\varepsilon(\gamma + \alpha))^{-1} |\alpha y|^{2-p} (1 + o(1))$. Hence $\zeta_1 \equiv \zeta_2$ near 0, and $\mathcal{T}_1 \equiv \mathcal{T}_2$. ■

## 6 The case $\varepsilon = 1$, $-\gamma \leqq \alpha$

In that Section and in Sections 7, 8 and 9 we describe the solutions of ($\mathbf{E}_w$). When we give a *uniqueness* result, we mean that $w$ is unique, *up to a scaling,* from Remark 2.1.

**Theorem 6.1** *Assume* $\varepsilon = 1$, $-\gamma \leqq \alpha$ $(\alpha \neq 0)$.

*Any solution* $w$ *of* ($\boldsymbol{E}_w$) *has a finite number of simple zeros, and satisfies (4.4) or (4.5) near* $\infty$ *or has a compact support. Either* $w$ *is regular, or* $|w|$ *satisfies (4.6),(4.8), (4.7),(4.9) or (4.10) near* $0$*, and there exist solutions of each type.*

*(1) Case* $\alpha < N$*. All regular solutions have a strict constant sign, and satisfy (4.4) or (4.5) near* $\infty$*. Moreover there exist (and exhaustively, up to a symmetry)*

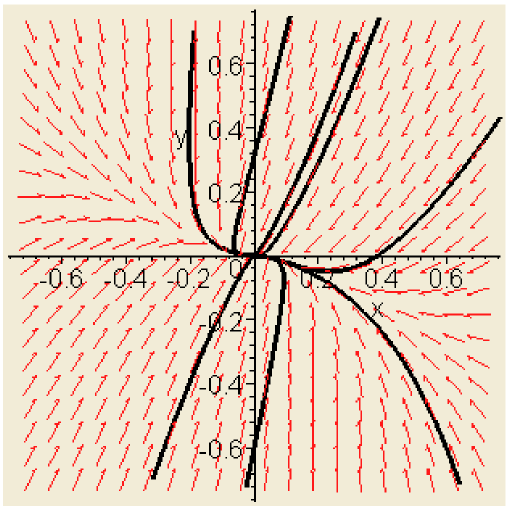
*(i) a unique nonnegative solution with (4.6)or (4.8) or (4.9)) near* $0$*, and compact support;*
*(ii) positive solutions with the same behaviour at* $0$ *and (4.4) or (4.5) near* $\infty$;
*(iii) solutions with one simple zero, and* $|w|$ *has the same behaviour at* $0$ *and* $\infty$;
*(iv) for* $p > N$, *a unique positive solution with (4.7) near* $0$*, and (4.4) or (4.5) near* $\infty$*;*
*(v) for* $p > N$, *positive solutions with (4.10) near* $0$*, and (4.4) or (4.5) near* $\infty$.

*(2) Case* $\alpha = N$*. Then the regular (Barenblatt) solutions have a constant sign with compact support. If* $p \leqq N$*, all the other solutions are of type (iii). If* $p > N$*, there exist also solutions of type (iv) and (v).*

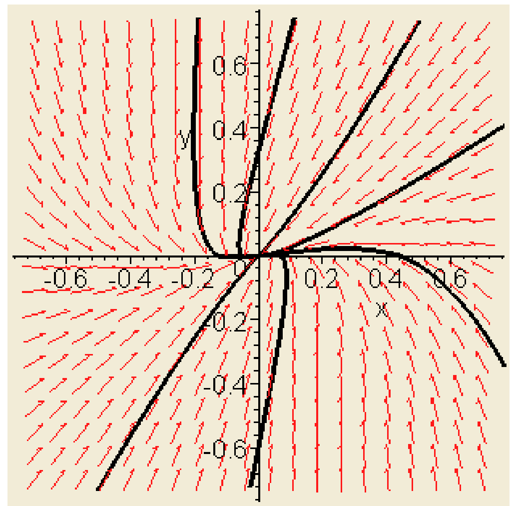
*(3) Case* $\alpha > N$.

*Either the regular solutions have* $m$ *simple zeros and satisfy satisfies (4.4) near* $\infty$*. Then there exist*
*(vi) a unique solution with* $m$ *simple zeros,* $|w|$ *satisfies (4.6), (4.8) or(4.9) near* $0$*, with compact support;*
*(vii) solutions with* $m+1$ *simple zeros,* $|w|$ *satisfies (4.6), (4.8) or (4.9) near* $0$*, and (4.4) or (4.5) near* $\infty$;
*(viii) for* $p > N$*, solutions with* $m$ *simple zeros,* $|w|$ *satisfies (4.9),(4.7) or (4.10) near* $0$*, and (4.4) or (4.5) near* $\infty$.

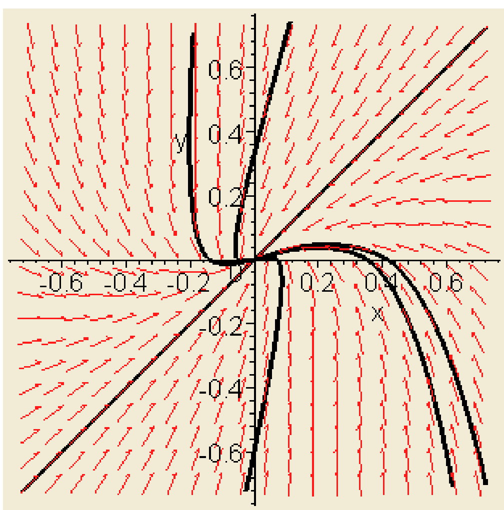
*Or the regular solutions have* $m$ *simple zeros and a compact support. Then the other solutions are of type (vii) or (viii).*



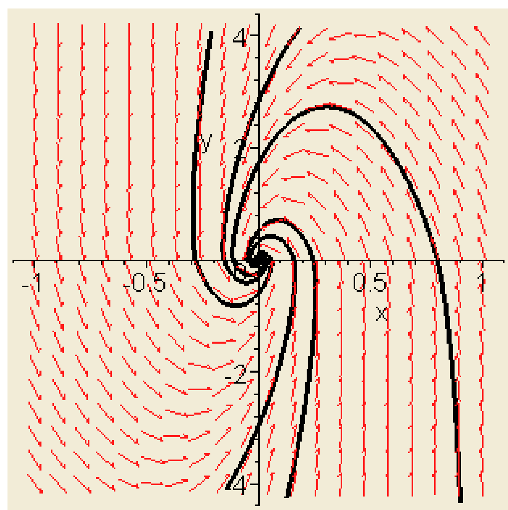

th 6.1,fig1: $\varepsilon = 1, N = 2, p = 3, \alpha = -2$




th 6.1,fig2: $\varepsilon = 1, N = 2, p = 3, \alpha = 1$




th 6.1,fig3: $\varepsilon = 1, N = 2, p = 3, \alpha = 2$




th 6.1,fig4: $\varepsilon = 1, N = 2, p = 3, \alpha = 50$

**Proof.** All the solutions $w$ have a finite number of simple zeros, from Proposition 3.7 and Theorem 3.9. Either they have a compact support. Or $y$ has a strict constant sign and is monotone near $\infty$, and converge to $(0, 0)$ at $\infty$, and (4.4) or (4.5) holds, from Propositions 3.8, 4.1.

In the phase plane $(y, Y)$, system $(\mathbf{S})$ admits only one stationary point $(0, 0)$. The trajectory $\mathcal{T}_r$ starts in $\mathcal{Q}_4$ when $\alpha < 0$, in $\mathcal{Q}_1$ when $\alpha > 0$, and $\lim_{\tau \to -\infty} y = \infty$, with an asymptotical direction of slope $\alpha/N$. From Propositions 4.1 and 4.2 all the nonregular solutions $\pm w$ satisfy (4.6), (4.8), (4.7), (4.9) or (4.10) near $-\infty$. The existence of solutions of any kind is proved at Theorems 5.1 and 5.2. When $p \leqq N$, they correspond to trajectories $\pm\mathcal{T}_\eta$ such that $\mathcal{T}_\eta$ starts in $\mathcal{Q}_1$ with an infinite slope, in any case above $\mathcal{T}_r$. When $p > N$, there is a unique trajectory $\mathcal{T}_u$ satisfying (4.7), starting in $\mathcal{Q}_4$, under $\mathcal{T}_r$; the trajectories $\mathcal{T}_+$ start from $\mathcal{Q}_1$, above $\mathcal{T}_r$; the trajectories $\mathcal{T}_-$ start in $\mathcal{Q}_4$ under $\mathcal{T}_r$. From Theorem 3.9, there exists a unique trajectory $\mathcal{T}_\varepsilon$ converging to $(0,0)$ in $\mathcal{Q}_1$ at $\infty$, with the slope 1.

(1) Case $\alpha < N$. From Proposition 3.6, all the solutions $w$ have at most one simple zero.

The regular solutions stay positive, and $\mathcal{T}_r$ stays in its quadrant, $\mathcal{Q}_4$ or $\mathcal{Q}_1$, from Remark 2.3 (see figures 1 and 2). Then $\mathcal{T}_\varepsilon$ stays in $\mathcal{Q}_1$, because it cannot meet $\mathcal{T}_r$ for $\alpha > 0$, or the line $\{Y = 0\}$ for $\alpha < 0$, from Remark 2.3; and the corresponding $w$ is of type (i).

Consider any trajectory $\mathcal{T}_{[P]}$ with $P \in \mathcal{Q}_1$ above $\mathcal{T}_\varepsilon$. It cannot stay in $\mathcal{Q}_1$ because it does not meet $\mathcal{T}_\varepsilon$ and converges to $(0,0)$ with a slope 0. Thus it enters $\mathcal{Q}_2$ from Remark 2.3. Then $y$ has a unique zero, and $\mathcal{T}_{[P]}$ stays in $\mathcal{Q}_1$ before $P$, and in $\mathcal{Q}_2 \cup \mathcal{Q}_3$ after $P$. Since $\mathcal{T}_{[P]}$ cannot meet $\pm\mathcal{T}_\varepsilon$, and $\lim_{\tau\to\infty} \zeta = \alpha$, $\mathcal{T}_{[P]}$ ends up in $\mathcal{Q}_3$ if $\alpha > 0$, in $\mathcal{Q}_2$ if $\alpha < 0$. It has the same behaviour as $\mathcal{T}_\varepsilon$ at $-\infty$, and $w$ is of type (iii).

Next consider $\mathcal{T}_{[P]}$ for any $P \in \mathcal{Q}_1 \cup\, \mathcal{Q}_4$ between $\mathcal{T}_\varepsilon$ and $\mathcal{T}_r$. Then $y$ stays positive, and $\mathcal{T}_{[P]}$ necessarily starts from $\mathcal{Q}_1$, and $w$ is of type (ii).

At least take any $P \in \mathcal{Q}_1 \cup \mathcal{Q}_4$ under $\mathcal{T}_r$. If $p \leqq N$, $\mathcal{T}_{[P]}$ starts from $\mathcal{Q}_3$ and $y$ has a unique zero, and $-w$ is of type (iii). If $p > N$, either $-w$ is of type (iii), or $\mathcal{T}_{[P]}$ stays in $\mathcal{Q}_4$. From Theorems 5.1, 5.2, either $\mathcal{T}_{[P]}$ coincides with $\mathcal{T}_u$, and $w$ is of type (iv), or with one of the trajectories $\mathcal{T}_-$, thus $w$ is of type (v).

(2) Case $\alpha = N$. All the solutions are given by (1.9), which is equivalent to $J_N \equiv C$, where $J_N$ is defined by (2.1). For $C = 0$, the regular (Barenblatt) solutions, given by (1.10), are nonnegative, with a compact support. In other words the trajectory $\mathcal{T}_\varepsilon$ given by Theorem 5.3 coincides with $\mathcal{T}_r$, it is given by $y \equiv Y$, $y > 0$ (see figure 3). The only change in the phase plane is the nonexistence of solutions of type (ii).

(3) Case $\alpha > N$.

The regular solutions have a number $m \geqq 1$ of simple zeros, from Proposition 3.6 (see figure 4). As above, $\mathcal{T}_r$ starts from $\mathcal{Q}_1$ with a finite slope $\alpha/N$.

Either $\mathcal{T}_r \neq \mathcal{T}_\varepsilon$. Then the regular solutions satisfy $\lim_{r\to\infty} r^\alpha w = L \neq 0$. Since $\mathcal{T}_\varepsilon$ cannot meet $\mathcal{T}_r$, $\mathcal{T}_\varepsilon$ also cuts the line $\{y = 0\}$ at $m$ points, and the corresponding $w$ is of type (vi). For any $P \in \mathcal{Q}_1$ above $\mathcal{T}_r$, the trajectory $\mathcal{T}_{[P]}$ cuts the line $\{y = 0\}$ at $m + 1$ points and $w$ is of type (vii). If $p > N$, there exist trajectories starting from $\mathcal{Q}_1$ between $\mathcal{T}_\varepsilon$ and $\mathcal{T}_r$, with (4.9), such that $w$ has $m$ simple zeros, and trajectories with (4.7) or (4.10), $m$ zeros, and $\lim_{r\to\infty} r^\alpha w = L \neq 0$.

Or $\mathcal{T}_r = \mathcal{T}_\varepsilon$, the regular solutions have a compact support, and we only find solutions of type (vii), (viii). ∎

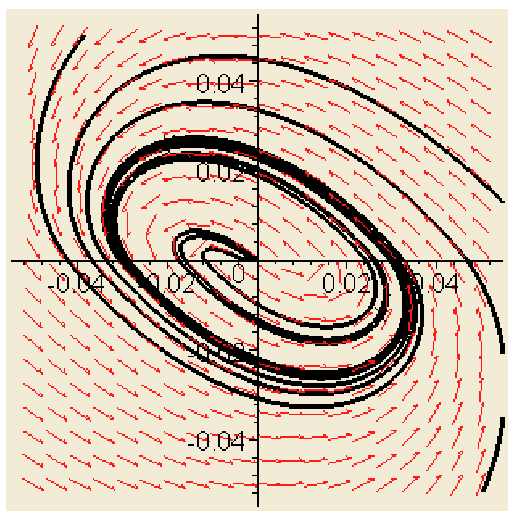
**Remark 6.2** *In the case $\alpha = \eta < 0$, the solutions (iv) are given by (1.11). In the case $N = 1$, $\alpha = -(p-1)/(p-2)$, the solutions of types (i) and (v) are given by (1.14).*

**Remark 6.3** *We conjecture that there exists an increasing sequence* $(\bar{\alpha}_m)$, *with* $\bar{\alpha}_0 = N$ *such that the regular solutions* $w$ *have* $m$ *simple zeros for* $\alpha \in (\bar{\alpha}_{m-1}, \bar{\alpha}_m)$, *with* $\lim_{r\to\infty} r^{\alpha} w = L \neq 0$, *and* $m$ *simple zeros and a compact support for* $\alpha = \bar{\alpha}_m$, *in which case* $\mathcal{T}_r = \mathcal{T}_\varepsilon$.

# 7 The case $\varepsilon = -1, \alpha \leqq -\gamma$

**Theorem 7.1** *Assume* $\varepsilon = -1$, $\alpha \leqq -\gamma$. *Then all the solutions* $w$ *of* $(\boldsymbol{E}_w)$, *among them the regular ones, are ocillating near* $\infty$ *and* $r^{-\gamma} w$ *is asymptotically periodic in* $\ln r$. *There exist*

*(i) solutions such that* $r^{-\gamma} w$ *is periodic in* $\ln r$;
*(ii) a unique solution with a hole;*
*(iii) flat solutions* $w$ *with (4.4) or (4.5) near* $0$;
*(iv) solutions with (4.6) or(4.8) or (4.9) or also (4.10) near* $0$;
*(v) for* $p > N$, *a unique solution with (4.7) near* $0$.




th 7.1,fig5: $\varepsilon = -1, N = 1, p = 3, \alpha = -4$

**Proof.** Here again, $(0, 0)$ is the unique stationary point in the plane $(y, Y)$. Any solution $y$ of $(\mathbf{E}_y)$ oscillates near $\infty$, and $(y, Y)$ is bounded from Proposition 4.3. From the strong form of the Poincaré-Bendixon theorem, see [7, p.239], all the trajectories have a limit cycle or are periodic. In particular $\mathcal{T}_r$ starts in $\mathcal{Q}_1$, since $\varepsilon\alpha > 0$, with the asymptotical direction $\varepsilon\alpha/N$. and it has a limit cycle $\mathcal{O}$. There exists a periodic trajectory of orbit $\mathcal{O}$, thus $w$ is of type (i) (see figure 5).

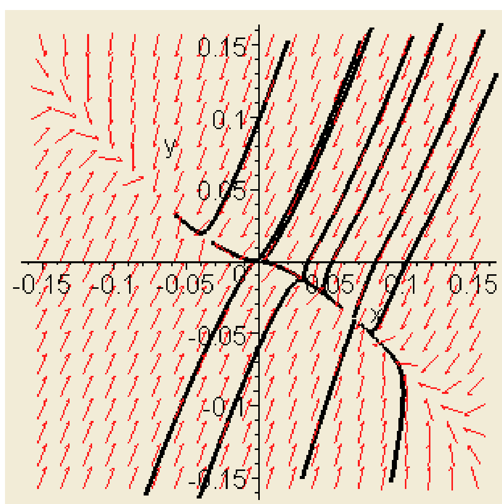
From Theorem 5.2 there exists a unique trajectory $\mathcal{T}_\varepsilon$ starting from $(0, 0)$ with the slope $-1$, $y > 0$; it has a limit cycle $\mathcal{O}_\varepsilon \subset \mathcal{O}$, and $w$ is of type (ii). For any $P$ in the bounded domain delimitated by $\mathcal{O}_\varepsilon$, not located on $\mathcal{T}_\varepsilon$, the trajectory $\mathcal{T}_{[P]}$ does not meet $\mathcal{T}_\varepsilon$, and admits $\mathcal{O}_\varepsilon$ as limit cycle; near $-\infty$, $y$ has a constant sign, is monotone and converges to $(0,0)$ from Propositions 3.8 and 4.1, and $\lim_{\tau\to-\infty}\zeta=\alpha$. This show again the existence of such trajectories, proved at Theorem 5.1, and there is an infinity of them; and $w$ is if type (iii).

From Theorems 5.1 and 5.2, there exist trajectories starting from infinity, with $\mathcal{O}$ as limit cycle, and $w$ is of type (iv) or (v). If $\mathcal{O}=\mathcal{O}_\varepsilon$, all the solutions are described. ■

## 8 Case $\varepsilon = 1, \alpha < -\gamma$.

**Theorem 8.1** *Assume* $\varepsilon=1$, $\alpha<-\gamma$. *Then* $w\equiv\pm\ell r^\gamma$ *is a solution of* ($\boldsymbol{E}_w$). *All regular solutions have a strict constant sign, and satisfy (4.3) near* $\infty$. *Moreover there exist (exhaustively, up to a symmetry)*

*(i) a unique positive flat solution with (4.4) near* $0$ *and (4.3) near* $\infty$;
*(ii) a unique nonnegative solution with (4.6) or (4.8) or (4.9) near* $0$, *and compact support;*
*(iii) positive solutions with the same behaviour near 0 and (4.3) near* $\infty$;
*(iv) solutions with one zero and the same behaviour near 0, and* $|w|$ *satisfies (4.3) near* $\infty$;
*(v) for* $p>N$, *positive solutions with (4.7) near 0 and (4.3) near* $\infty$;
*(vi) for* $p>N$, *positive solutions with (4.10) near 0 and (4.3) near* $\infty$.




th 8.1, fig6: $\varepsilon=1, N=2, p=3, \alpha=-6$

**Proof.** Here system (**S**) admits three stationary points in the plane $(y,Y)$, given at (2.8), thus $w\equiv\pm\ell r^\gamma$ is a solution; and $M_\ell$ is a sink (see figure 6). Any solution $y$ of ($\mathbf{E}_y$) has at most one zero, and is strictly monotone near $\pm\infty$, from Propositions 3.6 and 3.8.

From Theorems 3.9 and 5.3, there exists a unique trajectory $\mathcal{T}_\varepsilon$ converging to $(0,0)$ in $\mathcal{Q}_1$ at $\infty$, and a unique trajectory $\mathcal{T}_\alpha$ converging to $(0,0)$ in $\mathcal{Q}_4$ at $-\infty$. The trajectory $\mathcal{T}_r$ starts in $\mathcal{Q}_4$ with the asymptotical direction $-\left|\alpha\right|/N$. From Remark 2.3, $\mathcal{Q}_4$ is positively invariant, and $\mathcal{Q}_1$ negatively invariant. Then $\mathcal{T}_\varepsilon$ stays in $\mathcal{Q}_1$, and $\mathcal{T}_\alpha$ and $\mathcal{T}_r$ in $\mathcal{Q}_4$. From Proposition 4.1, all the trajectories, apart from $\pm\mathcal{T}_\varepsilon$, converge to $\pm M_\ell$ at $\infty$. Then $\mathcal{T}_r$ converges to $M_\ell$, and $w$ satisfies (4.3) near $\infty$. And $\mathcal{T}_\alpha$ also converges to $M_\ell$, and $w$ is of type (i).

From Propositions 4.1, Theorems 5.1 and 5.2, all the nonregular solutions which are positive near $-\infty$ satisfy (4.6), (4.8), (4.9), (4.10) or (4.7), and there exist such solutions. For $p<N$ (resp. $p=N$), they correspond to trajectories $\mathcal{T}_\eta$ (resp. $\mathcal{T}_+$) starting in $\mathcal{Q}_1$. For $p>N$, there is a unique trajectory $\mathcal{T}_u$ satisfying (4.7), starting in $\mathcal{Q}_4$ under $\mathcal{T}_r$; and the trajectories $\mathcal{T}_+$ satisfying (4.9) start from $\mathcal{Q}_1$; the trajectories $\mathcal{T}_-$ satisfying (4.10) and the unique trajectory $\mathcal{T}_u$ satisfying (4.7) start from $\mathcal{Q}_4$, under $\mathcal{T}_r$. Since $\mathcal{T}_\varepsilon$ stays in $\mathcal{Q}_1$, it defines solutions $w$ of type (ii).

Consider the basis of eigenvectors $(e_1,e_2)$ defined at (2.15), where $\nu(\alpha)>0$, associated to the eigenvalues $\lambda_1<\lambda_2$. One verifies that $\lambda_1<-\gamma<\lambda_2$; thus $e_1$ points towards $\mathcal{Q}_3$ and $e_2$ points towards $\mathcal{Q}_4$. There exist unique trajectories $\mathcal{T}_{e_1}$ and $\mathcal{T}_{-e_1}$ converging to $M_\ell$, tangentially to $e_1$ and $-e_1$. All the other trajectories converging to $M_\ell$ at $\infty$ are tangent to $\pm e_2$. Let

$$\mathcal{M}=\left\{\left|Y\right|^{(2-p)/(p-1)}Y=-\gamma y\right\}, \qquad \mathcal{N}=\left\{(N+\gamma)Y+\varepsilon\left|Y\right|^{(2-p)/(p-1)}Y=\varepsilon\alpha y\right\}$$

be the sets of extremal points of $y$ and $Y$.

The trajectory $\mathcal{T}_r$ starts above the curves $\mathcal{M}$ and $\mathcal{N}$, thus $y'<0$ and $Y'>0$ near $-\infty$. And $\mathcal{T}_r$ converges to $M_\ell$ at $\infty$, tangentially to $e_2$. Indeed if $\mathcal{T}_r=\mathcal{T}_{e_1}$, then $y$ has a minimal point such that $y<\ell$ and $Y<-\left(\gamma\ell\right)^{p-1}$, then $(y,Y)$ cannot be on $\mathcal{M}$. If $\mathcal{T}_r=\mathcal{T}_{-e_1}$, then $Y$ has a maximal point such that $y>\ell$ and $Y<-\left(\gamma\ell\right)^{p-1}$, then also $(y,Y)$ cannot be on $\mathcal{N}$. Finally $\mathcal{T}_r$ cannot end up tangentially to $-e_2$, it would intersect $\mathcal{T}_{e_1}$ or $\mathcal{T}_{-e_1}$.

The trajectory $\mathcal{T}_\alpha$ converge to $M_\ell$ tangentially to $-e_2$. Indeed if $\mathcal{T}_\alpha=\mathcal{T}_{e_1}$, then $Y$ has a maximal point such that $y<\ell$ and $Y<-\left(\gamma\ell\right)^{p-1}$; if $\mathcal{T}_\alpha=\mathcal{T}_{-e_1}$, then $y$ has a maximal point such that $y>\ell$ and $Y>-\left(\gamma\ell\right)^{p-1}$. In any case we reach a contradiction. Moreover $\mathcal{T}_{e_1}$ does not stay in $\mathcal{Q}_4$ : $y$ would have a minimal point such that $y<\ell$ and $Y<-\left(\gamma\ell\right)^{p-1}$, which is impossible; thus $\mathcal{T}_{e_1}$ starts in $\mathcal{Q}_3$, and enters $\mathcal{Q}_4$ at some point $(\xi_1,0)$ with $\xi_1<0$. And $-w$ is of type (iv).

Any trajectory $\mathcal{T}_{[P]}$, with $P$ in the domain of $\mathcal{Q}_1\cup\mathcal{Q}_4$ delimitated by $\mathcal{T}_r,\mathcal{T}_\alpha$ and $\mathcal{T}_\varepsilon$, comes from $\mathcal{Q}_1$, and converges to $M_\ell$ in $\mathcal{Q}_4$, in particular $\mathcal{T}_{-e_1}$; the corresponding $w$ are of type (iii).

Any trajectory $\mathcal{T}_{[P]}$, with $P$ in the domain of $\mathcal{Q}_3\cup\mathcal{Q}_4$ delimitated by $\mathcal{T}_{e_1},\mathcal{T}_\alpha$ and $-\mathcal{T}_\varepsilon$, goes from $\mathcal{Q}_3$ to $\mathcal{Q}_4$, and $\mathcal{T}_{[P]}$ converges to $M_\ell$ at $\infty$, and $-w$ is of type (iv). For any $\xi<\xi_1$, the trajectory $\mathcal{T}_{[(0,\xi)]}$ is of the same type. If $p\leqq N$, any trajectory in the domain under $\mathcal{T}_r$, and $\mathcal{T}_{e_1}$ is of the same type.

If $p > N$, moreover in this domain there exists a the unique trajectory $\mathcal{T}_u$ and trajectories of the type $\mathcal{T}_-$ corresponding to solutions $w$ of type (v) and (vi), from Theorems 5.1 and 5.2. Up to a symmetry, all the solutions are described, and all of them do exist. ■

# 9 Case $\varepsilon = -1, -\gamma < \alpha$

Here again System (S) admits the three stationary points (2.8), thus $w \equiv \pm\ell r^\gamma$ is a solution of $(\mathbf{E}_w)$. The behaviour is very rich: it depends on the position of $\alpha$ with respect to $\alpha^*$ defined at (1.5), and 0, $-p'$, and $\eta$ (in case $p > N$), and also $\alpha_1, \alpha_2$ defined at (2.14). We start from some general remarks.

**Remark 9.1** *(i) There exists a unique trajectory $\mathcal{T}_\varepsilon$ starting from $(0,0)$ in $\mathcal{Q}_4$ with the slope $-1$, from Theorem 3.9.*
*(ii) There exists a unique trajectory $\mathcal{T}_\alpha$ converging to $(0,0)$ at $\infty$, in $\mathcal{Q}_1$ if $\alpha > 0$, in $\mathcal{Q}_4$ if $\alpha < 0$, with a slope $0$ at $(0,0)$, and $\lim_{\tau\to\infty} \zeta = \alpha$, from Theorem 5.3.*
*(iii) From Remark 2.3, if $\alpha > 0$, $\mathcal{Q}_4$ is positively invariant and $\mathcal{Q}_1$ negatively invariant. If $\alpha < 0$, at any point $(0,\xi), \xi < 0$, the vector field points to $\mathcal{Q}_4$, and at any point $(\varphi, 0), \varphi > 0$, it points to $\mathcal{Q}_1$. Thus if $\mathcal{T}_\varepsilon$ does not stay in $\mathcal{Q}_1$, then $\mathcal{T}_\alpha$ stays in the bounded domain delimitated by $\mathcal{Q}_4 \cap \mathcal{T}_\varepsilon$. If $\mathcal{T}_\alpha$ does not stay in $\mathcal{Q}_4$, then $\mathcal{T}_\varepsilon$ stays in the bounded domain delimitated by $\mathcal{Q}_4 \cap \mathcal{T}_\alpha$. If $\mathcal{T}_\varepsilon$ is homoclinic, in other words $\mathcal{T}_\varepsilon = \mathcal{T}_\alpha$, it stays in $\mathcal{Q}_4$.*

**Remark 9.2** *From Propositions 4.1, Theorems 5.1 and 5.2, all the nonregular solutions positive near $-\infty$ satisfy (4.6) for $p < N$, (4.8) for $p = N$, corresponding to trajectories $\mathcal{T}_\eta, \mathcal{T}_+$ starting from$\mathcal{Q}_1$; and (4.9), (4.10) or (4.7) for $p > N$, corresponding to trajectories $\mathcal{T}_+$ starting from $\mathcal{Q}_1$, and $\mathcal{T}_-, \mathcal{T}_u$ starting from $\mathcal{Q}_4$.*

**Remark 9.3** *Any trajectory $\mathcal{T}$ is bounded near $\infty$ from Proposition 4.3. From the strong form of the Poincaré-Bendixon theorem, any trajectory $\mathcal{T}$ bounded at $\pm\infty$ converges to $(0,0)$ or $\pm M_\ell$, or its limit set $\Gamma_\pm$ at $\pm\infty$ is a cycle, or it is homoclinic, namely $\mathcal{T}_\varepsilon = \mathcal{T}_\alpha$. If there exists a limit cycle surrounding $(0,0)$, it also surrounds the points $\pm M_\ell$, from Proposition 3.8.*

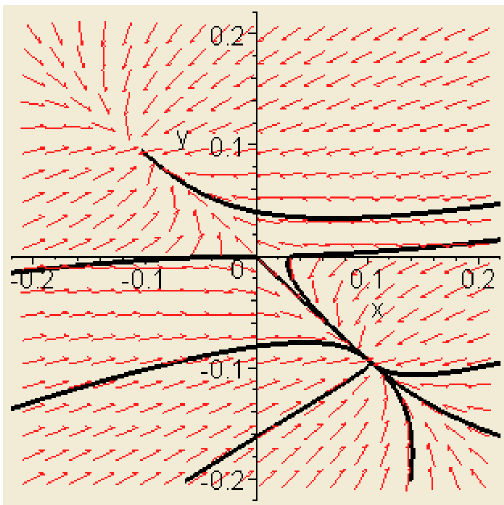
The simplest case is $\alpha > 0$.

**Theorem 9.4** *Assume $\varepsilon = -1$, $\alpha > 0$.*

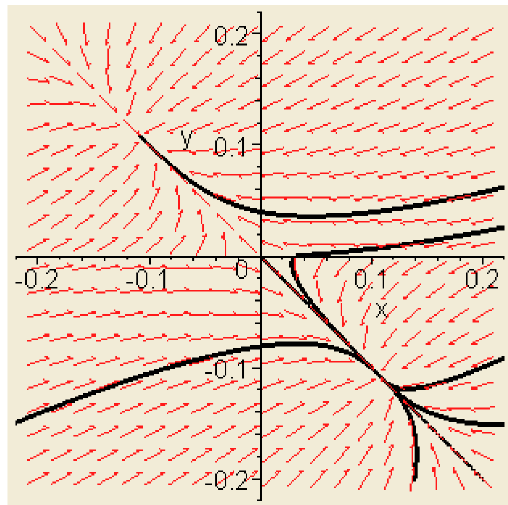
*Then $w \equiv \ell r^\gamma$ is a solution $w$ of $(\boldsymbol{E}_w)$. All regular solutions have a strict constant sign; and satisfy (4.3) near $\infty$. There exist (exhaustively, up to a symmetry)*

*(i) a unique nonnegative solution with a hole, and (4.3) near $\infty$;*
*(ii) a unique positive solution with (4.6), or (4.8) or (4.9), and (4.4) near $\infty$;*
*(iii) positive solutions with the same behaviour near $0$, and (4.3) near $\infty$;*

*(iv) solutions with one zero, the same behaviour near* $0$*, and* $|w|$ *satisfies (4.3) near* $\infty$*;*
*(v) for* $p > N$*, a unique positive solution with (4.7) near* $0$*, and (4.3) near* $\infty$*;*
*(vi) for* $p > N$*, positive solutions with (4.10) near* $0$*, and (4.3) near* $\infty$*.*




th 9.4, fig7: $\varepsilon = -1, N = 1, p = 3, \alpha = 0.7$




th 9.4, fig8: $\varepsilon = -1, N = 1, p = 3, \alpha = 1$

**Proof.** Any solution $y$ of $(\mathbf{E}_y)$ has at most one zero, and $y$ is strictly monotone near $\infty$, from Propositions 3.6 and 4.4. The point $M_\ell$ is a sink and a node point, since $\alpha > 0 \geqq \alpha_2$ (see figure 7). Consider the basis eigenvectors $(e_1,\, e_2)$, defined at (2.15), where $\nu(\alpha) < 0$, associated to the eigenvalues $\lambda_1 < \lambda_2 < 0$. One verifies that $\lambda_1 < -\gamma < \lambda_2$, thus $e_1$ points towards $\mathcal{Q}_3$ and $e_2$ points towards $\mathcal{Q}_4$. There exist unique trajectories $\mathcal{T}_{e_1}$ and $\mathcal{T}_{-e_1}$ tangent to $e_1$ and $-e_1$ at $\infty$. All the other trajectories which converge to $M_\ell$ end up tangentially to $\pm e_1$.

The trajectory $\mathcal{T}_\alpha$ stays in $\mathcal{Q}_1$ from Remark 9.1; near $-\infty$ it is of type $\mathcal{T}_\eta$ for $p < N$, and $\mathcal{T}_+$ for $p \geqq N$; it defines the solution of type (ii). Since $\mathcal{T}_\alpha$ is the unique trajectory converging to $(0,0)$ at $\infty$, all the trajectories, apart from $\pm\mathcal{T}_\alpha$, converge to $\pm M_\ell$ at $\infty$, from Propositions 3.8 and 4.1.

The trajectories $\mathcal{T}_r$ and $\mathcal{T}_\varepsilon$ start in $\mathcal{Q}_4$, and stay in it from Remark 9.1, and both converge to $M_\ell$ at $\infty$, then $w$ satisfies (4.3); and $\mathcal{T}_r$ starts with the asymptotical direction $-\alpha/N$. And $\mathcal{T}_\varepsilon$ defines the solution of type (i).

As in the proof of Theorem 8.1, $\mathcal{T}_r$ ends up tangentially to $e_2$, and $\mathcal{T}_\varepsilon$ tangentially to $-e_2$. Moreover $\mathcal{T}_{e_1}$ does not stay in $\mathcal{Q}_4$, it starts in $\mathcal{Q}_3$, and converges to $M_\ell$ in $\mathcal{Q}_4$, and $-w$ is of type (iv). Any trajectory $\mathcal{T}_{[P]}$, with $P$ in the domain of $\mathcal{Q}_4$ between $\mathcal{T}_{e_1}$, $\mathcal{T}_\varepsilon$, starts from $\mathcal{Q}_3$, enters $\mathcal{Q}_4$ at some point $(0,\xi), \xi > \xi_1$, and has the same type as $\mathcal{T}_{e_1}$. Any trajectory $\mathcal{T}_{[(0,\xi)]}$ with $\xi < \xi_1$ is of the same type.

Any trajectory $\mathcal{T}_{[P]}$, with $P$ in the domain of $\mathcal{Q}_1 \cup \mathcal{Q}_4$ above $\mathcal{T}_r \cup \mathcal{T}_\varepsilon$, starts from $\mathcal{Q}_1$, and converges to $M_\ell$ in $\mathcal{Q}_4$, in particular $\mathcal{T}_{-e_1}$; the corresponding $w$ are of type (iii). If $p \leqq N$, all the solutions are described. If $p > N$, moreover there exist trajectories staying in $\mathcal{Q}_4 : \mathcal{T}_u$ and the $\mathcal{T}_-$, starting under $\mathcal{T}_r$, corresponding to types (v) and (vi). ∎

**Remark 9.5** *For $\alpha = N$, $\mathcal{T}_r$ and $\mathcal{T}_\varepsilon$ are given by (1.10), respectively with $K > 0$ and $K < 0$. The trajectory $\mathcal{T}_\varepsilon$ describes the portion $0M_\ell$ of the line $\{Y = -y\}$, and $\mathcal{T}_r$ the complementary half-line in $\mathcal{Q}_4$ (see figure 8).*

Next we assume $-p' \leqq \alpha < 0$. The case $p > N$ is delicate: indeed the special value $\alpha = \eta$ is involved, because $\eta < 0$.

**Theorem 9.6** *Assume $\varepsilon = -1, p \leqq N$, and $-p' \leqq \alpha < 0$. Then $w \equiv \ell r^\gamma$ is a solution $w$ of $(\boldsymbol{E}_w)$.*

*There exist a unique nonnegative solution with a hole, satisfying (4.3) at $\infty$.*

*(1) If $\alpha \neq -p'$, all regular solutions have one zero, and $|w|$ satisfies (4.3) near $\infty$. There exist (exhaustively, up to a symmetry)*

• *for $p \leqq N$,*

*(i) a unique solution with one zero, with (4.6) or (4.8) near $0$,and (4.4) near $\infty$;*
*(ii) solutions with one zero, with (4.6) or (4.8) near $0$, and $|w|$ satisfies (4.3) near $\infty$;*
*(iii) solutions with two zeros, with (4.6) or (4.8) near $0$, and (4.3) near $\infty$;*

• *for $p > N$, $\eta < \alpha$,*

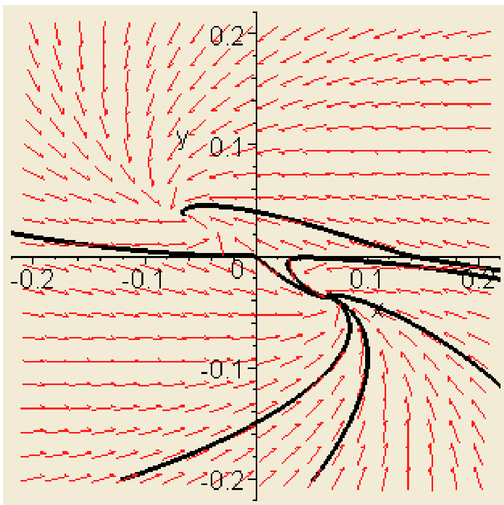
*(iv) a unique positive solution, with (4.10) near $0$, and (4.4) near $\infty$;*
*(v) a unique positive solution, with (4.7) near $0$, and (4.3) near $\infty$;*
*(vi) positive solutions, with (4.10) near $0$, and (4.3) near $\infty$;*
*(vii) solutions with one zero with (4.10) or (4.9) near $0$, and (4.3) near $\infty$;*

• *for $p > N, \alpha < \eta$,*

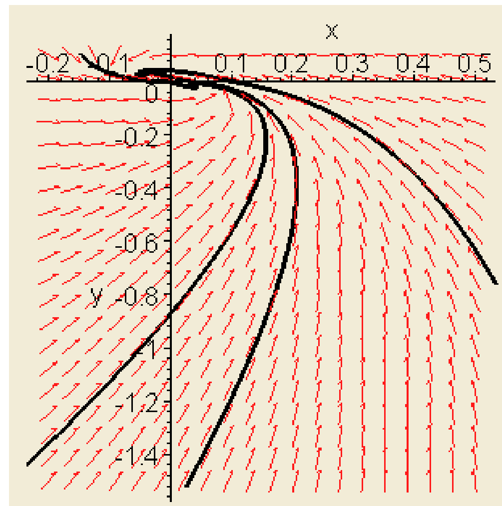
*(viii) a unique solution with one zero, with (4.9) near $0$, and (4.4) near $\infty$;*
*(ix) a unique solution with one zero, with (4.7) near $0$, and $|w|$ satisfies (4.3) near $\infty$;*
*(x) solutions with one zero, with (4.9) or (4.9) near $0$, and $|w|$ satisfies (4.3) near $\infty$;*
*(xi) solutions with two zeros, with (4.9) near $0$, and (4.3) near $\infty$.*

• *for $p > N, \alpha = \eta$, solutions of the form $w = cr^{|\eta|}$ $(c > 0)$. The other solutions are of type (vii).*

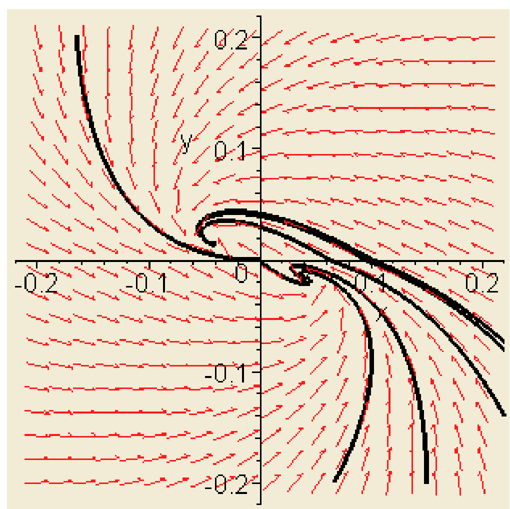
*(2) If $\alpha = -p'$, all regular solutions have one zero and satisfy (4.4) near $\infty$. The solutions without hole are of types (ii), (iii) for $p \leqq N$, (ix), (x), (xi) for $p > N$.*




th 9.6, fig9: $\varepsilon = -1, N = 1, p = 3, \alpha = -0.7$




th 9.6, fig10: $\varepsilon = -1, N = 1, p = 3, \alpha = -1.49$




th 9.6, fig11: $\varepsilon = -1, N = 1, p = 3, \alpha = -3/2$

**Proof.** Here again $M_\ell$ is a sink; but it is a node point only if $\alpha \geqq \alpha_{2.}$. The phase plane $(y, Y)$ does not contain any cycle, from Proposition 4.4. From Proposition 3.6, any solution $y$ has at most two zeros, and $Y$ at most one.

The unique trajectory $\mathcal{T}_\alpha$ ends up in $\mathcal{Q}_4$ with the slope 0. From the uniqueness of $\mathcal{T}_\alpha$ and $\mathcal{T}_\varepsilon$, all the trajectories, apart from $\pm\mathcal{T}_\alpha$, converge to $\pm M_\ell$ at $\infty$, from Proposition 4.1 and Remark 9.3. Since $\varepsilon\alpha > 0$, the trajectory $\mathcal{T}_r$ starts in $\mathcal{Q}_1$, and $y$ has at most one zero. Then $\mathcal{T}_r$ converges to $-M_\ell$ in $\mathcal{Q}_2$, or $\mathcal{T}_r = -\mathcal{T}_\alpha$.

The trajectory $\mathcal{T}_\varepsilon$ starts in $\mathcal{Q}_4$ with the slope $-1$, satisfies $y \geqq 0$ from Proposition 3.6. If $\mathcal{T}_\varepsilon$ converge to $(0,0)$, then $\mathcal{T}_\varepsilon = \mathcal{T}_\alpha$, thus it is homoclinic. Then $M_\ell$ is in the bounded component defined by $\mathcal{T}_\varepsilon$, and $\mathcal{T}_\varepsilon$ meets $\mathcal{T}_r$, which is impossible. Hence $\mathcal{T}_\varepsilon$ converges to $M_\ell$ in $\mathcal{Q}_4$, and $w$ is nonnegative with a hole and satisfies (4.3) near $\infty$.

If $\alpha \neq -p'$, we claim that $\mathcal{T}_r \neq -\mathcal{T}_\alpha$. Indeed suppose $\mathcal{T}_r = -\mathcal{T}_\alpha$. Consider the functions $y_\alpha, Y_\alpha$, defined by (2.3) with $d = \alpha$. Then $Y_\alpha$ stays positive, and $Y_\alpha = O(e^{(\alpha(p-1)+p)\tau})$ at $\infty$, thus

$$\lim_{\tau\to\infty} Y_\alpha = 0, \quad \lim_{\tau\to\infty} Y_\alpha = c > 0, \quad \lim_{\tau\to-\infty} y_\alpha = \infty, \quad \lim_{\tau\to\infty} y_\alpha = L < 0.$$

Moreover $y_\alpha, Y_\alpha$ have no extremal point: at such a point, from (3.2), (3.3) the second derivatives have a strict constant sign; then $Y'_\alpha > 0 > y'_\alpha$. If $\alpha < \eta$ (in particular if $p \leqq N$), from (4.13), near $\infty$,

$$(p-1)Y''_\alpha/Y'_\alpha \geqq |Y|^{(2-p)/(p-1)}\,(1+o(1)),$$

thus $Y''_\alpha > 0$ near $\infty$, which is contradictory; if $\alpha > \eta$, from (4.12)

$$(p-1)y''_\alpha/y'_\alpha \geqq |Y|^{(2-p)/(p-1)}\,(1+o(1)),$$

thus $y''_\alpha < 0$ near $\infty$, still contradictory. If $\alpha = \eta$, $\mathcal{T}_\alpha = \mathcal{T}_u$ from (1.11), thus again $\mathcal{T}_r \neq -\mathcal{T}_\alpha$.

If $p > N$ and $\alpha \neq \eta$, we claim that $\mathcal{T}_\alpha \neq \mathcal{T}_u$. Indeed suppose $\mathcal{T}_\alpha = \mathcal{T}_u$. This trajectory stays $\mathcal{Q}_4$, the function $\zeta$ stays negative, and $\lim_{\tau\to-\infty}\zeta = \eta$, $\lim_{\tau\to\infty}\zeta = \alpha$. If $\zeta$ has an extremal point $\vartheta$, then $\vartheta \in (\alpha, \eta)$ from System (**Q**), and $\zeta''$ has a constant sign, the sign of $\alpha - \zeta$; it is impossible. Thus $\zeta$ is monotone; then $(\alpha - \eta)\zeta' > 0$, which contradicts System (**Q**).

(1) Case $\alpha \neq -p'$. Since $\mathcal{T}_r \neq -\mathcal{T}_\alpha$, $\mathcal{T}_r$ converges to $-M_\ell$, and $y$ has one zero, and $|w|$ satisfies (4.3).

• Case $p \leqq N$. All the other trajectories start in $\mathcal{Q}_3$ or $\mathcal{Q}_1$, from Remarks 9.1 and 9.2. For any $\varphi > 0$, the trajectory $\mathcal{T}_{[(\varphi,0)]}$ goes from $\mathcal{Q}_4$ into $\mathcal{Q}_1$, and converges to $-M_\ell$ in $\mathcal{Q}_2$, since it cannot meet $\mathcal{T}_r$ and $-\mathcal{T}_\varepsilon$; thus $y$ has two zeros, and $w$ is of type (iii). The trajectory $\mathcal{T}_\alpha$ cannot meet $\mathcal{T}_{[(\varphi,0)]}$, thus $y$ has one zero, and it has the same behaviour at $-\infty$, and $w$ is of type (i). All the trajectories $\mathcal{T}_{[P]}$ with $P$ in the interior domain of $\mathcal{Q}_1$ delimitated by $-\mathcal{T}_\varepsilon$ and $\mathcal{T}_r$ start from $\mathcal{Q}_1$ and converge to $-M_\ell$, $y$ has precisely one zero, and has the same behaviour at $-\infty$, and $w$ is of type (ii).

• Case $p > N$, $\eta < \alpha$ (see figure 9). Any solution $y$ has at most one simple zero. The trajectory $\mathcal{T}_\alpha$ stays in $\mathcal{Q}_4$. Indeed if it started in $\mathcal{Q}_3$, then for any trajectory $\mathcal{T}_{[(0,\xi)]}$ with $(0,\xi)$ above $-\mathcal{T}_\alpha$, the function $y$ would have two zeros. Since $\mathcal{T}_\alpha \neq \mathcal{T}_u$, we have $\mathcal{T}_\alpha \in \mathcal{T}_-$, and $w$ is of type (iv). The trajectory $\mathcal{T}_u$ necessarily stays in $\mathcal{Q}_4$ and converges to $M_\ell$, and $w$ is of type (v). The trajectories $\mathcal{T}_{[P]}$, with $P$ in the domain delimitated by $\mathcal{T}_u, \mathcal{T}_\alpha$ and $\mathcal{T}_\varepsilon$, are of type $\mathcal{T}_-$ and converge in $\mathcal{Q}_4$ to $M_\ell$, and $w$ is of type (vi). The trajectories $\mathcal{T}_{[P]}$, with $P$ in the domain delimitated by $\mathcal{T}_r, \mathcal{T}_\alpha$ and $-\mathcal{T}_\varepsilon$, are of type $\mathcal{T}_-$, and converge to $-M_\ell$, and $y$ has one zero. The trajectories $\mathcal{T}_{[P]}$, with $P$ in the domain delimitated by $\mathcal{T}_r$ and $-\mathcal{T}_u$, are of type $\mathcal{T}_+$,converge to $-M_\ell$, and $y$ has one zero. Both define solutions $w$ of type (vii).

• Case $p > N, \alpha < \eta$ (see figure 10). We have seen that $\mathcal{T}_r \neq -\mathcal{T}_\alpha$. If $\mathcal{T}_\alpha \in \mathcal{T}_+$, then $\zeta$ decreases from 0 to $\alpha$, which contradicts System (**Q**) at $\infty$. Then $\mathcal{T}_\alpha$ does not stay in $\mathcal{Q}_4$, it starts in $\mathcal{Q}_3$ and $-\mathcal{T}_\alpha \in \mathcal{T}_-$, hence $y$ has a zero, and $w$ is of type (viii). Then $\mathcal{T}_u$ and the trajectories $\mathcal{T}_-$ converge to $-M_\ell$, and $y$ has one zero. The trajectories $\mathcal{T}_{[P]}$, with $P$ in the domain delimitated by $\mathcal{T}_r, -\mathcal{T}_\alpha$ and $-\mathcal{T}_\varepsilon$, are of type $\mathcal{T}_+$ and converge to $-M_\ell$, $y$ has one zero. They correspond to $w$ is of type (ix) or (x). The trajectories $\mathcal{T}_{[P]}$, with $P$ in $\mathcal{Q}_4$ above $\mathcal{T}_r$, cut the line $\{y = 0\}$ twice, and converge to $M_\ell$, and $w$ is of type (xi).

• Case $p > N, \alpha = \eta$. Then $\mathcal{T}_\alpha = \mathcal{T}_u$, the functions $w = cr^{-\eta}$ $(c > 0)$ are particular solutions. The phase plane study is the same, and gives only solutions of type (vii).

(2) Case $\alpha = -p'$ (see figure 11). Here $\mathcal{T}_r = -\mathcal{T}_\alpha$, since the regular solutions are given by (1.12). Thus there exist no more solutions of type (ii) or (viii). ■

Next we study the behaviour of all the solutions when $\alpha < -p'$. In particular we prove the existence and uniqueness of an $\alpha_c$ for which there exists an homoclinic trajectory. Thus we find again some results obtained in [8], with new detailed proofs. We also improve the bounds for $\alpha_c$, in particular $\alpha^* < \alpha_c$.

**Lemma 9.7** *Let*

$$\alpha_p := -(p-1)/(p-2).$$

*If $N = 1$, for $\alpha = \alpha_p$, then there exists an homoclinic trajectory in the phase plane $(y, Y)$. If $N \geqq 2$, for $\alpha = \alpha_p$, there is no homoclinic trajectory, moreover $\mathcal{T}_\alpha$ converges to $M_\ell$ at $-\infty$ or has a limit cycle in $\mathcal{Q}_4$.*

**Proof.** In the case $N = 1$, $\alpha = \alpha_p$, the explicit solutions (1.14) define an homoclinic trajectory in the phase plane $(y, Y)$, namely $\mathcal{T}_\varepsilon = \mathcal{T}_\alpha$. In the phase plane $(g, s)$ of System **(R)**, from Remark 2.6, they correspond to the line $s \equiv 1 + \alpha g$, joining the stationary points $(0, 1)$ and $(-1/\alpha, 0)$.

Next assume $N \geqq 2$ and consider the trajectory $\mathcal{T}_\alpha$ in the plane $(y, Y)$. In the plane $(g, s)$ of System (**R**), the corresponding trajectory $\mathcal{T}'_\alpha$ ends up at $(-1/\alpha, 0)$, as $\nu$ tends to $\infty$ from (2.18), with the slope $-k_p$. If $\mathcal{T}_\alpha$ is homoclinic, then $\mathcal{T}'_\alpha$ converges to $(0, 1)$ as $\nu$ tends to $-\infty$. Consider the segment

$$T = \{(g, -k(g + 1/\alpha_p) : g \in [0, 1/|\alpha_p|]\}, \quad \text{with} \quad k = p'\alpha_p^2/(N + 2/(p-2)) > k_p.$$

Its extremity $(0, k/|\alpha_p|)$ is strictly under $(0, 1)$. The domain $\mathcal{R}$ delimitated by the axes, which are particular orbits, and $T$, is negatively invariant: indeed, at any point of $T$, we find

$$k\frac{dg}{d\nu} + \frac{ds}{d\nu} = (N-1)p'ks(g - \frac{1}{\gamma})^2.$$

The trajectory $\mathcal{T}'_\alpha$ ends up in $\mathcal{R}$, thus it stays in it, hence $\mathcal{T}'_\alpha$ cannot join $(0,1)$. In the phase plane $(y,Y)$, $\mathcal{T}_\alpha$ is not homoclinic, and $\mathcal{T}_\alpha$ stays in $\mathcal{Q}_4$, and Remark 9.3 applies. ■

**Remark 9.8** *Notice that* $\alpha^* \leqq \alpha_p \Leftrightarrow N \leqq p$.

**Theorem 9.9** *Assume* $\varepsilon = -1$, *and* $\alpha < -p'$. *There exists a unique* $\alpha_c < 0$ *such that there exists an homoclinic trajectory in the plane* $(y,Y)$; *in other words* $\mathcal{T}_\varepsilon = \mathcal{T}_\alpha$. *If* $N = 1$, *then* $\alpha_c = \alpha_p$. *If* $N \geqq 2$, *then*

$$max(\alpha^*, \alpha_p) < \alpha_c < \min(\alpha_2, -p'). \tag{9.1}$$

**Proof.** In order to prove the existence of an homoclinic orbit for System (**S**), we could consider a Poincaré application as in [4], but it does not give uniqueness. Thus we consider the system $(\mathbf{R}_\beta)$ obtained from $(\mathbf{R})$ by setting $s = \beta S$:

$$\left.\begin{array}{lll} \frac{dg}{d\nu} = gF(g,S), & & F(g,S) := \beta S(1+\eta g) - \frac{1}{p-1}(1+\alpha g), \\ & & \\ \frac{dS}{d\nu} = SG(g,S), & & G(g,S) := 1 + \alpha g - \beta(1+Ng)S. \end{array}\right\} \qquad (\mathbf{R}_\beta)$$

Its stationary points are

$$(0,0), \qquad A' = (1/|\alpha|, 0), \qquad B' = (0, 1/\beta), \qquad M' = (1/\gamma, 1/(N+\gamma)(p-2)),$$

where $M'$ corresponds to $M_\ell$. The existence of homoclinic trajectory for System (**S**) resumes to the existence of a trajectory for System $(\mathbf{R}_\beta)$ in the plane $(g,S)$, starting from $B'$ and ending at $A'$.

(i) *Existence*. We can assume that $\alpha \in (\alpha_1, \min(\alpha_2, -p'))$, from Proposition 4.4. In the plane $(g,S)$, consider the trajectories $\mathcal{T}'_\varepsilon$ and $\mathcal{T}'_\alpha$ corresponding to $\mathcal{T}_\varepsilon \cap \mathcal{Q}_4$ and $\mathcal{T}_\alpha \cap \mathcal{Q}_4$ in the plane $(y,Y)$. Then $\mathcal{T}'_\varepsilon$ starts from $B'$ and $\mathcal{T}'_\alpha$ ends up at $A'$. From Remark 9.1, for any $\alpha \in (\alpha_1, \alpha_2)$, with $\alpha \leqq -p'$, we have three possibilities:

• $\mathcal{T}'_\varepsilon$ is converging to $M'$ as $\nu$ tends to $\infty$ and turns around this point, since $\alpha$ is a spiral point, or it has a limit cycle in $\mathcal{Q}_1$ around $M'$. And $\mathcal{T}'_\alpha$ admits the line $g = 0$ as an asymptote as $\nu$ tends to $-\infty$, which means that $\mathcal{T}_\alpha$ does not stay in $\mathcal{Q}_4$ in the plane $(y,Y)$. Then $\mathcal{T}'_\varepsilon$ meats the line

$$L := \{g = 1/\gamma\}$$

at a first point $(1/\gamma, S_0(\alpha))$. And $\mathcal{T}'_\alpha$ meats $L$ at a last point $(1/\gamma, S_1(\alpha))$, such that $S_0(\alpha) - S_1(\alpha) < 0$;

• $\mathcal{T}'_\alpha$ is converging to $M'$ at $-\infty$ or it has a limit cycle in $\mathcal{Q}_1$ around $M'$. And $\mathcal{T}'_\varepsilon$ admits the line $S = 0$ as an asymptote at $\infty$, which means that $\mathcal{T}_\varepsilon$ does not stay in $\mathcal{Q}_4$. Then with the same notations, $S_0(\alpha) - S_1(\alpha) > 0$.

• $\mathcal{T}'_\varepsilon = \mathcal{T}'_\alpha$, equivalently $S_0(\alpha) - S_1(\alpha) = 0$.

The function $\alpha \mapsto \varphi(\alpha) = S_0(\alpha) - S_1(\alpha)$ is continuous, from Theorems 3.9 and 5.3. If $-p' < \alpha_2$, then $\varphi(-p')$ is well defined and $\varphi(-p') < 0$; indeed $\mathcal{T}_\alpha = -\mathcal{T}_r$, thus $\mathcal{T}_\alpha$ does not stay in $\mathcal{Q}_4$ from Theorem 9.6. If $\alpha_2 \leqq -p'$, in the plane $(y, Y)$, the trajectory $\mathcal{T}_{\alpha_2}$ leaves $\mathcal{Q}_4$, from Proposition 4.4, because $\alpha_2$ is a sink, and transversally from Remark 9.1. The same happens for $\mathcal{T}_{\alpha_{2-\upsilon}}$ for $\upsilon > 0$ small enough, by continuity, thus $\varphi(\alpha_2 - \upsilon) < 0$. From Lemma 9.7, $\varphi(\alpha_p) > 0$ if $N \geqq 2$, and $\varphi(\alpha_p) = 0$ if $N = 1$. In any case there exists at least an $\alpha_c$ satisfying (9.1), such that $\varphi(\alpha_c) = 0$.

(ii) *Uniqueness*. First observe that $1 + \eta g > 0$; indeed $1 + \eta/|\alpha| > (p' + \eta)/|\alpha| > 0$. Now

$$(p-1)F + G = p\beta S(1/\gamma - g) = (p-2)\beta S(1 - \gamma g),$$

hence the curves $\{F = 0\}$ and $\{G = 0\}$ intersect at $M'$ and $A'$, $\{G = 0\}$ contains $B'$ and is above $\{F = 0\}$ for $g \in (0, 1/\gamma)$ and under it for $g \in (1/\gamma, 1/|\alpha|)$. Moreover $\mathcal{T}'_\varepsilon$ has a negative slope at $B'$, thus $F > 0 > G$ near 0 from $(\mathbf{R}_\beta)$. And $\mathcal{T}'_\varepsilon$ cannot meet $\{G = 0\}$ for $(0, 1/\gamma)$, because on this curve the vector field is $(gF, 0)$ and $F > 0$. Thus $\mathcal{T}'_\varepsilon$ satisfies $F > 0 > G$ on $(0, 1/\gamma)$. In the same way $\mathcal{T}'_\alpha$ has a negative slope $-\theta\alpha^2/(p-1)(\eta + |\alpha|) < 0$ at $1/|\alpha|$, thus $F > 0 > G$ near $1/|\alpha|$. And $\mathcal{T}'_\alpha$ cannot meet $\{F = 0\}$, because the vector field on this curve is $(0, SG)$ and $G < 0$. Thus $\mathcal{T}'_\alpha$ satisfies $F > 0 > G$ on $(1/\gamma, 1/|\alpha|)$.

Let $\alpha < \bar{\alpha}$. Then $\mathcal{T}'_\varepsilon$ is above $\bar{\mathcal{T}}'_\varepsilon$ near $g = 0$, and $\mathcal{T}'_\alpha$ is at the left of $\mathcal{T}'_{\bar{\alpha}}$ near $S = 0$. We show that $\varphi(\alpha) > \varphi(\bar{\alpha})$. First suppose that $\mathcal{T}'_\varepsilon$ and $\bar{\mathcal{T}}'_\varepsilon$ (or $\mathcal{T}'_\alpha$ and $\bar{\mathcal{T}}'_{\bar{\alpha}}$) intersect at a first point $P_1$ (or a last point) such $g \neq 1/\gamma$. Then at this point

$$\frac{1}{p-1}\frac{g}{S}\frac{dS}{dg} + 1 = \frac{(p-2)(1-\gamma g)S}{(p-1)S(1+\eta g) - \beta^{-1}(1+\alpha g)} = \frac{(p-2)(1-\gamma g)S}{h_S(g) - \beta^{-1}(1-\gamma g)} \tag{9.2}$$

with $h_S(g) = (p-1)S(1+\eta g) - g/(p-2)$. Thus the denominator, which is positive, is increasing in $\alpha$ on $(0, 1/\gamma)$, decreasing on $(1/\gamma, 1/|\alpha|)$; in any case $dS/dg > d\overline{S}/dg$ at $P_1$, which is contradictory. Next suppose that there is an intersection on $L$. At such a point $P_1 = (1/\gamma, S_1) = (1/\gamma, \overline{S}_1)$ the derivatives are equal from (9.2), and $P_1$ is above $M'$, because $F > 0$. At any points $(g, S(g)) \in \mathcal{T}'_\varepsilon$ (or $\mathcal{T}'_\alpha$), $(g, \overline{S}(g)) \in \bar{\mathcal{T}}'_\varepsilon$ (or $\bar{\mathcal{T}}'_{\bar{\alpha}}$), setting $g = 1/\gamma + u$,

$$\Phi(u) = (\frac{1}{p-1}\frac{g}{S}\frac{dS}{dg} + 1)\frac{1}{(p-2)S} = -\frac{\gamma}{h_S(1/\gamma)}u + \frac{1}{h_S^2(1/\gamma)}(\frac{\gamma}{\beta} + h'_S(1/\gamma))u^2(1 + o(1)),$$

$$\bar{\Phi}(u) = (\frac{1}{p-1}\frac{g}{\overline{S}}\frac{d\overline{S}}{dg} + 1)\frac{1}{(p-2)\overline{S}} = -\frac{\gamma}{h_{\overline{S}}(1/\gamma)}u + \frac{1}{h_{\overline{S}}^2(1/\gamma)}(\frac{\gamma}{\beta} + h'_{\overline{S}}(1/\gamma))u^2(1 + o(1)),$$

And $h_S(1/\gamma) = h_{\overline{S}}(1/\gamma) > 0$, and $h'_S(1/\gamma) = h'_{\overline{S}}(1/\gamma)$, then

$$(\Phi - \bar{\Phi})(u) = \frac{\gamma u^2(1/\beta - 1/\bar{\beta})}{h(1/\gamma)}(1 + o(1)).$$

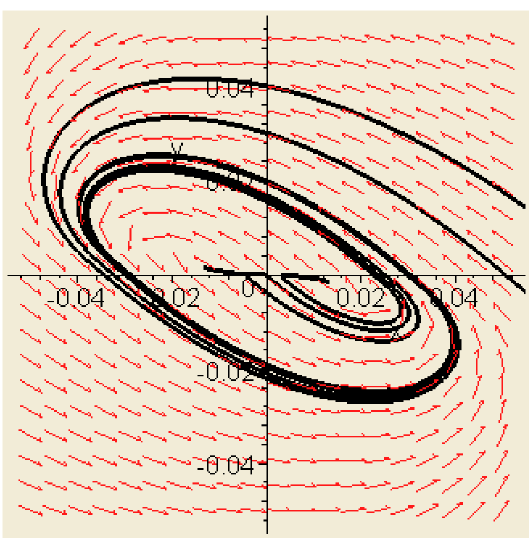
This implies $d^2(S-\overline{S})/dg^2 = 0$ and $d^3(S-\overline{S})/dg^3 = 2S_1\gamma^2(p-1)(p-2)(1/\beta - 1/\bar{\beta}) > 0$, which is a contradiction. Then $\mathcal{T}'_\varepsilon$ and $\bar{\mathcal{T}}'_\varepsilon$ cannot intersect on this line, similarly for $\mathcal{T}'_\alpha$ and $\bar{\mathcal{T}}'_{\bar{\alpha}}$. Hence $\varphi(\alpha) > \varphi(\bar{\alpha})$, which proves the uniqueness.

As a consequence, for $\alpha < \alpha_c$, $\varphi(\alpha) > 0$, in the plane $(y, Y)$, $\mathcal{T}_\varepsilon$ does not stay in $\mathcal{Q}_4$; for $\alpha > \alpha_c$, $\varphi(\alpha) < 0$, $\mathcal{T}_\alpha$ does not stay in $\mathcal{Q}_4$. From Lemma 9.7, it follows that $\alpha_p < \alpha_c$ if $N \geqq 2$. Moreover $\alpha^* < \alpha_c$. Indeed $\alpha^*$ is a weak source from Proposition 2.5, thus for $\alpha > \alpha^*$ small enough, there exists a unique cycle $\mathcal{O}$ around $M_\ell$, which is unstable. For such an $\alpha$, $\mathcal{T}_\varepsilon$ cannot stay in $\mathcal{Q}_4$ : it would have $\mathcal{O}$ as a limit cycle at $\infty$, which contradicts the unstability. ∎

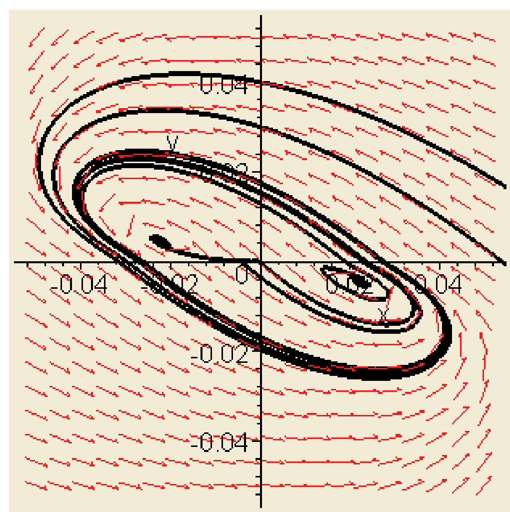
Next we discuss according to the position of $\alpha$ with respect to $\alpha^*$ and $\alpha_c$.

**Theorem 9.10** *Assume $\varepsilon = -1$, and $\alpha \leqq \alpha^*$. Then*

*(i) there exist a unique flat positive solution $w$ of ($\boldsymbol{E}_w$) with (4.3) near 0, and (4.4) near $\infty$;*
*(ii) All the other solutions are oscillating at $\infty$, among them the regular ones, and $r^{-\gamma}w$ is asymptotically periodic in $\ln r$. There exist solutions with a hole, also with (4.3), (4.6) or (4.9) or (4.9) or (4.7) near 0. There exist solutions such that $r^{-\gamma}w$ is periodic in $\ln r$.*




th 9.10,fig 12: $\varepsilon = -1, N = 1, p = 3, \alpha = -2.53$




th 9.10, fig 13: $\varepsilon = -1, N = 1, p = 3, \alpha = -2.2$

**Proof.** Here $\alpha < \alpha_c$, from Theorem 9.9, and the trajectory $\mathcal{T}_\alpha$ stays in $\mathcal{Q}_4$. From Proposition 4.4, it converges at $-\infty$ to $M_\ell$, and $w$ is of type (i).

The trajectory $\mathcal{T}_\varepsilon$ leaves $\mathcal{Q}_4$, and cannot converge either to $(0,0)$ since $\mathcal{T}_\varepsilon \neq \mathcal{T}_\alpha$, or to $\pm M_\ell$, because this point is a source, or a weak source. Recall that $M_\ell$ is a node point for $\alpha \leqq \alpha_1$ (see figure 12,, where $\alpha_1 \cong -2.50$), or a spiral point (see figure 13). And $\mathcal{T}_\varepsilon$ is bounded at $\infty$ from Proposition 4.3. Then it has a limit cycle $\mathcal{O}_\varepsilon$ surrounding $(0,0)$ from Proposition 4.4, and $\pm M_\ell$ from Remark 9.3. Thus $w$ is oscillating around 0 near $\infty$, $r^{-\gamma}w$ is asymptotically periodic in $\ln r$.

The solutions $w$ corresponding to $\mathcal{O}_\varepsilon$ are oscillating and $r^{-\gamma}w$ is periodic in $\ln r$. Any trajectory $\mathcal{T}_{[P]}$ with $P$ in the interior domain delimitated by $\mathcal{O}_\varepsilon$ converges to $M_\ell$ at $-\infty$ and has the same limit cycle at $\infty$. The trajectory $\mathcal{T}_r$ starts in $\mathcal{Q}_1$, with $\lim_{\tau\to-\infty} y = \infty$ and cannot converge to any stationary point at $\infty$. It is bounded, thus has a limit cycle $\mathcal{O}_r$ surrounding $\mathcal{O}_0$. For any $P \notin \mathcal{T}_r$ in the exterior domain to $\mathcal{O}_r$, the trajectory $\mathcal{T}_{[P]}$ admits $\mathcal{O}_r$ as a limit cycle at $\infty$, and $y$ is necessarily monotone at $-\infty$, thus (4.6) or (4.9) or (4.9) or (4.7) near 0; all those solutions exist. The question of the uniqueness of the cycle ($\mathcal{O}_r = \mathcal{O}_\varepsilon$) is open.

**Theorem 9.11** *Let $\alpha_c$ be defined by Theorem 9.9.*

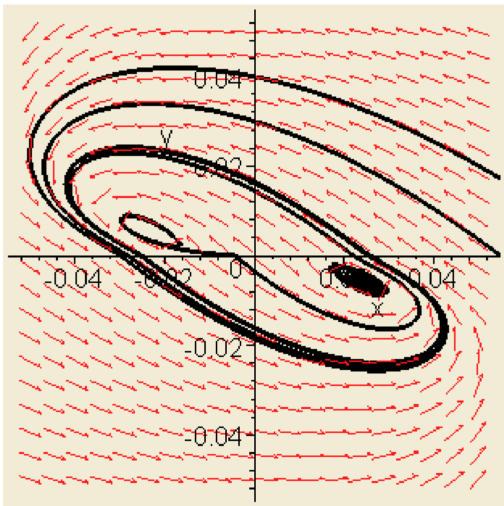
*(1) Let $\alpha^* < \alpha < \alpha_c$. Then all regular solutions $w$ of ($\boldsymbol{E}_w$) are oscillating around* 0 *near* $\infty$, *and $r^{-\gamma}w$ is asymptotically periodic in* $\ln r$. *There exist*

*(i)* ***positive*** *solutions, such that $r^{-\gamma}w$ is periodic in* $\ln r$;
*(ii) a unique positive solution such that $r^{-\gamma}w$ is asymptotically periodic in* $\ln r$ *near* 0, *with (4.4) near* $\infty$;
*(iii) positive solutions such that $r^{-\gamma}w$ is asymptotically periodic in* $\ln r$ *near* 0, *with (4.3) near* $\infty$;
*(iv) solutions oscillating around* 0 *such that $r^{-\gamma}w$ is periodic in* $\ln r$;
*(v) solutions with a hole, oscillating near* $\infty$, *such that $r^{-\gamma}w$ is asymptotically periodic in* $\ln r$;
*(vi) solutions satisfying (4.6) or (4.9) or (4.9) or (4.7) near* 0, *oscillating around* 0 *near* $\infty$, *such that $r^{-\gamma}w$ is asymptotically periodic in* $\ln r$;
*(vii) solutions positive near* 0, *oscillating near* $\infty$, *such that $r^{-\gamma}w$ is asymptotically periodic in* $\ln r$ *near* 0 *and* $\infty$.

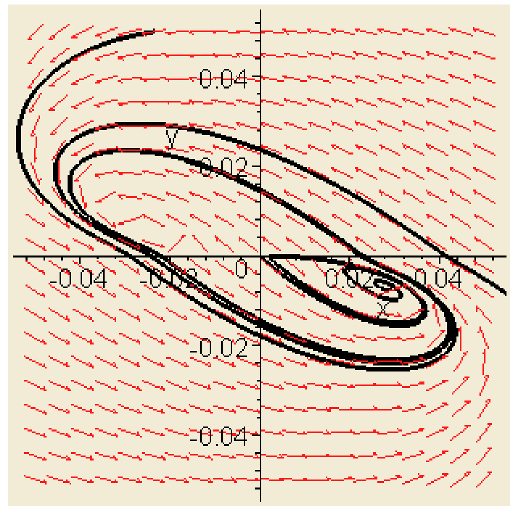
*(2) Let $\alpha = \alpha_c$.*

*(viii) There exist a* ***unique nonnegative*** *solution* ***with a hole****, with (4.4) near* $\infty$.
*The regular solutions are as above. There exist solutions of types (iv), (vi), and*
*(ix) positive solutions such that $r^{-\gamma}w$ is bounded from above near* 0, *with (4.3) near* $\infty$.

■




th 9.11,fig 14: $\varepsilon = -1, N = 1, p = 3, \alpha = -2.1$




th 9.11,fig 15: $\varepsilon = -1, N = 1, p = 3, \alpha = -2$

**Proof.** (1) Let $\alpha^* < \alpha < \alpha_c$ (see figure 14). Then $\mathcal{T}_\alpha$ stays in $\mathcal{Q}_4$, but cannot converge neither to $M_\ell$ which is a sink, nor to $(0,0)$ since $\mathcal{T}_\alpha \neq \mathcal{T}_\varepsilon$. It has a limit cycle $\mathcal{O}_\alpha$ in $\mathcal{Q}_4$ at $-\infty$, surrounding $M_\ell$, and $w$ is of type (ii). The orbit $\mathcal{O}_\alpha$ corresponds to solutions of type (i). There exist positive solutions converging to $M_\ell$ at $\infty$, with a limit cycle $\mathcal{O}_\ell$ at $-\infty$ surrounded by $\mathcal{O}_\alpha$, and $w$ is of type (iii). This cycle is unique ($\mathcal{O}_\ell = \mathcal{O}_\alpha$) for $\alpha - \alpha^*$ small enough, from Proposition 2.5. The trajectory $\mathcal{T}_\varepsilon$ still cannot stay in $\mathcal{Q}_4$. As in the case $\alpha \leqq \alpha^*$, $\mathcal{T}_\varepsilon$ has a limit cycle $\mathcal{O}_\varepsilon$ surrounding the three stationary points, $w$ is of type (v), and $\mathcal{T}_r$ is oscillating around 0, and there exist solutions of type (vi). Any trajectory $\mathcal{T}_{[P]}$ with $P \notin \mathcal{T}_\varepsilon$ in $\mathcal{Q}_4$ in the domain delimitated by $\mathcal{O}_\alpha$ and $\mathcal{O}_\varepsilon$ admits $\mathcal{O}_\alpha$ as a limit cycle at $-\infty$ and $\mathcal{O}_\varepsilon$ at $\infty$, and $w$ is of type (vii).

(2) Let $\alpha = \alpha_c$ (see figure 15). The homoclinic trajectory $\mathcal{T}_\varepsilon = \mathcal{T}_\alpha$ corresponds to the solution $w$ of type (viii). The trajectory $\mathcal{T}_r$ has a limit cycle $\mathcal{O}_r$ surrounding the three points. Thus there exist solutions of types (iv) or (vi). Any trajectory ending up at $M_\ell$ at $\infty$ is bounded, contained in the domain delimitated by $\mathcal{T}_\varepsilon$, and its limit set at $-\infty$ is the homoclinic trajectory $\mathcal{T}_\varepsilon$, or a cycle around $M_\ell$, and $w$ is of type (ix).

**Theorem 9.12** *Assume* $\varepsilon = -1$, *and* $\alpha_c < \alpha < -p'$.

*There exist a unique nonnegative solution* $w$ *of* $(\boldsymbol{E}_w)$ *with a hole, with* $r^{-\gamma} w$ *bounded from above and below at* $\infty$. *The regular solutions have at least two zeros.*

*(1) Either there exist oscillating solutions such that* $r^{-\gamma} w$ *is periodic in* $\ln r$. *Then the regular solutions have an infinity of zeros, and* $r^{-\gamma} w$ *is asymptotically periodic in* $\ln r$. *There exist*

*(i) solutions satisfying (4.6) or (4.9) or (4.9) or (4.7) near* $0$*, oscillating near* $\infty$*, such that* $r^{-\gamma}w$ *is asymptotically periodic in* $\ln r$*;*
*(ii) a unique solution oscillating near 0, such that* $r^{-\gamma}w$ *is asymptotically periodic in* $\ln r$*, and with (4.4) near* $\infty$*;*
*(iii) solutions positive near* $0$*, with* $r^{-\gamma}w$ *bounded, and oscillating near* $\infty$*, such that* $r^{-\gamma}w$ *is asymptotically periodic in* $\ln r$*.*

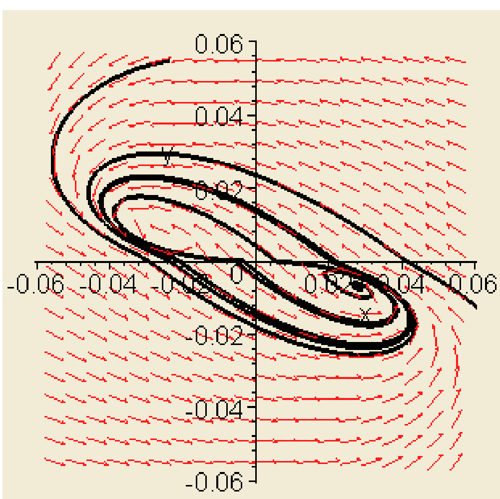
*(2) Or all the solutions have a finite number of zeros, and at least two. Two cases may occur:*

- *Either regular solutions have* $m$ *zeros and* $r^{-\gamma}w$ *bounded from above and below at* $\infty$*. Then there exist*

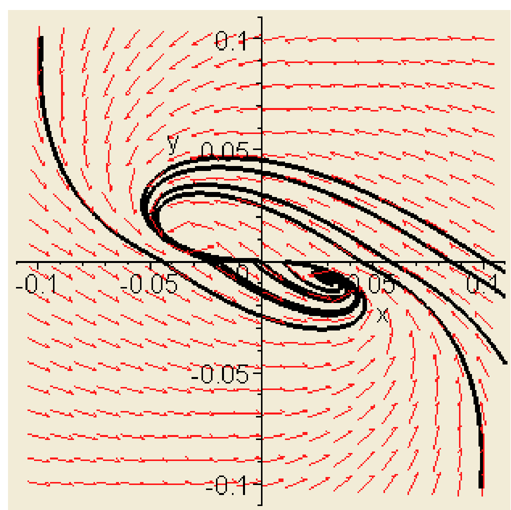
*(iv) solutions with* $m$ *zeros, with (4.6) or (4.9), with (4.4) near* $\infty$*;*
*(v) solutions with* $m$ *zeros with (4.6) or (4.9) and* $r^{-\gamma}w$ *bounded from above and below at* $\infty$*;*
*(vi) solutions with* $m+1$ *zeros with (4.6) or (4.9) and* $r^{-\gamma}w$ *bounded from above and below at* $\infty$*;*
*(vii) (for* $p > N$*) a unique solution with* $m$ *zeros,with (4.7) or (4.10) and* $r^{-\gamma}w$ *bounded from above and below at* $\infty$*.*

- *Or regular solutions have* $m$ *zeros and (4.4) holds near* $\infty$*. Then there exist solutions of type (vi) or (vii).*

■




th 9.12,fig 16: $\varepsilon = -1, N = 1, p = 3, \alpha = -1.98$




th 9.12, fig 17: $\varepsilon = -1, N = 1, p = 3, \alpha = -1.90$

**Proof.** Here $\mathcal{T}_\varepsilon$ stays in $\mathcal{Q}_4$, converges to $M_\ell$ or has a limit cycle around $M_\ell$, thus $w$ has a hole and $r^{-\gamma}w$ bounded from above and below at $\infty$. If $\alpha \geqq \alpha_2$, there is no cycle in $\mathcal{Q}_4$, from Proposition 4.4, thus $\mathcal{T}_\varepsilon$ converges to $M_\ell$.

(1) Either there exists a cycle surrounding $(0, 0)$ and $\pm M_\ell$, thus solutions $w$ oscillating around 0, such that $r^{-\gamma}w$ is periodic in $\ln r$. Then $\mathcal{T}_r$ has such a limit cycle $\mathcal{O}_r$, and $w$ is oscillating around 0. The trajectory $\mathcal{T}_\alpha$ has a limit cycle at $-\infty$ of the same type $\mathcal{O}_\alpha \subset \mathcal{O}_r$, and $w$ is of type (ii). For any $P \notin \mathcal{T}_\varepsilon$ in the interior domain in $\mathcal{O}_\alpha$, $\mathcal{T}_{[P]}$ admits $\mathcal{O}_\alpha$ as a limit cycle at $-\infty$ and converges to $M_\ell$ at $\infty$, or has a limit cycle in $\mathcal{Q}_4$; and $w$ is of type (iii). For any $P \notin \mathcal{T}_r$, in the domain exterior to $\mathcal{O}_r, \mathcal{T}_{[P]}$ has $\mathcal{O}_\alpha$ as limit cycle at $\infty$, and w is of type (i).

(2) Or no such cycle exists. Then any trajectory converges at $\infty$, any trajectory, apart from $\pm\mathcal{T}_\alpha$, converges to $\pm M_\ell$ or has a limit cycle in $\mathcal{Q}_1$. All the trajectories end up in $\mathcal{Q}_2$ or $\mathcal{Q}_4$. Since $\mathcal{T}_r$ starts in $\mathcal{Q}_1$, $y$ has at least one zero. Suppose that it is unique. Then $\mathcal{T}_r$ converges to $-M_\ell$, thus $Y$ stays positive. Consider the function $Y_\alpha = e^{(\alpha+\gamma)(p-1)\tau}Y$ defined by (2.3) with $d = \alpha$. From Theorem 3.3, $Y_\alpha = (a\,|\alpha|\,/N)e^{(\alpha(p-1)+p)\tau}(1+o(1))$ near $-\infty$; thus $Y_\alpha$ tends to $\infty$, since $\alpha < p'$. And $Y_\alpha = (\gamma\ell)^{p-1}e^{(\alpha+\gamma)(p-1)\tau}$ near $\infty$, thus also $Y_\alpha$ tends to $\infty$; then it has a minimum point $\tau$, and from (2.6), $Y_\alpha''(\tau) = (p-1)^2(\eta-\alpha)(p'+\alpha)Y_\alpha < 0$, which is contradictory. Thus $y$ has a number $m \geqq 2$ of zeros.

Either $\mathcal{T}_r \neq \mathcal{T}_\alpha$. Since the slope of $\mathcal{T}_\alpha$ near $-\infty$ is infinite and the slope of $\mathcal{T}_r$ is finite, $\mathcal{T}_\alpha$ cuts the line $\{y = 0\}$ at $m$ points, starts from $\mathcal{Q}_1$, and $w$ is of type (iv). For any $P$ in the domain of $\mathcal{Q}_1$ between $\mathcal{T}_r$ and $\mathcal{T}_\alpha$, $\mathcal{T}_{[P]}$ cuts $\{y = 0\}$ at $m+1$ points, and $w$ is of type (v). For any $P$ in the domain of $\mathcal{Q}_1$ above $\mathcal{T}_r$, $\mathcal{T}_{[P]}$ cuts the line $\{y = 0\}$ at $m+1$ points, and $w$ is of type (vi). If $p > N$, the trajectories $\mathcal{T}_-$ and $\mathcal{T}_u$ cut the line $\{y = 0\}$ at $m$ points, and $w$ is of type (vii).

Or $\mathcal{T}_r = \mathcal{T}_\alpha$, and then we find only trajectories with $w$ of type (vi) or (vii). ∎

**Remark 9.13** *Consider the regular solutions in the range $\alpha_c < \alpha < -p'$. We conjecture that there exists a decreasing sequence $(\bar{\alpha}_n)$, with $\bar{\alpha}_0 = -p'$ and $\alpha_c < \bar{\alpha}_n$ such that for $\alpha \in (\bar{\alpha}_m, \bar{\alpha}_{m-1})$, $y$ has $m$ zeros and converges to $\pm M_\ell$; and for $\alpha = \bar{\alpha}_m$, $y$ has $m+1$ zeros and converges to $(0, 0)$, thus $\mathcal{T}_r = \mathcal{T}_\alpha$. We presume that $(\bar{\alpha}_m)$ has a limit $\bar{\alpha} > \alpha_c$. And for $\alpha < \bar{\alpha}$, $y$ has an infinity of zeros, in other words there exists a cycle $\mathcal{O}_r$ surrounding $\{0\}$ and $\pm M_\ell$.*

*Numerically, for $\alpha = \alpha_c$, the cycle $\mathcal{O}_r$ seems to be the unique cycle surrounding the three points. But for $\alpha > \alpha_c$ and $\alpha - \alpha_c$ small enough, there exist **two different cycles** $\mathcal{O}_\alpha \subset \mathcal{O}_r$ (see figure 15). As $\alpha$ increases, we observe the coalescence of those cycles; they disappear after some value $\bar{\alpha}$ (see figure 16).*

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