# Initial Trace of Solutions of Some Quasilinear Parabolic Equations with Absorption 

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We study the existence of an initial trace of nonnegative solutions of the problem

$$
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+u^{q}=0 \quad \text { in } Q_{T}=\Omega \times(0, T) .
$$

We prove that the initial trace is an outer regular Borel measure which may not be locally bounded for some values of the parameters $p$ and $q$. We study also the corresponding Cauchy problems with a given generalized Borel measure as initial data. © 2002 Elsevier Science (USA)

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## 1. INTRODUCTION

Let $\Omega$ be a domain in $\mathbb{R}^{N}(N \geqslant 1)$, possibly unbounded. The aim of this article is to investigate the initial trace problem for the following class of quasilinear equations with absorption:

$$
\begin{equation*}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+u^{q}=0 \quad \text { in } Q_{T}=\Omega \times(0, T), \tag{1.1}
\end{equation*}
$$

in the range $q>0, p>1$. We prove the existence of an initial trace in the class $\mathscr{B}_{+}^{\text {reg }}(\Omega)$ of outer regular positive Borel measures in $\Omega$, not necessarily locally bounded. Conversely, given a measure $v \in \mathscr{B}_{+}^{\mathrm{reg}}(\Omega)$, we study the initial-boundary value problem for Eq. (1.3) with initial data $v$.

The initial trace problem for a nonnegative solution $u$ of the mere heat equation ( $p=2$, no absorption) in $Q_{T}$ is classical. It is easy to establish that there exists a unique positive Radon measure $\mu$ in $\Omega$ which is the initial trace of $u$, in the sense that

$$
\lim _{t \rightarrow 0} \int_{\Omega} u(x, t) \xi(x) d x=\int_{\Omega} \xi(x) d \mu(x), \quad \forall \xi \in C_{\mathrm{c}}(\Omega)
$$

Moreover, when $\Omega=\mathbb{R}^{N}$, the initial trace $\mu$ satisfies the following growth condition:

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} e^{-|x|^{2} / 4 T} d \mu(x)<\infty \tag{1.2}
\end{equation*}
$$

Conversely, if $\mu$ is a nonnegative Radon measure in $\mathbb{R}^{N}$ satisfying (1.2), there exists a unique solution of the heat equation in $\mathbb{R}^{N} \times(0, T)$ with initial trace $\mu$.

Concerning the semilinear heat equation with absorption

$$
\begin{equation*}
\partial_{t} u-\Delta u+u^{q}=0 \tag{1.3}
\end{equation*}
$$

a rather complete picture of the initial trace problem is provided by Marcus and Véron [33] who prove that the initial trace is well defined in the class $\mathscr{B}_{+}^{\mathrm{reg}}(\Omega)$. Their first result asserts that there exists a relatively closed subset $\mathscr{S} \subseteq \Omega$ and a nonnegative Radon measure $\mu$ on $\mathscr{R}=\Omega \backslash \mathscr{S}$ such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathscr{R}} u(x, t) \xi(x) d x=\int_{\mathscr{R}} \xi(x) d \mu(x), \quad \forall \xi \in C_{\mathrm{c}}(\mathscr{R}) \tag{1.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{V} u(x, t) d x=\infty, \quad \forall V \text { open s.t. } \mathscr{S} \cap V \neq \emptyset \tag{1.5}
\end{equation*}
$$

Moreover, $\mathscr{S}$ and $\mu$ are uniquely determined. They define in a unique way an outer regular Borel measure $v$ which is the initial trace of $u$ :

$$
\operatorname{tr}_{\Omega}(u)=(\mathscr{S}, \mu) \approx v
$$

The reverse problem is to reconstruct the solution from a given outer regular Borel measure $v=(\mathscr{S}, \mu) \in \mathscr{B}_{+}^{\text {reg }}(\Omega)$. The role of the critical exponent

$$
\begin{equation*}
q_{*}=1+2 / N \tag{1.6}
\end{equation*}
$$

is pointed out by the fact that if $1<q<q_{*}$ there exist no removable singularity for Eq. (1.3), in the sense that there exist nontrivial solutions of (1.3) continuous in $\bar{Q} \backslash\{(a, 0)\} \quad$ (where $a \in \Omega$ ) and vanishing on $\bar{\Omega} \backslash\{a\} \times\{0\} \cup \partial \Omega \times[0, T)$. From this result follows that, always in the subcritical case $1<q<q_{*}$, any outer regular Borel measure $v$ is eligible for being the initial trace of a nonnegative solution of (1.3). Moreover, under minor additional assumptions, uniqueness of the solution of the generalized Cauchy-Dirichlet problem

$$
\left\{\begin{array}{l}
\partial_{t} u-\Delta u+u^{q}=0 \quad \text { in } Q_{T}  \tag{1.7}\\
u=f \in L_{+}^{1}(\partial \Omega \times(0, T)) \\
\operatorname{tr}_{\Omega}(u)=v
\end{array}\right.
$$

is obtained. On the opposite, if $q \geqslant q_{*}$ the Cauchy-Dirichlet problem (1.7) becomes much more difficult: necessary and sufficient conditions on the concentration of the Radon part $\mu$ and the singular part $\mathscr{S}$ of $v$ are expressed in terms of Bessel capacities $C_{2 / q, q^{\prime}}$. When these conditions are fulfilled existence of a maximal solution follows, but usually uniqueness does not hold.

When the heat equation is replaced by the porous medium equation

$$
\begin{equation*}
\partial_{t} u-\Delta u^{m}=0 \quad \text { in } \mathbb{R}^{N} \times(0, T) \tag{1.8}
\end{equation*}
$$

the existence of an initial trace $\mu$ in the set of positive Radon measures in $\mathbb{R}^{N}$ is obtained by Aronson and Caffarelli [2]. They also prove that $\mu$ has to satisfy some polynomial growth at infinity. The corresponding initial value problem is studied in [10]. The trace question for the equation with absorption

$$
\begin{equation*}
\partial_{t} u-\Delta u^{m}+u^{q}=0 \quad \text { in } Q_{T} \tag{1.9}
\end{equation*}
$$

is considered by the author in $[20,21]$ in the range $1<m \leqslant q$. If $u$ is a nonnegative solution of (1.9) there exists an outer regular Borel measure $v=(\mathscr{S}, \mu)$ such that properties (1.4) and (1.5) hold in the same way as above. This defines the initial trace $\operatorname{tr}_{\Omega}(u)$. There are two critical values

$$
\begin{equation*}
q=m \quad \text { and } \quad q=q_{m}=m+2 / N \tag{1.10}
\end{equation*}
$$

If $q=m$ the diffusion is comparable to the absorption, which is superlinear. Therefore there are only two possibilities:
(i) either $\mathscr{S}=\emptyset$ and $v$ is a reduced to its regular part which is a Radon measure with some growth at infinity,
(ii) or $\mathscr{S}=\Omega$ and $u$ is the flat solution, that is

$$
u(x, t)=\left(\frac{1}{t(m-1)}\right)^{1 /(m-1)}
$$

If $m<q$ the absorption dominates the diffusion but the situation differs according to $m<q<q_{m}$ or $q \geqslant q_{m}$. If $m<q<q_{m}$ any outer regular Borel measure $v$ is eligible for being the initial trace of a nonnegative solution $u$ of (1.9). An existence result of a solution of the generalized Cauchy-Dirichlet problem

$$
\left\{\begin{array}{l}
\partial_{t} u-\Delta u^{m}+u^{q}=0 \quad \text { in } Q_{T}  \tag{1.11}\\
u=f \in L_{+}^{m}(\partial \Omega \times(0, T)) \\
\operatorname{tr}_{\Omega}(u)=v \in \mathscr{B}_{+}^{\mathrm{reg}}(\Omega)
\end{array}\right.
$$

also holds, but uniqueness is not proved except in the case where $v$ has no singular part and $f=0$. If $q \geqslant q_{m}$, the situation is again more difficult as a result of the fact that isolated points at $t=0$ are removable singularities for solutions of (1.9). Some sufficient conditions for the existence of a solution to (1.11) exist, however up to now they have not been proved to be necessary.

The situation for the $p$-Laplacian-type equation is in some sense more difficult than for the porous-medium-type equation. One of the reasons is that the full duality argument in a $L^{1}$ framework, which was at the core of Marcus and Véron or Chasseigne's results, does not fit with the $p$-Laplace operator. As a consequence the proof of the existence of an initial trace is much more difficult to obtain. The notion of weak solution is not at all straightforward, even in the case of the equation without absorption

$$
\begin{equation*}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)=0 \quad \text { in } Q_{T}=\Omega \times(0, T) \tag{1.12}
\end{equation*}
$$

Two critical values for $p$ exist for expressing the notion of weak solution to (1.12) with a measure or an integrable function as initial data,

$$
\begin{equation*}
p_{0}=\frac{2 N}{N+1} \quad \text { and } \quad p_{1}=\frac{2 N+1}{N+1} \tag{1.13}
\end{equation*}
$$

If $1<p \leqslant p_{1}$ the gradient is not well defined, in particular is not an integrable function. This is why the different authors who studied as well stationary or time-dependent equations involving the $p$-Laplacian were led to introduce the notion of renormalized or entropy solutions which are solutions such that $\nabla T_{k}(u) \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right)$ for any $k>0$, where

$$
T_{k}(s)=\left\{\begin{array}{l}
s \quad \text { if }|s| \leqslant k  \tag{1.14}\\
k \operatorname{sign}(s) \quad \text { if }|s|>k
\end{array}\right.
$$

Using this definition, Di Benedetto and Herrero proved in [7, 8] that any nonnegative solution of $(1.12)$ in $\mathbb{R}^{N} \times(0, T)$ admits an initial trace $\mu$ which is a nonnegative Radon measure, and the proof heavily relies on the parabolic Harnack inequality. When $p>2$, it is also derived that the measure $\mu$ has to satisfy a precise polynomial growth at infinity. Moreover, the same authors study the corresponding initial value problem with an initial data $\mu \in L_{\text {loc }}^{1}\left(\mathbb{R}^{N}\right)$ and obtain existence and uniqueness results. If $\Omega$ is bounded, the initial value problem is considered in [3,16] when $p>p_{1}$ and in [14] (see also [15]) in the general case. When $\mu$ is a Radon measure, new difficulties appear in the range $1<p \leqslant p_{0}$. The a priori estimates show that $u$ may not be an integrable function, and some questions dealing with the concentration of the measure are unavoidable.

The same kind of difficulties are present in Eq. (1.1). Since it is only assumed that $q$ is positive, three critical values for $q$ appear

$$
q=1, \quad q=p-1, \quad q=q_{\mathrm{c}}=p-1+\frac{p}{N}
$$

The case $q>p-1$ is the analogous of the superlinearity case in the case $p=2$. It means that the absorption term is dominant with respect to the diffusion operator. The exponent $q=q_{\mathrm{c}}$ is the mere extension of the critical values $q_{*}$ and $q_{m}$ for Eqs. (1.3) and (1.9). When $q>1$, the absorption term is dominant with respect to the diffusion operator, and (1.1) admits a particular singular solution in $\mathbb{R}^{N} \times(0, \infty)$, called the flat solution, which is the function

$$
\begin{equation*}
(x, t) \mapsto W(x, t)=\left(\frac{1}{t(q-1)}\right)^{1 /(q-1)} \tag{1.15}
\end{equation*}
$$

This particular solution plays an important role since it dominates any nonnegative solution of $(1.1)$ which is locally bounded in $(0, \infty) \times \mathbb{R}^{N}$.

### 1.1. Existence Results for the Initial Trace

Our main result which states the basis of the definition of the initial trace of nonnegative solution of (1.1) is the following:

- Let $u$ be a nonnegative weak solution of (1.1) in $Q_{T}$. Then for any $y \in \Omega$ the following alternative occurs:
(i) either for any open subset $U \subset \Omega$ containing $y$

$$
\lim _{t \rightarrow 0} \int_{U} u(x, t) d x=\infty
$$

(ii) or there exist an open neighborhood $U \subset \Omega$ of $y$ and a Radon measure $\ell_{U}$ on $U$ such that for any $\zeta \in C_{\mathrm{c}}(U)$,

$$
\lim _{t \rightarrow 0} \int_{U} u(x, t) \zeta(x) d x=\ell_{U^{*}}(\zeta)
$$

Owing to this result it follows the definition of the initial trace of $u$ as an outer regular Borel measure $v=(\mathscr{S}, \mu) \in \mathscr{B}_{+}^{\text {reg }}(\Omega)$. Notice that this result also applies to Eq. (1.12) with $v=(\emptyset, \mu)$, and our proof does not require the Harnack inequality, which, among others, brings some simplification to the proofs in [7, 8].

The first case in which the singular set $\mathscr{S}$ is empty is the following:

- When $0<q \leqslant 1$, and $p<2$ the initial trace is reduced to a Radon measure.

If $0<q \leqslant p-1$ and $p>2$, the absorption term is dominated and the initial trace derives from the diffusion operator. However, the situation differs completely according to $0<q \leqslant 1$ or $q>1$, and $\Omega$ is bounded or $\mathbb{R}^{N}$.

- When $0<q \leqslant 1, p>2, \Omega=\mathbb{R}^{N}$, and $u$ is continuous in $Q_{T}$, the initial trace is a Radon measure.

The proof of the last result is based upon the notion of concentration of mass on a single point.

- When $0<q<1, p>2$, and $\Omega$ is bounded, then
(i) either the initial trace of $u$ is a Radon measure in $\Omega$,
(ii) or

$$
\liminf _{t \rightarrow 0} t^{1 /(p-2)} u(x, t) \geqslant v_{\Omega}(x)
$$

where $v=v_{\Omega}$ is the unique nonnegative solution of

$$
\left\{\begin{array}{l}
-\nabla \cdot\left(|\nabla v|^{p-2} \nabla v\right)=\frac{1}{p-2} v \quad \text { in } \Omega \\
v=0 \quad \text { on } \partial \Omega
\end{array}\right.
$$

Notice that in that case the function $V_{\Omega}(x, t)=t^{-1 /(p-2)} v_{\Omega}(x)$ is a particular singular solution of

$$
\left\{\begin{array}{ll}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)=0 \\
u=0 & \text { on } \partial \Omega \times(0, \infty)
\end{array} \quad \text { in } \Omega \times(0, \infty)\right.
$$

When $q=p-1$ the same dichotomy holds with $v_{\Omega}$ replaced by $w=w_{\Omega}$ unique nonnegative solution of

$$
\left\{\begin{array}{l}
-\nabla \cdot\left(|\nabla w|^{p-2} \nabla w\right)+w^{p-1}=\frac{1}{p-2} w \quad \text { in } \Omega \\
w=0 \quad \text { on } \partial \Omega
\end{array}\right.
$$

When $1<q \leqslant p-1$ the absorption term is still dominated but the existence of the particular flat singular solution $W$ creates a more complicated situation. In the case of $\mathbb{R}^{N}$ we prove the following.

- When $1<q \leqslant p-1$ and $\Omega=\mathbb{R}^{N}$, then
(i) either the initial trace of $u$ is a Radon measure in $\mathbb{R}^{N}$,
(ii) or

$$
\liminf _{t \rightarrow 0} t^{1 /(q-1)} u(x, t) \geqslant\left(\frac{1}{(q-1)}\right)^{1 /(q-1)}
$$

If $u(., t)$ is bounded for any $t>0$, assertion (ii) becomes
(ii) or

$$
u(x, t) \equiv W(x, t)
$$

When $\Omega$ is bounded, the situation is very intricated because of the existence of many different types of singular solutions, and the role of the boundary conditions, if any, may be dominant.

### 1.2. The Generalized Initial Value Problem

The reverse problem is to construct a solution $u$ of (1.1) with a given initial trace $v=(\mathscr{S}, \mu)$.

We begin with the case where the initial trace is a Radon measure, i.e. $\mathscr{S}=\emptyset$. Similarly to the case $p=2, q_{\mathrm{c}}$ is a critical exponent. If $q \geqslant q_{\mathrm{c}}$ and $q>1$, Eq. (1.1) admits no solution with isolated singularities, see [27] when $p>2$, and [23] when $p<2$. If $1<q<q_{\mathrm{c}}$, for any $k>0$, there exists a unique
solution $\tilde{w}_{k}$ of problem

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{w}_{k}-\nabla \cdot\left(\left|\nabla \tilde{w}_{k}\right|^{p-2} \nabla \tilde{w}_{k}\right)+\tilde{w}_{k}^{q}=0 \quad \text { in } \mathbb{R}^{N} \times(0, \infty),  \tag{1.16}\\
\tilde{w}_{k}(., 0)=k \delta_{0} \quad \text { in } \mathbb{R}^{N}
\end{array}\right.
$$

from [32] when $p>2$ and [23] when $p<2$. Our existence result is the following:

- Assume $q<q_{\mathrm{c}}, \quad p>p_{0}$, and $p-1<q$, or $p<2$. For any nonnegative Radon measure $\mu$ on $\mathbb{R}^{N}$, there exists a weak solution $u$ of (1.1) in $\mathbb{R}^{N} \times$ $(0, \infty)$ with initial trace $\mu$.

Notice that no growth condition is assumed, since $p-1<q$, or $p<2$. In the case where $\Omega$ is bounded, we also extend the results of [3] from $p>p_{1}$ to $p>p_{0}$.

Next, we come to the general problem of existence of a solution for a given outer regular nonnegative Borel measure $v=(\mathscr{S}, \mu)$ as initial trace. Since in the case $q \leqslant p-1$, where the diffusion term is dominant, either the solutions are almost explicit and $\mathscr{S}=\mathbb{R}^{N}$, or the initial trace is a Radon measure and $\mathscr{S}=\emptyset$, we concentrate on the case

$$
\begin{equation*}
q>\max (1, p-1) \tag{1.17}
\end{equation*}
$$

where the absorption is dominant. In the subcritical case $q<q_{\mathrm{c}}$, we construct solutions by approximating $v$ by a sequence of Radon measures $\left\{\mu_{\kappa}\right\}$. Since $q>1$, the corresponding sequence of solutions $\left\{u_{k}\right\}$ is dominated by the flat solution $W$. Hence $\left\{u_{k}\right\}$ remains locally uniformly bounded and the truncation plays no role. Moreover, since $q>p-1$, no growth condition is needed. Nevertheless, the problem is still complicated by the possible nonuniqueness of the $u_{\kappa}$, and we show that we can construct a nondecreasing sequence of such solutions, that we call constructive. We get the following.

- When $\max (1, p-1)<q<q_{\mathrm{c}}$, for any given nonnegative Borel measure $v$ in $\mathbb{R}^{N}$ there exists a nonnegative solution $u$ of (1.1) in $\mathbb{R}^{N} \times$ $(0, \infty)$ with initial trace $v$.

In the supercritical case $q \geqslant q_{\mathrm{c}}$, the situation is much more delicate and we give a sufficient condition on $v$ to be an initial trace. This condition is the natural extension of the one discovered by Marcus and Véron [33] in the case $p=2$. In their case this condition is expressed in terms of Bessel capacity.

## 2. PRELIMINARIES

In this section, we assume that $\Omega$ is an any domain in $\mathbb{R}^{N}, T>0$, and set

$$
Q_{T}=\Omega \times(0, T)
$$

We denote by $\mathscr{M}^{+}(\Omega)$ the set of nonnegative Radon measures on $\Omega$. In the sequel, for any open set $U$, we write $U \subset \subset \Omega$ whenever $U$ has a compact closure in $\Omega$.

We consider a more general equation than (1.1), namely

$$
\begin{equation*}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+a(x) u^{q}=0 \tag{2.1}
\end{equation*}
$$

where $a \in L_{\mathrm{loc}}^{\infty}(\Omega)$.
Definition 2.1. A nonnegative function $u$ is said to be a weak solution of (2.1) in $Q_{T}$, if

$$
u \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right), \quad a u^{q} \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right), \quad \nabla u \in L_{\mathrm{loc}}^{p}\left(Q_{T}\right)
$$

and

$$
\begin{equation*}
\int_{0}^{T} \int_{\Omega}\left(-H(u) \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u \cdot \nabla(h(u) \varphi)+h(u) u^{q} \varphi\right) d x d t=0 \tag{2.2}
\end{equation*}
$$

for any $\varphi \in C_{\mathrm{c}}^{\infty}\left(Q_{T}\right)$ and any function $h \in C(\mathbb{R}) \cap W^{1, \infty}(\mathbb{R})$ where $H^{\prime}(r)=h(r)$.

It follows from the definition that for any $\eta \in C_{\mathrm{c}}(\Omega)$,

$$
t \mapsto \int_{\Omega} H(u(x, t)) \eta(x) d x
$$

is continuous and there holds

$$
\begin{align*}
& \int_{t}^{\theta} \int_{\Omega}\left(-H(u) \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u \cdot \nabla(h(u) \varphi)+h(u) a u^{q} \varphi\right) d x d t \\
& \quad=\int_{\Omega} H(u(x, t)) \varphi(x, t) d x-\int_{\Omega} H(u(x, \theta)) \varphi(x, \theta) d x \tag{2.3}
\end{align*}
$$

for any $0<t<\theta<T$ and $\varphi \in C_{\mathrm{c}}^{\infty}(\Omega \times[0, T])$.

As a consequence of (2.3) there holds

$$
\begin{align*}
& \int_{t}^{\theta} \int_{\Omega}\left(h(u)|\nabla u|^{p-2} \nabla u \cdot \nabla \xi+h^{\prime}(u)|\nabla u|^{p} \xi+h(u) a u^{q} \xi\right) d x d t \\
& \quad=\int_{\Omega} H(u(x, t)) \xi(x) d x-\int_{\Omega} H(u(x, \theta)) \xi(x) d x \tag{2.4}
\end{align*}
$$

for any $0<t<\theta<T$ and $\xi \in C_{\mathrm{c}}^{\infty}(\Omega)$.
Remark 2.1. In the definition of the weak solution, we can impose to the function $h$ to be constant outside $[-k, k]$ for some $k>0$. Therefore the term $|\nabla u|^{p}$ in identity (2.2) is only considered on the set

$$
\left\{(x, t) \in Q_{T}:|u(x, t)| \leqslant k\right\} .
$$

The assumption $\nabla u \in L_{\mathrm{loc}}^{p}\left(Q_{T}\right)$ can be replaced by

$$
|\nabla u|^{p-1} \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right), \quad \nabla T_{k}(u) \in L_{\mathrm{loc}}^{p}\left(Q_{T}\right),
$$

where $T_{k}$ is defined in (1.14), and the gradient of $u$, denoted by $y=\nabla u$, is defined by

$$
\nabla\left(T_{k}(u)\right)=y \times 1_{\{|u| \leqslant k\}} \quad \text { a.e. in } Q_{T} .
$$

We will again say that $u$ is a weak solution. This class corresponds to the classes of entropy or renormalized solutions defined in [1, 44], and is closely related to the ones of $[8,14,15,39,40]$, see Remark 2.4.

### 2.1. Integral Estimates

We give below some integral estimates valid for the solutions of Eq. (1.1). The method employed here is formally settled upon multiplying successively the equation by $(1+u)^{\alpha} \eta$ with $\alpha<0$ and $\eta \in C_{\mathrm{c}}^{\infty}(\Omega)$ and by $\eta$. This technique has been used in [8] for proving the Harnack inequality in the case of the mere diffusion equation

$$
\begin{equation*}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)=0 \tag{2.5}
\end{equation*}
$$

in the case $1<p<2$. It has also been employed by Bidaut-Véron and Pohozaev [13] and Mitidieri and Pohozaev [37] for getting estimates and nonexistence results of solutions of the stationary equation

$$
\begin{equation*}
-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+u^{q}=0 \tag{2.6}
\end{equation*}
$$

Proposition 2.1. Let $\alpha<0, \alpha \neq-1$, and $0<t<\theta<T$. Let $u$ be a nonnegative weak solution of (2.1) in $Q_{T}$. For any nonnegative function $\zeta \in C_{c}^{\infty}(\Omega)$,
and any $\tau>p$,

$$
\begin{align*}
& \frac{1}{\alpha+1} \int_{\Omega}(1+u(x, t))^{\alpha+1} \zeta^{\tau}(x) d x+\frac{|\alpha|}{2} \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p \zeta^{\tau} d x d t} \\
& \quad \leqslant \frac{1}{\alpha+1} \int_{\Omega}(1+u(x, \theta))^{\alpha+1} \zeta^{\tau}(x) d x+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{q+\alpha} \zeta^{\tau} d x d t \\
& \quad+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha+p-1} \zeta^{\tau-p}|\nabla \zeta|^{p} d x d t \tag{2.7}
\end{align*}
$$

and

$$
\begin{align*}
\int_{\Omega}(1+u(x, t)) \zeta^{\tau}(x) d x \leqslant & \int_{\Omega}(1+u(x, \theta)) \zeta^{\tau}(x) d x+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{q} \zeta^{\tau} d x d t \\
& +C \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t \\
& +C \int_{t}^{\theta} \int_{\Omega}|\nabla \zeta|^{p}(1+u)^{(1-\alpha)(p-1)} \zeta^{\tau-p} d x d t \tag{2.8}
\end{align*}
$$

where $C=C\left(\alpha, p, q, \tau,\|a\|_{L^{\infty}(\Omega)}\right)$. Conversely, if $a(x) \geqslant 0$ a.e. in $\Omega$, then

$$
\begin{align*}
& \frac{1}{4} \int_{\Omega}(1+u(x, \theta)) \zeta^{\tau}(x) d x+\frac{1}{2} \int_{t}^{\theta} \int_{\Omega} a u^{q \zeta^{\tau}} d x d t \\
& \leqslant \int_{\Omega}(1+u(x, t)) \zeta^{\tau}(x) d x+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t \\
& \quad+C \int_{t}^{\theta} \int_{\Omega}|\nabla \zeta|^{p}(1+u)^{(1-\alpha)(p-1)} \zeta^{\tau-p} d x d t, \tag{2.9}
\end{align*}
$$

where $C=C(\alpha, p, q, \tau)$.
Proof. We set for any $\alpha \leqslant 0, \alpha \neq-1$,

$$
g_{\alpha}(u)=(1+u)^{\alpha}, \quad G_{\alpha}(u)=\frac{(1+u)^{1+\alpha}}{1+\alpha}
$$

and take

$$
\begin{equation*}
\eta=\zeta^{\tau} \quad \text { and } \quad H(u)=G_{\alpha}(u) \tag{2.10}
\end{equation*}
$$

for test functions in (2.4). Then

$$
\begin{aligned}
& \int_{t}^{\theta} \int_{\Omega}\left(\tau g_{\alpha}(u) \zeta^{\tau-1}|\nabla u|^{p-2} \nabla u \cdot \nabla \zeta+g_{\alpha}^{\prime}(u) \zeta^{\tau}|\nabla u|^{p}+g_{\alpha}(u) \zeta^{\tau} a u^{q}\right) d x d t \\
& \quad=\int_{\Omega} G_{\alpha}(u(x, t)) \zeta^{\tau}(x) d x-\int_{\Omega} G_{\alpha}(u(x, \theta)) \zeta^{\tau}(x) d x
\end{aligned}
$$

or, equivalently,

$$
\begin{align*}
& \frac{1}{\alpha+1} \int_{\Omega}(1+u(x, t))^{\alpha+1} \zeta^{\tau}(x) d x+|\alpha| \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1} \zeta^{\tau}|\nabla u|^{p} d x d t \\
& \quad=\frac{1}{\alpha+1} \int_{\Omega}(1+u(x, \theta))^{\alpha+1} \zeta^{\tau}(x) d x+\int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha} \zeta^{\tau} a u^{q} d x d t \\
& \quad+\tau \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha} \zeta^{\tau-1}|\nabla u|^{p-2} \nabla u \cdot \nabla \zeta d x d t \tag{2.11}
\end{align*}
$$

Writing

$$
\begin{equation*}
|\nabla u|^{p-1} \zeta^{\tau-1}|\nabla \zeta|=|\nabla u|^{p-1}(1+u)^{(\alpha-1) / p^{\prime}}(1+u)^{(1-\alpha) / p^{\prime}} \zeta^{\tau-1}|\nabla \zeta| \tag{2.12}
\end{equation*}
$$

we derive for $\alpha<0$,

$$
\begin{align*}
& \tau \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha} \zeta^{\tau-1}|\nabla u|^{p-2} \nabla u \cdot \nabla \zeta d x d t \\
& \quad \leqslant \frac{|\alpha|}{2} \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t \\
& \quad+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha+p-1} \zeta^{\tau-p}|\nabla \zeta|^{p} d x d t \tag{2.13}
\end{align*}
$$

Hence (2.7) holds.
As a particular case of (2.11) (with $\alpha_{0}=0$ ),

$$
\begin{align*}
\int_{\Omega}(1+u(x, t)) \zeta^{\tau}(x) d x= & \int_{\Omega}(1+u(x, \theta)) \zeta^{\tau}(x) d x+\int_{t}^{\theta} \int_{\Omega} \zeta^{\tau} a u^{q} d x d t \\
& +\tau \int_{t}^{\theta} \int_{\Omega} \zeta^{\tau-1}|\nabla u|^{p-2} \nabla u \cdot \nabla \zeta d x d t \tag{2.14}
\end{align*}
$$

From (2.12) we derive, for any $\alpha<0$,

$$
\begin{align*}
\int_{t}^{\theta} & \int_{\Omega}|\nabla u|^{p-1} \zeta^{\tau-1}|\nabla \zeta| d x d t \\
\leqslant & \int_{t}^{\theta} \int_{\Omega}|\nabla u|^{p}(1+u)^{\alpha-1} \zeta^{\tau} d x d t \\
\quad & +\int_{t}^{\theta} \int_{\Omega}|\nabla \zeta|^{p}(1+u)^{(1-\alpha)(p-1)} \zeta^{\tau-p} d x d t \tag{2.15}
\end{align*}
$$

Estimate (2.8) follows from (2.14), (2.15) and (2.7).
Next, we set

$$
d_{\alpha}(u)=1-g_{\alpha}(u)=1-(1+u)^{\alpha} \quad \text { and } \quad D_{\alpha}=(1+u)-G_{\alpha} .
$$

Subtracting (2.11) from (2.14) infers

$$
\begin{align*}
\int_{\Omega} & D_{\alpha}(u(x, \theta)) \zeta^{\tau}(x) d x+|\alpha| \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p \zeta^{\tau}} d x d t \\
& +\int_{t}^{\theta} \int_{\Omega} a u^{q} d_{\alpha}(u) \zeta^{\tau} d x d t \\
= & \int_{\Omega} D_{\alpha}(u(x, t)) \zeta^{\tau}(x) d x-\tau \int_{t}^{\theta} \int_{\Omega} d_{\alpha}(u) \zeta^{\tau-1}|\nabla u|^{p-2} \nabla u \cdot \nabla \zeta d x d t \tag{2.16}
\end{align*}
$$

But $d_{\alpha}(u)>0$ and $1+u \geqslant D_{\alpha}(u)$. Moreover,

$$
D_{\alpha}(u) \geqslant u / 4 \quad \text { and } \quad d_{\alpha}(u) \geqslant 1 / 2 \quad \text { on } \quad\left\{(x, t) \in Q_{T}: u(x, t) \geqslant 2\left(2^{-1 / \alpha}-1\right)\right\}
$$

Henceforth, since $a$ is nonnegative,

$$
\begin{align*}
& \frac{1}{4} \int_{\Omega} u(x, \theta) \zeta^{\tau}(x) d x+\frac{1}{2} \int_{t}^{\theta} \int_{\Omega} a u^{q} \zeta^{\tau} d x d t \\
& \quad \leqslant C+\int_{\Omega} u(x, t) \zeta^{\tau}(x) d x+\tau \int_{t}^{\theta} \int_{\Omega} \zeta^{\tau-1}|\nabla u|^{p-1}|\nabla \zeta| d x d t \tag{2.17}
\end{align*}
$$

Thus (2.9) follows from (2.15) and(2.17).
Remark 2.2. In case $\alpha=-1$, the term $\frac{1}{\alpha+1} \int_{\Omega}(1+u(x, t))^{\alpha+1} \zeta^{\tau}(x) d x$ at time $t$ in (2.7) has to be replaced by $\int_{\Omega} \ln (1+u(x, t)) \zeta^{\tau}(x) d x$, similarly for the term at time $\theta$.

Remark 2.3. If we consider the notion of weak solutions as it is defined in Remark 2.1, we have to replace the test function $g_{\alpha}$ with $\alpha \leqslant 0$ by $g_{\alpha, k}=g_{\alpha} \circ T_{k}$. Up to this change, relations (2.11), (2.14), (2.15) and estimates
(2.7), (2.8) remain valid with $u$ replaced by $T_{k}(u)$. The term $u^{q}$ remaining unchanged.

### 2.2. Regularity Properties

Here we derive regularity properties of the solutions under assumptions of boundedness for some integrals of $u$. We assume that $\Omega$ is an any domain in $\mathbb{R}^{N}$, and recall that $Q_{T}=\Omega \times(0, T)$.

Proposition 2.2. Let $u$ be a nonnegative solution of (2.1) in $Q_{T}$, with $a \in L_{\mathrm{loc}}^{\infty}(\Omega)$. Let $0<\theta<T$. Assume that two of the three following conditions hold, for any open set $U \subset \subset \Omega$ :

$$
\begin{gather*}
\sup _{t \in(0, \theta]} \int_{U} u(x, t) d x<\infty  \tag{2.18}\\
\int_{0}^{\theta} \int_{U}\left(|a| u^{q}+u^{p-1}\right) d x d t<\infty  \tag{2.19}\\
\int_{0}^{\theta} \int_{U}|\nabla u|^{p-1} d x d t<\infty \tag{2.20}
\end{gather*}
$$

Then the third one holds for any $U \subset \subset \Omega$. Moreover,

$$
\begin{equation*}
\int_{0}^{\theta} \int_{U} u^{\sigma} d x d t<\infty, \quad \forall \sigma \in\left(0, q_{\mathrm{c}}\right) \tag{2.21}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{0}^{\theta} \int_{U}|\nabla u|^{r} d x d t<\infty, \quad \forall r \in\left(0, \frac{N}{N+1} q_{\mathrm{c}}\right)=\left(0, p-\frac{N}{N+1}\right) \tag{2.22}
\end{equation*}
$$

Finally, there exists a Radon measure $\ell \in \mathscr{M}^{+}(\Omega)$ such that for any $\xi \in C_{\mathrm{c}}^{\infty}(\Omega)$,

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\Omega} u(x, t) \xi(x) d x=\ell(\xi) \tag{2.23}
\end{equation*}
$$

and $u$ satisfies

$$
\begin{align*}
& \left.\int_{0}^{\theta} \int_{\Omega}\left(-u \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u \cdot \nabla u \varphi\right)+a u^{q} \varphi\right) d x d t \\
& \quad=\int_{\Omega} \varphi(x, 0) d \ell(x)-\int_{\Omega} u(x, \theta) \varphi(x, \theta) d x \tag{2.24}
\end{align*}
$$

for any $0<\theta<T$ and $\varphi \in C_{c}^{\infty}(\Omega \times[0, T))$.

Proof. (i) Assume that (2.18) and (2.20) hold. Let $\zeta$ and $\tau$ be as in Proposition 2.1. From (2.14) it follows that $a u^{q} \in L^{1}\left((0, \theta), L_{\mathrm{loc}}^{1}(\Omega)\right)$. Now let us prove that $u^{p-1} \in L^{1}\left((0, \theta), L_{\text {loc }}^{1}(\Omega)\right)$. Let $U$ be any open subset such that $U \subset \subset \Omega$. If $p \leqslant 2$, taking $\zeta \equiv 1$ in $U$, we observe that

$$
\int_{t}^{\theta} \int_{U}(1+u)^{p-1} d x d t \leqslant \int_{t}^{\theta} \int_{\Omega}(1+u) \zeta^{\tau} d x d t \leqslant C
$$

If $p>2$, we use Poincaré's inequality,

$$
\int_{U}|u(x, t)-\bar{u}(t)|^{p-1} d x \leqslant C \int_{U}|\nabla u(x, t)|^{p-1} d x
$$

where, for any $v \in L^{1}(U)$, we have set $\bar{v}=|U|^{-1} \int_{U} v d x$. Hence

$$
\int_{t}^{\theta} \int_{U}|u(x, t)|^{p-1} d x \leqslant C \int_{t}^{\theta} \int_{U}|\nabla u(x, t)|^{p-1} d x+C
$$

follows from (2.18), and (2.19) holds.
(ii) Assume (2.19) and (2.20). Then (2.18) holds: indeed, we can chose $\theta>0$ such that $\left.\int_{\Omega} u(x, \theta)\right) \zeta^{\tau}(x) d x$ is finite, since $u \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right)$.
(iii) Assume that (2.18) and (2.19) hold. Let $\alpha \in(\max (1-p,-1), 0)$ be fixed. From (2.7) we get for any $0<t<\theta$,

$$
\begin{align*}
& \frac{|\alpha|}{2} \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t \\
& \leqslant \frac{1}{\alpha+1} \int_{\Omega}(1+u(x, \theta))^{\alpha+1} \zeta^{\tau}(x) d x+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{p-1+\alpha} \zeta^{\tau} d x d t \\
& \quad+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{q+\alpha} \zeta^{\tau} d x d t \tag{2.25}
\end{align*}
$$

Since $(1+u)^{q+\alpha} \leqslant(1+u)^{q}$, and $(1+u)^{p-1+\alpha} \leqslant(1+u)^{p-1}$, we find

$$
\int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p \zeta^{\tau}} d x d t \leqslant C
$$

(and in fact for any $\alpha<0$ ), hence

$$
\begin{equation*}
\int_{0}^{\theta} \int_{U}(1+u)^{\alpha-1}|\nabla u|^{p} d x d t<\infty \tag{2.26}
\end{equation*}
$$

We recall the Gagliardo-Nirenberg estimate: for any $p, m \geqslant 1, \gamma \in[1,+\infty)$ and $s \in[0,1]$ such that

$$
\begin{equation*}
\frac{1}{\gamma}=s\left(\frac{1}{p}-\frac{1}{N}\right)+\frac{1-s}{m} \tag{2.27}
\end{equation*}
$$

there exists $c=c(N, p, m, s, \Omega)>0$ such that for any $v \in W^{1, p}(U) \cap L^{m}(U)$,

$$
\begin{equation*}
\|v-\bar{v}\|_{L^{\gamma}(U)} \leqslant c\| \| \nabla v\| \|_{L^{p}(U)}^{s}\|v-\bar{v}\|_{L^{m}(U)}^{1-s} . \tag{2.28}
\end{equation*}
$$

We apply it to $v$ where

$$
\begin{equation*}
v(x, t)=(1+u(x, t))^{\beta}, \quad \beta=(\alpha+p-1) / p . \tag{2.29}
\end{equation*}
$$

Choosing

$$
\begin{equation*}
\gamma=p+\frac{p}{N \beta}, \quad s=\frac{p}{\gamma}, \quad m=\frac{1}{\beta}, \tag{2.30}
\end{equation*}
$$

for which (2.27) holds, then

$$
\begin{equation*}
\|\bar{v}(\cdot)\|_{L^{\infty}((0, \theta))} \leqslant C \tag{2.31}
\end{equation*}
$$

is derived from (2.18) and the Hölder inequality, because $\beta \in(0,1)$. Therefore, for almost all $t \in(0, \theta)$,

$$
\begin{aligned}
\int_{U}\left|(1+u(x, t))^{\beta}-\bar{v}(t)\right|^{\gamma} d x \leqslant & C\left(\int_{U}(1+u(x, t))^{\alpha-1}|\nabla u(x, t)|^{p} d x\right)^{s \gamma / p} \\
& \times\left(\int_{U}\left|(1+u(x, t))^{\beta}-\bar{v}(t)\right|^{1 / \beta} d x\right)^{(1-s) \gamma \beta} \\
\leqslant & C\left(\int_{U}(1+u(x, t))^{\alpha-1}|\nabla u(x, t)|^{p} d x\right) \\
& \times\left(\int_{U}\left|(1+u(x, t))^{\beta}-\bar{v}(t)\right|^{1 / \beta} d x\right)^{(1-s) \gamma \beta}
\end{aligned}
$$

and

$$
\int_{U}\left|(1+u(x, t))^{\beta}-\bar{v}(t)\right|^{1 / \beta} d x \leqslant C, \quad \int_{U} \bar{v}(t)^{\gamma} d x=|U| \overline{v_{n}}(t)^{\gamma} \leqslant C
$$

from (2.18) and (2.31). Therefore

$$
\int_{U}(1+u(x, t))^{\beta \gamma} d x \leqslant C \int_{U}(1+u(x, t))^{\alpha-1}|\nabla u(x, t)|^{p} d x+C .
$$

Integrating on $(0, \theta)$ we get

$$
\int_{0}^{\theta} \int_{U}(1+u(t))^{\beta \gamma} d x \leqslant C \int_{0}^{\theta} \int_{U}(1+u)^{\alpha-1}|\nabla u|^{p} d x d t+C \theta \leqslant C
$$

from (2.26). If we set $\sigma=\beta \gamma$, this means

$$
\int_{0}^{\theta} \int_{U}(1+u(t))^{\sigma} d x \leqslant C
$$

for any $\sigma$ such that $p / N+\max (p-2,0)<\sigma<p-1+p / N=q_{\mathrm{c}}$, and for any $0<\sigma<q_{\mathrm{c}}$ by the Hölder inequality. This proves (2.21).

Next for any $0<r<p$, and any $\alpha<0$, we find

$$
\begin{align*}
\int_{0}^{\theta} \int_{U}|\nabla u|^{r} d x \leqslant & \left(\int_{0}^{\theta} \int_{U}(1+u)^{\alpha-1}|\nabla u|^{p} d x d t\right)^{r / p} \\
& \times\left(\int_{0}^{\theta} \int_{U}(1+u)^{(1-\alpha) r /(p-r)} d x\right)^{(p-r) / p} \tag{2.32}
\end{align*}
$$

Hence

$$
\int_{0}^{\theta} \int_{U}|\nabla u|^{r} d x \leqslant C
$$

holds if $r$ is such that $0<r<N q_{\mathrm{c}} /(N+1)$. This proves (2.22), which implies (2.20) in particular, since $p-1<p-N /(N+1)$. Now from (2.4) with $h=1$, for any $\xi \in C_{\mathrm{c}}^{\infty}(\Omega)$ and any $0<t<\theta<T$,

$$
\begin{align*}
\int_{\Omega} u(x, t) \xi(x) d x= & \int_{\Omega} u(x, \theta) \xi(x) d x \\
& +\int_{t}^{\theta} \int_{\Omega}\left(|\nabla u|^{p-2} \nabla u \cdot \nabla \xi+u^{q} \xi\right) d x d t \tag{2.33}
\end{align*}
$$

Because the right-hand side of (2.11) has a finite limit when $t \rightarrow 0$, the same holds with $t \mapsto \int_{\Omega} u(x, t) \xi(x) d x$. The mapping $\xi \mapsto \lim _{t \rightarrow 0} \int_{\Omega} u(x, t)$ $\xi(x) d x$ is clearly a positive linear functional over the space $C_{\mathrm{c}}^{\infty}(\Omega)$. It can be extended in a unique way as a Radon measure $\ell$ on $\Omega$, and (2.23) holds in $\Omega$.

Finally, let $0<t<\theta$ be fixed. Taking $h=1$ in (2.3) with $\varphi \in C_{\mathrm{c}}^{\infty}$ $(\Omega \times[0, T))$, infers

$$
\begin{aligned}
& \left.\int_{t}^{\theta} \int_{\Omega}\left(-u \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u \cdot \nabla u \varphi\right)+u^{q} \varphi\right) d x d t \\
& \quad=\int_{\Omega} u(x, t) \varphi(x, t) d x-\int_{\Omega} u(x, \theta) \varphi(x, \theta) d x
\end{aligned}
$$

Letting $t$ go to 0 in the left-hand side of the above equality and using (2.18)-(2.20) yields to

$$
\mid \int_{\Omega} u(x, t)\left(\varphi(x, t)-\varphi(x, 0) d x \mid \leqslant C t \int_{\Omega} u(x, t) d x\right.
$$

Therefore,

$$
\int_{\Omega} u(x, t) \varphi(x, t) d x \rightarrow \int_{\Omega} \varphi(x, 0) d \ell(x)
$$

from (2.23). This proves (2.24), and the proof is complete.

Remark 2.4. Estimates (2.21) and (2.22) are still valid for the weak solutions defined in Remark 2.1, by considering $T_{k}(u)$ and going to the limit as $k \rightarrow \infty$, and using the definition of the gradient. Next, we consider the class $\Sigma$ of solutions introduced in [8] with $a \equiv 0$ and $p<2$, defined by the conditions

$$
u \in C\left((0, T), L_{\mathrm{loc}}^{1}(\Omega), \quad \nabla T_{k}(u) \in L_{\mathrm{loc}}^{p}\left(Q_{T}\right), \quad \partial_{t}\left(T_{k}(u)\right) \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right)\right.
$$

and satisfying

$$
\begin{equation*}
\int_{0}^{T} \int_{\Omega}\left((\varphi-u)_{+} \partial_{t} u+|\nabla u|^{p-2} \nabla u \cdot \nabla\left((\varphi-u)_{+}\right)\right) d x d t=0 \tag{2.34}
\end{equation*}
$$

for any $\varphi \in C_{\mathrm{c}}^{\infty}\left(Q_{T}\right)$. By a straightforward argument, it is shown in [8, Lemma 3.2] that any $u \in \Sigma$ satisfies $(1+u)^{\alpha-1}|\nabla u|^{p} \in L_{\text {loc }}^{1}\left(Q_{T}\right)$ for any $\alpha \in(1-p, 0)$. Then, by the Gagliardo-Nirenberg estimate applied on $(t, \theta)$ instead of $(0, \theta)$, we deduce that $u^{\sigma} \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right)$ for any $\sigma \in\left(0, q_{\mathrm{c}}\right)$, and $|\nabla u|^{r} \in$ $L_{\mathrm{loc}}^{1}\left(Q_{T}\right)$ for any $r \in\left(0, N q_{\mathrm{c}} /(N+1)\right)$. In particular $|\nabla u|^{p-1} \in L_{\mathrm{loc}}^{1}\left(Q_{T}\right)$. By [8, Proposition 2.1],

$$
\int_{0}^{T} \int_{\Omega}\left(\varphi \partial_{t}(H(u))+|\nabla u|^{p-2} \nabla u \cdot \nabla(h(u) \varphi)\right) d x d t=0
$$

for any $\varphi \in C_{\mathrm{c}}^{\infty}\left(Q_{T}\right)$ and any function $h \in C_{\mathrm{c}}^{1}(\mathbb{R}) \cap W^{1, \infty}(\mathbb{R})$ where $H^{\prime}(r)=h(r)$. Consequently,

$$
\begin{equation*}
\int_{0}^{T} \int_{\Omega}\left(-H(u) \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u \cdot \nabla(h(u) \varphi)\right) d x d t=0 \tag{2.35}
\end{equation*}
$$

In order to prove uniqueness of solutions with a given initial data $\mu \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N}\right)$, they introduced in [8, Proposition 2.1] a class $\Sigma^{*} \subset \Sigma$ of functions which satisfies

$$
\lim _{k \rightarrow \infty}\left(\int_{0}^{T} \int_{K \cap\{k<u<k+1\}}|\nabla u|^{p} d x d t\right)=0
$$

for any compact $K \subset Q_{T}$. Since it can be proved that (2.35) remains true for any $h \in C(\mathbb{R}) \cap W^{1, \infty}(\mathbb{R})$ constant outside $[-k, k]$ for some $k>0$, any $u \in \Sigma^{*}$ is a weak solution.

Next, we consider the class of solutions introduced in [15] for initial data $\mu \in L^{1}(\Omega)$ for bounded $\Omega$. They satisfy $u \in C\left([0, T], L^{1}(\Omega)\right)$ and (2.35) for any $\varphi \in C_{\mathrm{c}}^{\infty}\left(Q_{T}\right)$ and any function $h \in C_{\mathrm{c}}^{1}(\mathbb{R}) \cap W^{1, \infty}(\mathbb{R})$ where $H^{\prime}(r)=h(r)$, with

$$
\lim _{k \rightarrow \infty}\left(\int_{0}^{T} \int_{\{k<u<k+1\}}|\nabla u|^{p} d x d t\right)=0
$$

Because these above integrals are bounded, for any $\alpha<0$, there holds

$$
\begin{aligned}
& \int_{0}^{\theta} \int_{U}(1+u)^{\alpha-1}|\nabla u|^{p} d x d t \\
& \quad \leqslant \sum_{k=0}^{\infty}(1+k)^{\alpha-1} \int_{0}^{\theta} \int_{\{k<u<k+1\}}|\nabla u|^{p} d x d t \leqslant C \sum_{k=0}^{\infty}(1+k)^{\alpha-1}
\end{aligned}
$$

By (2.32) in the proof of Proposition 2.2, it follows that (2.20) is valid, hence also (2.19), and $u$ is a weak solution of the problem.

## 3. EXISTENCE OF THE INITIAL TRACE

The main result of this section which settled the basis of the definition of the initial trace of a nonnegative solution of (1.1) is the following.

THEOREM 3.1. Let $u$ be a nonnegative weak solution of (1.1) in $Q_{T}$. Then for any $y \in \Omega$ the following alternative holds:
(i) either for any open subset $U \subset \Omega$ containing $y$

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{U} u(x, t) d x=\infty \tag{3.1}
\end{equation*}
$$

(ii) or there exist an open neighborhood $U^{*} \subset \Omega$ of $y$ and a nonnegative Radon measure $\ell_{U^{*}} \in \mathscr{M}^{+}\left(U^{*}\right)$ such that for any $\xi \in C_{\mathrm{c}}\left(U^{*}\right)$,

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{U^{*}} u(x, t) \xi(x) d x=\ell_{U^{*}}(\xi) \tag{3.2}
\end{equation*}
$$

and (2.21) and (2.22) hold in any open set $U \subset \subset U^{*}$.
It will be proved below and precised according to the different values of $q$ with respect to $p-1$. With Theorem 3.1 we can define a set $\mathscr{R}$ (depending on $u$ ) by

$$
\begin{equation*}
\mathscr{R}=\left\{y \in \Omega: \exists U \text { open } \subset \Omega, \quad y \in U, \quad \limsup _{t \rightarrow 0} \int_{U} u(x, t) d x<\infty\right\} \tag{3.3}
\end{equation*}
$$

Then $\mathscr{R}$ is an open subset of $\Omega$ and there exists a unique Radon measure $\mu \in \mathscr{M}^{+}(\mathscr{R})$ such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathscr{R}} u(x, t) \xi(x) d x=\int_{\mathscr{R}} \xi(x) d \mu(x), \quad \forall \xi \in C_{\mathrm{c}}(\mathscr{R}) \tag{3.4}
\end{equation*}
$$

By Proposition 2.2, $u$ satisfies

$$
\begin{align*}
& \left.\int_{0}^{\theta} \int_{\mathscr{R}}\left(-u \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u \cdot \nabla u \varphi\right)+u^{q} \varphi\right) d x d t \\
& \quad=\int_{\mathscr{R}} \varphi(x, 0) d \mu(x)-\int_{\mathscr{R}} u(x, \theta) \varphi(x, \theta) d x \tag{3.5}
\end{align*}
$$

for any $0<\theta<T$ and $\varphi \in C_{\mathrm{c}}^{\infty}(\mathscr{R} \times[0, T))$, and (2.21) and (2.22) hold in any open set $U \subset \subset \mathscr{R}$.

The next definition is parallel to the one introduced by Marcus and Véron [33-35].

Definition 3.1. Let $u$ be a nonnegative weak solution of (1.1) in $Q_{T}$. A point $y \in \Omega$ is called a regular point if $y \in \mathscr{R}$. Otherwise it is called a singular point. The set of singular points is denoted $\mathscr{S}=\Omega \backslash \mathscr{R}$; it is a relatively closed subset of $\Omega$. We shall denote

$$
\operatorname{tr}_{\Omega}(u)=(\mathscr{S}, \mu)
$$

where $\mu$ is the Radon measure in (3.4) and call it the initial trace of $u$ at $t=0$.

The initial trace can also be represented in a unique way by a positive, outer regular Borel measure $v$ with the following values on any Borel subset $A$ :

$$
v(A)= \begin{cases}\infty & \text { if } A \cap \mathscr{S} \neq \emptyset  \tag{3.6}\\ \mu(A) & \text { if } A \subset \mathscr{R}\end{cases}
$$

It is known that there is a one-to-one correspondence between the sets

$$
\mathscr{B}_{\text {reg }}^{+}(\Omega)=\{v \text { outer regular Borel measure on } \Omega, v \geqslant 0\}
$$

and

$$
C M^{+}(\Omega)=\left\{(\mathscr{S}, \mu): \mathscr{S} \text { relatively closed in } \Omega, \mu \in \mathscr{M}^{+}(\Omega \backslash \mathscr{S})\right\}
$$

When considering the initial trace of $u$ as an element of $\mathscr{B}_{\text {reg }}^{+}(\Omega)$, we shall denote it by $\operatorname{Tr}_{\Omega}(u)=v$. If $\mathscr{S}=\emptyset$, we will say that $u$ admits for initial trace the Radon measure $\mu$.

### 3.1. The Case $q>p-1>0$

In this case the proof of Theorem 3.1 is based upon the following lemma.
Lemma 3.2. Let $q>p-1>0$. Under the assumptions of Theorem 3.1, let $\tau>p q /(q-p-1)$. For any nonnegative function $\zeta \in C_{\mathrm{c}}^{\infty}(\Omega)$ the following dichotomy occurs:
(i) either $\int_{0}^{T} \int_{\Omega} u^{q} \zeta^{\tau} d x d t<\infty$, then

$$
t \mapsto \int_{\Omega} u(x, t) \zeta^{\tau}(x) d x
$$

remains bounded near $t=0$.
(ii) or $\int_{0}^{T} \int_{\Omega} u^{q} \zeta^{\tau} d x d t=\infty$, then

$$
\lim _{t \rightarrow 0} \int_{\Omega} u(x, t) \zeta^{\tau}(x) d x=\infty
$$

Proof. Step 1. Let $\alpha \in\left(\max \left((1-p),-\frac{1}{2}\right), 0\right)$ and $0<\theta<T$. We start from (2.7). Since $d=(q+\alpha) /(\alpha+p-1)>1$ and $\tau>p d^{\prime}=p d /(d-1)$, it follows
by Hölder's inequality

$$
\begin{align*}
& \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha+p-1} \zeta^{\tau-p}|\nabla \zeta|^{p} d x d t \\
& \quad \leqslant \int_{t}^{\theta} \int_{\Omega}\left((1+u)^{q+\alpha} \zeta^{\tau}+\zeta^{\tau-p d^{\prime}}|\nabla \zeta|^{p d^{\prime}}\right) d x d t \\
& \quad \leqslant C+\int_{t}^{\theta} \int_{\Omega}(1+u)^{q+\alpha} \zeta^{\tau} d x d t \tag{3.7}
\end{align*}
$$

where $C=C(\alpha, p, q, \tau, \zeta)$. Therefore from (2.7), since $\alpha>-1$,

$$
\begin{align*}
& \frac{|\alpha|}{2} \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t \\
& \quad \leqslant \frac{1}{\alpha+1} \int_{\Omega}(1+u(x, \theta))^{\alpha+1} \zeta^{\tau}(x) d x+C+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{q+\alpha} \zeta^{\tau} d x d t \tag{3.8}
\end{align*}
$$

with another $C=C(\alpha, p, q, \tau, \zeta)$.
Step 2. Next we use (2.15). Since $q>p-1$, we can choose $\alpha$ such that $\delta=q /(1-\alpha)(p-1)>1$, and $|\alpha|$ small enough in order to have $\tau>p \delta^{\prime} \geqslant p d^{\prime}$ (remember that from assumption $\tau>p q /(q-p+1)$, and notice that $\delta^{\prime} \geqslant d^{\prime}$ is equivalent to $p q+1>p-\alpha p+\alpha$ since $\alpha<0$ ). Therefore, by Hölder's inequality,

$$
\begin{align*}
& \int_{t}^{\theta} \int_{\Omega}|\nabla \zeta|^{p}(1+u)^{(1-\alpha)(p-1)} \zeta^{\tau-p} d x d t \\
& \quad \leqslant \int_{t}^{\theta} \int_{\Omega}(1+u)^{q} \zeta^{\tau} d x d t+\int_{t}^{\theta} \int_{\Omega}|\nabla \zeta|^{p \delta^{\prime}} \zeta^{\tau-p \delta^{\prime}} d x d t \tag{3.9}
\end{align*}
$$

Combining this estimate with (2.8) and (3.7), (3.8) yields

$$
\begin{equation*}
\int_{\Omega} u(x, t) \zeta^{\tau}(x) d x \leqslant C+\int_{\Omega} u(x, \theta) \zeta^{\tau}(x) d x+C \int_{t}^{\theta} \int_{\Omega} u^{q \zeta^{\tau}} d x d t \tag{3.10}
\end{equation*}
$$

where $C=C(\tau, p, q, \alpha, \zeta)$.
Step 3. We start from (3.8) and the converse estimate (2.9). Using (3.8) and Hölder's inequality as in (2.15), (3.9), for any $\varepsilon>0$,

$$
\begin{aligned}
& \tau \int_{t}^{\theta} \int_{\Omega} \zeta^{\tau-1}|\nabla u|^{p-1}|\nabla \zeta| d x d t \\
& \quad \leqslant \varepsilon \frac{|\alpha|}{2} \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t+\varepsilon \int_{t}^{\theta} \int_{\Omega} u^{q} \zeta^{\tau} d x d t+C_{\varepsilon} \\
& \quad \leqslant 2 \varepsilon \int_{\Omega}(u(x, \theta)) \zeta^{\tau}(x) d x+\varepsilon(1+C) \int_{t}^{\theta} \int_{\Omega} u^{q} \zeta^{\tau} d x d t+C_{\varepsilon}
\end{aligned}
$$

where $C_{\varepsilon}=C_{\varepsilon}(\tau, p, q, \alpha, \zeta, \varepsilon)$. If we choose $\varepsilon$ such that $\frac{1}{4}-2 \varepsilon>\frac{1}{8}$ and $\frac{1}{2}-\varepsilon(1+C)>\frac{1}{4}$, and combine this inequality with (2.9), we obtain

$$
\begin{equation*}
\int_{\Omega} u(x, \theta) \zeta^{\tau}(x) d x+\int_{t}^{\theta} \int_{\Omega} u^{q} \zeta^{\tau} d x d t \leqslant 8 \int_{\Omega} u(x, t) \zeta^{\tau}(x) d x+C \tag{3.11}
\end{equation*}
$$

where $C=C(\tau, p, q, \alpha, \zeta)$.
Step 4. Assume first that $\int_{0}^{T} \int_{\Omega} u^{q} \zeta^{\tau} d x d t<\infty$, and let $0<\theta<T$ such that $\int_{\Omega} u(x, \theta) \zeta^{\tau} d x<\infty$. Then

$$
\begin{equation*}
\int_{\Omega} u(x, t) \zeta^{\tau}(x) d x<\infty \tag{3.12}
\end{equation*}
$$

for almost all $0<t<\theta$, from (3.10). Since $\theta$ can be taken arbitrarily close to $T$, (2.11) holds a.e. on $(0, T)$. Suppose now that $\int_{0}^{T} \int_{\Omega} u^{q} \zeta^{\tau} d x d t=\infty$, and that there exists a sequence $\left\{t_{n}\right\}$ converging to 0 such that $\int_{\Omega} u\left(x, t_{n}\right) \zeta^{\tau} \times$ $(x) d x \leqslant M$ for some constant $M$. Taking $t=t_{n}$ and $T=\theta$ in (2.11) yields a contradiction.

Remark 3.1. The result is still valid for the weaker solutions defined in Remark 2.2. Indeed, estimates (3.8) and (3.9) remain valid with $u$ replaced by $T_{k}(u)$, the term $u^{q}$ remaining unchanged. Therefore (3.10) holds under the form

$$
\begin{align*}
& \int_{\Omega} T_{k}(u(x, t)) \zeta^{\tau}(x) d x \\
& \quad \leqslant C+C \int_{\Omega} T_{k}(u(x, \theta)) \zeta^{\tau}(x) d x+C \int_{t}^{\theta} \int_{\Omega} u^{q} \zeta^{\tau} d x d t \tag{3.13}
\end{align*}
$$

In estimate (2.16) $u^{q}$ can be replaced by $T_{k}(u)^{q}$ which is smaller. Therefore (2.17) holds with $u$ replaced by $T_{k}(u)$. Since (3.9) holds with $(1+u)^{q}$ replaced by $\left(1+T_{k}(u)\right)^{q}$, (2.11) is valid under the form

$$
\begin{align*}
& \frac{1}{8} \int_{\Omega} T_{k}(u(x, \theta)) \zeta^{\tau}(x) d x+\frac{1}{4} \int_{t}^{\theta} \int_{\Omega}\left(T_{k}(u)\right)^{q} \zeta^{\tau} d x d t \\
& \quad \leqslant \int_{\Omega} T_{k}(u(x, t)) \zeta^{\tau}(x) d x+C \tag{3.14}
\end{align*}
$$

Letting $k \rightarrow \infty$ in (3.13) and (2.11) implies that (3.10) and (2.11) are still valid and the conclusion of Lemma 3.2 follows.

Remark 3.2. The results make use essentially of the term $u^{q}$ in Eq. (1.1). They extend to Eq. (2.1) under the condition $a(x) \geqslant C>0$ a.e. in $\Omega$.

Proof of Theorem 3.1. Case $q>p-1>0$. We first assume that for any open subset $U$ of $\Omega$ containing $y$ and any nonnegative $\zeta \in C_{\mathrm{c}}^{\infty}(U)$ with value 1 in a neighborhood of $y$, and $\tau>p q /(q-p-1)$,

$$
\int_{0}^{T} \int_{U} u^{q} \zeta^{\tau} d x d t=\infty
$$

Then (3.1) holds from Lemma 3.2.
Assume now that there exists an open neighborhood $\tilde{U} \subset \Omega$ of $y$ and a $C_{\mathrm{c}}^{\infty}(\tilde{U})$-function $\zeta$ with value in $[0,1]$ and value 1 in a neighborhood $U_{\zeta}=U^{*}$ of $y$ such that

$$
\int_{0}^{T} \int_{\tilde{U}} u^{q \zeta^{\tau}} d x d t<\infty
$$

Then

$$
t \mapsto \int_{U^{*}} u(x, t) d x
$$

remains bounded in a neighborhood of $t=0$ from Lemma 3.2. Moreover, we have

$$
\begin{equation*}
\int_{0}^{T} \int_{U^{*}}|\nabla u|^{p-1} d x d t<\infty \tag{3.15}
\end{equation*}
$$

Indeed from (2.12),

$$
\begin{align*}
\int_{t}^{\theta} \int_{\Omega}|\nabla u|^{p-1} \zeta^{\tau} d x d t \leqslant & \int_{t}^{\theta} \int_{\Omega}|\nabla u|^{p}(1+u)^{\alpha-1} \zeta^{\tau} d x d t \\
& +\int_{t}^{\theta} \int_{\Omega}(1+u)^{(1-\alpha)(p-1)} \zeta^{\tau} d x d t \tag{3.16}
\end{align*}
$$

and, without using Proposition 2.2 but the assumption $q>p-1$,

$$
\int_{t}^{\theta} \int_{\Omega}(1+u)^{(1-\alpha)(p-1)} \zeta^{\tau} d x d t \leqslant \int_{t}^{\theta} \int_{\Omega}(1+u)^{q} \zeta^{\tau} d x d t
$$

by choosing $\alpha$ such that $(1-\alpha)(p-1) \leqslant q$. Then (3.15) follows from (3.8). Thus we can apply Proposition 2.2 in $U^{*}$, and deduce (3.2) and the regularity results.

### 3.2. The Case $q \leqslant 1$ and $p<2$

In this case we prove that any solution of (1.1) admits an initial trace reduced to a Radon measure. Moreover, we can treat the case of (2.1).

THEOREM 3.3. Let $0<q \leqslant 1$ and $p<2$. Let $u$ be a nonnegative weak solution of (2.1) in $Q_{T}$, with $a \in L_{\mathrm{loc}}^{\infty}(\Omega)$. Then there exists a Radon measure $\mu \in \mathscr{M}^{+}(\Omega)$ such that

$$
\lim _{t \rightarrow 0} \int_{\Omega} u(x, t) \xi(x) d x=\int_{\Omega} \xi(x) d \mu(x), \quad \forall \xi \in C_{\mathrm{c}}(\Omega)
$$

Then (2.24) holds, (2.21) and estimates (2.22) are satisfied in any open set $U \subset \subset \Omega$.

Proof. We claim that for any open subset $U$ with compact closure $\bar{U} \subset \Omega$,

$$
t \mapsto \int_{U} u(x, t) d x
$$

remains bounded near $t=0$. Let $\alpha, t, \theta, \zeta$ and $\tau$ be as in Proposition 2.1 with $\alpha>-1$ and $\tau>p /(2-p)$. We have

$$
\begin{align*}
& \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha+p-1} \zeta^{\tau-p}|\nabla \zeta|^{p} d x d t \\
& \leqslant \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha+1} \zeta^{\tau} d x d t \\
& \quad+\int_{t}^{\theta} \int_{\Omega} \zeta^{\tau-(1+\alpha) p /(2-p)}|\nabla \zeta|^{(1+\alpha) p /(2-p)} d x d t \tag{3.17}
\end{align*}
$$

since $p<2$, and

$$
\int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha+q} \zeta^{\tau} d x d t \leqslant \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha+1} \zeta^{\tau} d x d t
$$

since $q \leqslant 1$. Owing to (2.7) it follows that

$$
\begin{align*}
& \int_{\Omega}(1+u(x, t))^{\alpha+1} \zeta^{\tau}(x) d x \\
& \quad \leqslant \int_{\Omega}(1+u(x, \theta))^{\alpha+1} \zeta^{\tau}(x) d x+C \int_{t}^{\theta} \int \Omega(1+u)^{1+\alpha} \zeta^{\tau} d x d t+C^{\prime} \tag{3.18}
\end{align*}
$$

where $C, C^{\prime}$ depend on $\alpha, \zeta,\|a\|_{L^{\infty}(\Omega)}$ and the exponents. If we set

$$
X(t)=\int_{t}^{\theta} \int_{\Omega}(1+u)^{1+\alpha} \zeta^{\tau} d x d t
$$

then with another $C>0$,

$$
\begin{equation*}
X^{\prime}+C X+C+\int_{\Omega} u(x, \theta) \zeta^{\tau}(x) d x \geqslant 0 \tag{3.19}
\end{equation*}
$$

Integrating (3.19) between $t$ and $\theta$ implies that $X(t)$ is bounded on $(0, \theta)$. Now we use (2.8). Since $p<2$ we can choose $\alpha$ such that $(1-\alpha)(p-1) \leqslant 1$ and we derive

$$
\begin{align*}
\int_{\Omega}(1+u(x, t)) \zeta^{\tau}(x) d x \leqslant & \int_{\Omega}(1+u(x, \theta)) \zeta^{\tau}(x) d x \\
& +C \int_{t}^{\theta} \int_{\Omega}(1+u) \zeta^{\tau} d x d t \tag{3.20}
\end{align*}
$$

Putting

$$
Y(t)=\int_{t}^{\theta} \int_{\Omega}(1+u) \zeta^{\tau} d x d t
$$

and integrating differential inequality (2.11) yields to

$$
\int_{\Omega}(1+u(x, t)) \zeta^{\tau}(x) d x \leqslant\left(e^{C(\theta-t)}+1\right) \int_{\Omega}(1+u(x, \theta)) \zeta^{\tau}(x) d x
$$

for $0<t \leqslant \theta$, which implies the claim. Then (2.18) holds. It implies (2.19), since $q \leqslant 1$ and $p<2$, hence we can apply Proposition 2.2 in $\Omega$.

Remark 3.3. Theorem 3.3 shows, in particular, that in the case $p-1<q \leqslant 1$, the first alternative of Theorem 3.1 cannot happen, which means $\mathscr{S}=\emptyset$ and $\mathscr{R}=\Omega$.

### 3.3. The Case $q \leqslant p-1, p>2$

In this range of exponents we can deal with the more general equation (2.1). The proof of Theorem 3.1 is a consequence of the following lemma.

Lemma 3.4. Assume $0<q \leqslant p-1$ and $p>2$. Let $u$ be any nonnegative weak solution of (2.1) in $Q_{T}$, with $a \in L_{\mathrm{loc}}^{\infty}(\Omega)$. Assume that for any open set $U \subset \subset \Omega$,

$$
t \mapsto \int_{U} u(x, t) d x
$$

remains bounded near $t=0$. Then for any $0<\theta<T$,

$$
\int_{0}^{\theta} \int_{U} u^{p-1}(x, t) d x+\int_{0}^{\theta} \int_{U}|\nabla u(x, t)|^{p-1} d x<\infty
$$

Proof. Let $\alpha \in(1-p, 0)$ be fixed, $\alpha \neq-1$, and $\zeta \in C_{\mathrm{c}}^{\infty}(\Omega)$ as above. We start from (2.25) where possibly $\alpha+1 \leqslant 0$. Therefore, by (2.18), since $q \leqslant p-1$,

$$
\begin{equation*}
\frac{|\alpha|}{2} \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t \leqslant C+C \int_{t}^{\theta} \int_{\Omega}(1+u)^{p-1+\alpha} \zeta^{\tau} d x d t \tag{3.21}
\end{equation*}
$$

If $U, U^{*}$ are any open sets with $U \subset \subset U^{*} \subset \subset \Omega$, taking $\zeta \in C_{\mathrm{c}}^{\infty}(\Omega)$ with $0 \leqslant \zeta \leqslant 1$, with value 1 on $U$, and 0 outside of $U^{*}$, we get

$$
\begin{align*}
\int_{t}^{\theta} \int_{U}(1+u)^{\alpha+p-1+p / N} d x d t & \leqslant C+C \int_{t}^{\theta} \int_{U}(1+u)^{\alpha-1}|\nabla u|^{p} d x d t \\
& \leqslant C+C \int_{t}^{\theta} \int_{U^{*}}(1+u)^{\alpha+p-1} d x d t \tag{3.22}
\end{align*}
$$

Hence any estimate of $(1+u)^{\alpha+p-1}$ in $L^{1}\left((0, \theta), L_{\text {loc }}^{1}(\Omega)\right)$ implies the same estimate for $(1+u)^{\alpha+p-1+p / N}$. We first take $\alpha_{0}=2-p$. Since $p+\alpha_{0}-1=1$,

$$
u^{\sigma_{1}} \in L^{1}\left((0, \theta), L_{\mathrm{loc}}^{1}(\Omega)\right)
$$

with

$$
\sigma_{1}=\alpha_{0}+p-1+p / N=\sigma_{1}=1+p / N
$$

Defining by induction,

$$
\alpha_{n+1}=\alpha_{n}+p / N, \quad \sigma_{n}=\alpha_{n}+p-1, \quad \forall n \in \mathbb{N}
$$

we get

$$
(1+u)^{\sigma_{n+1}} \in L^{1}\left((0, \theta) \times L_{\mathrm{loc}}^{1}(\Omega)\right)
$$

as long as $\alpha_{n}=n p / N+2-p<0$. Let $n_{0}$ be the largest integer such that $\alpha_{n}<0$. Then $(1+u)^{\sigma_{n_{0}+1}} \in L^{1}\left((0, \theta) \times L_{\mathrm{loc}}^{1}(\Omega)\right)$. And $\sigma_{n_{0}+1} \geqslant p-1$. Then in (3.21) one can now take $\alpha<0$ arbitrarily close to 0 and deduce that

$$
\frac{|\alpha|}{2} \int_{t}^{\theta} \int_{\Omega}(1+u)^{\alpha-1}|\nabla u|^{p} \zeta^{\tau} d x d t \leqslant C .
$$

Then from (3.22) we get $(1+u)^{\alpha+p-1+p / N} \in L^{1}\left((0, \theta) \times L_{\text {loc }}^{1}(\Omega)\right)$. In particular $u^{p-1} \in L^{1}(0, \theta), L_{\mathrm{loc}}^{1}(\Omega)$. From (2.32), we deduce that $|\nabla u| \in L^{r}\left((0, \theta) \times L_{\mathrm{loc}}^{1}\right.$ $(\Omega)$ ) for any $r$ such that

$$
(1-\alpha) r /(p-r) \leqslant \alpha+p-1+p / N
$$

Since $\alpha$ is arbitrary, this happens for any $r<q_{\mathrm{c}}$, in particular $|\nabla u| \in L^{p-1}\left((0, \theta) \times L_{\mathrm{loc}}^{1}(\Omega)\right)$.

Proof of Theorem 3.1. Case $q \leqslant p-1, p>2$ Let $y \in \Omega$. Then either statement (i) of Theorem 3.1 holds, or there exists an open subset $U^{*} \subset \Omega$ containing $y$ such that $\int_{U^{*}} u(x, t) d x$ is bounded near $t=0$. By Lemma 3.4 and Proposition 2.2 in $U^{*}$ statement (ii) holds.

## 4. COMPLEMENTARY PROPERTIES FOR $p>2$

### 4.1. The Case $q \leqslant 1<p-1, \Omega=\mathbb{R}^{N}$

In this section, we prove that the absorption term is negligible, therefore the initial trace of a solution is reduced to a Radon measure, equivalently $\mathscr{S}=\emptyset$.

Theorem 4.1. Assume $q \leqslant 1<p-1$ and $u \in C\left(\mathbb{R}^{N} \times(0, T)\right)$ is a nonnegative weak solution of $(1.1)$ in $\mathbb{R}^{N} \times(0, T)$. Then the initial trace of $u$ is a Radon measure $\mu \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$.

Proof. We show that for any $b \in \mathbb{R}^{N}$ there exists $\rho>0$ such that

$$
\begin{equation*}
\limsup _{t \rightarrow 0} \int_{B_{\rho}(b)} u(x, t) d x<\infty \tag{4.1}
\end{equation*}
$$

By contradiction we assume that it does not hold. Then there exists some $b \in \mathbb{R}^{N}$ such that for any $\rho>0$ there exists a sequence $\left\{t_{n, \rho}\right\}$ converging to 0 with the property that

$$
\begin{equation*}
\lim _{t_{n, \rho} \rightarrow 0} \int_{B_{\rho}(b)} u\left(x, t_{n, \rho}\right) d x=\infty \tag{4.2}
\end{equation*}
$$

Let $k>0$ be fixed. For any $\rho>0$ there exists $n_{\rho}$ such that for any $n \geqslant n_{\rho}$,

$$
\begin{equation*}
\int_{B_{\rho}(b)} u\left(x, t_{n, \rho}\right) d x \geqslant k \tag{4.3}
\end{equation*}
$$

By continuity of the integral with respect to the domain, there exists some $0<\tilde{\rho} \leqslant \rho$ such that

$$
\begin{equation*}
\int_{B_{\tilde{\rho}}(b)} u\left(x, t_{n_{\rho}, \rho}\right) d x=k \tag{4.4}
\end{equation*}
$$

Moreover, $\tilde{\rho}$ is uniquely determined if we impose it to be the largest as possible (in order to avoid the axiom of the choice in what follows). When $\rho \rightarrow 0, t_{n_{\rho}, \rho} \rightarrow 0$ since $t \mapsto u(., t)$ is continuous from $(0, T)$ into $L_{\text {loc }}^{1}\left(\mathbb{R}^{N}\right)$. Let
$w=w_{\rho}$ be the solution of

$$
\left\{\begin{array}{l}
\partial_{t} w-\nabla \cdot\left(|\nabla w|^{p-2} \nabla w\right)+w^{q}=0 \quad \text { in } \mathbb{R}^{N} \times(0, \infty),  \tag{4.5}\\
w(., 0)=u\left(., t_{n_{\rho}, \rho}\right) \chi_{B_{B_{\bar{\rho}}(b)}}, \quad \text { in } \mathbb{R}^{N} .
\end{array}\right.
$$

Since $u$ is nonnegative it follows by the comparison principle that

$$
u\left(x, t+t_{n, \rho}\right) \geqslant w_{\rho}(x, t) \quad \text { in } \mathbb{R}^{N} \times\left(0, T-t_{n, \rho}\right)
$$

When $\rho \rightarrow 0, w_{\rho}$ converges [29-32] to $\tilde{w}_{k}$ solution of

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{w}_{k}-\nabla \cdot\left(\left|\nabla \tilde{w}_{k}\right|^{p-2} \nabla \tilde{w}_{k}\right)+\tilde{w}_{k}^{q}=0 \quad \text { in } \mathbb{R}^{N} \times(0, \infty),  \tag{4.6}\\
\tilde{w}_{k}(., 0)=k \delta_{b} \quad \text { in } \mathbb{R}^{N}
\end{array}\right.
$$

By construction and uniqueness, the function $\tilde{w}_{k}$ is radial with respect to $b$, and radially decreasing for $t>0$. Moreover,

$$
u(x, t) \geqslant \tilde{w}_{k}(x, t) \quad \text { in } \mathbb{R}^{N} \times(0, T)
$$

By the comparison principle, the mapping $k \mapsto \tilde{w}_{k}$ is increasing. The key point is to notice that, since $0<q \leqslant 1, \lim _{k \rightarrow \infty} \tilde{w}_{k}=w_{\infty}$ a.e. cannot exist.

When $0<q<1$ there exists a scaling invariance of the equation, namely, if $w$ is a solution of (1.1), $N_{\ell}(w)$ defined by

$$
\begin{equation*}
N_{\ell}(w)(x, t)=\ell^{1 /(q-1)} w\left(b+\ell^{\gamma}(x-b), \ell t\right) \tag{4.7}
\end{equation*}
$$

with $\gamma=(q+1-p) /(p(q-1))$ and $\ell>0$, is also a solution, and consequently

$$
N_{\ell}\left(w_{k}\right)=w_{k \ell^{1 /(q-1)-\gamma N}} .
$$

Letting $k \rightarrow \infty$ leads to the invariance property

$$
N_{\ell}\left(w_{\infty}\right)=w_{\infty}, \quad \forall \ell>0
$$

from which follows the self-similar form of $w_{\infty}$,

$$
\begin{equation*}
w_{\infty}(x, t)=t^{1 /(1-q)} f\left(t^{-\gamma}(x-b)\right), \quad \forall(x, t) \in \mathbb{R}^{N} \times(0, \infty) \tag{4.8}
\end{equation*}
$$

This estimate implies, in particular, that $f(0)$ is finite and

$$
\tilde{w}_{k}(b, t) \leqslant t^{1 /(1-q)} f(0) \leqslant u(b, t), \quad \forall t \in(0, T)
$$

which contradicts the fact that $\tilde{w}_{k}(b, t) \rightarrow \infty$ when $t \rightarrow 0$, since $q<1$.
When $q=1$ (and $p \neq 2$ otherwise the result is well known), Eq. (1.1) is invariant with respect to the transformation $M_{\ell}(w)$ defined (for $\ell>0$ ) by

$$
M_{\ell}(w)(x, t)=\ell w\left(b+\ell^{(2-p) / p}(x-b), t\right)
$$

which yields to

$$
M_{\ell}\left(w_{k}\right)=w_{k \ell^{1+N(p-2) / p}}
$$

As in the previous case $w_{\infty}$ satisfies the invariant property

$$
M_{\ell}\left(w_{\infty}\right)=w_{\infty}, \quad \forall \ell>0
$$

This estimate implies

$$
0<w_{k}(b, T / 2) \leqslant w_{\infty}(b, T / 2)=\ell w_{\infty}(b, T / 2) \leqslant u(b, T / 2), \quad \forall \ell>0
$$

which is again a contradiction. If $U$ is any bounded open of $\mathbb{R}^{N}$ it follows by compactness that

$$
t \mapsto \int_{U} u(x, t) d x
$$

remains bounded near $t=0$, and the singular set is empty.
Remark 4.1. Actually this proof works for any $p>p_{0}$. But the case $p<2$ is still covered by Theorem 3.3.

Remark 4.2. The argument used below implies that if $u$ is any continuous, nonnegative solution of (1.12) in $\mathbb{R}^{N} \times(0, T)$, then for any relatively compact subset $U \in \mathbb{R}^{N}, \int_{U} u(x, t) d x$ remains bounded when $t$ remains bounded.

### 4.2. The Case $1<q \leqslant p-1, \Omega=\mathbb{R}^{N}$

Since $q>1$, the main difference with the previous section comes from the existence of the flat singular solution $W$ defined in (1.15).

Theorem 4.2. Let $1<q \leqslant p-1$. Assume $u \in C\left(\mathbb{R}^{N} \times(0, T)\right)$ is a nonnegative weak solution of $(1.1)$ in $\mathbb{R}^{N} \times(0, T)$. Then
(i) either $\mathscr{S}=\mathbb{R}^{N}$ and in fact $u \geqslant W$,
(ii) $\mathscr{S}=\emptyset$ and the initial trace of $u$ is a Radon measure $\mu \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$.

Proof. The scheme of the proof is very similar to the one of Theorem 4.1. Either for any $b \in \mathbb{R}^{N}$ there exists $\rho>0$ such that (4.1) holds, hence (2.18) holds. And Lemma 3.4 applies as above, since $q<p-1$ and $p>2$, thus (2.20) holds, and statement (ii) follows by Proposition 2.2. Or there exists some $b \in \mathbb{R}^{N}$ such that for any $\rho>0$ there exists a sequence $\left\{t_{n, \rho}\right\}$ converging to 0 satisfying (4.2). Let $k>0$ be fixed. For any $\rho$ there exists $t_{n_{\rho}, \rho}>0$ such that (4.3) holds, and there exists some $0<\tilde{\rho} \leqslant \rho$ such that (4.4) holds. Letting $\rho \rightarrow 0$, implies that $u$ is bounded from below by the solution
of (4.6), which does exist since $1<q<q_{\mathrm{c}}$, and the convergence result quoted in Theorem 4.1 applies. Since this estimate from below holds for any $k>0$, it follows that

$$
u \geqslant \tilde{w}_{\infty}=\lim _{k \rightarrow \infty} \tilde{w}_{k}
$$

Such a $\tilde{w}_{\infty}$ exists and is a solution of (1.1) by the classical theory. Moreover,

$$
\begin{equation*}
\tilde{w}_{k} \leqslant \tilde{w}_{\infty} \leqslant W \tag{4.9}
\end{equation*}
$$

Using the same scaling operator $N_{\ell}$ defined in (4.7) infers that $\tilde{w}_{\infty}$ is a selfsimilar solution of (1.1) which means

$$
\begin{equation*}
\tilde{w}_{\infty}(x, t)=t^{-1 /(q-1)} f\left(t^{-\gamma}(x-b)\right), \quad \forall(x, t) \in \mathbb{R}^{N} \times(0, \infty) \tag{4.10}
\end{equation*}
$$

with $\gamma=(q+1-p) /(p(q-1))$. The function $f$ is nonnegative and radially symmetric. Because of $(4.9) f$ is smaller than $(1 /(q-1))^{1 /(q-1)}$. Moreover, it satisfies the following equation:

$$
\left\{\begin{array}{l}
r^{1-N}\left(r^{N-1}\left|f^{\prime}\right|^{p-2} f^{\prime}\right)^{\prime}+\gamma r f^{\prime}+\frac{1}{q-1} f-f^{q}=0 \quad \text { in }(0, \infty)  \tag{4.11}\\
f^{\prime}(0)=0
\end{array}\right.
$$

Clearly $f \equiv(1 /(q-1))^{1 /(q-1)}$ is a solution and any solution satisfies in particular $f(0) \leqslant(1 /(q-1))^{1 /(q-1)}$. It remains to show that in any case, $f$ is constant. Suppose it were not true. Since $\gamma \leqslant 0 f$ is a nonincreasing neighborhood of $r=0$, and it remains decreasing on $(0, \infty)$ (by contradiction). If $\gamma<0$, then, writing

$$
\left(r^{N-1}\left|f^{\prime}\right|^{p-2} f^{\prime}\right)^{\prime} \leqslant-\gamma r^{N} f^{\prime}
$$

a straightforward but lengthy computation implies that

$$
f^{\prime}(r) \leqslant-|\gamma|^{1 /(p-2)} C_{N, p} r^{2 /(p-2)}
$$

for some $C_{N, p}>0$, contradicting the nonnegativity of $f$. If $\gamma=0$ (or equivalently $q=p-1$ ), $f$ admits a nonnegative limit $l$ at $\infty$. If this limit is not zero, there holds

$$
\left(r^{N-1}\left|f^{\prime}\right|^{p-2} f^{\prime}\right)^{\prime} \leqslant r^{N-1}\left(f^{p-1}-\frac{1}{p-2} f\right) \leqslant m r^{N-1}
$$

for some $m<0$ and $r \geqslant 1$. This inequality implies

$$
\left|f^{\prime}\right|^{p-2} f^{\prime}(r) \leqslant m \frac{r}{N}+C_{1}
$$

for $r>1$ and some real $C_{1}$, contradicting the nonnegativity of $f$. If $l=0$ there exists $R>0$ such that

$$
-\nabla \cdot\left(|\nabla f|^{p-2} \nabla f\right) \geqslant \frac{1}{2(p-2)} f \quad \text { in }\left\{y \in \mathbb{R}^{N}:|y|>R\right\}
$$

Such an inequality admits no nonnegative, nontrivial solution [12, 28]. Actually, this is a particular case of a more general result stating that any nonnegative solution of

$$
-\nabla \cdot\left(|\nabla f|^{p-2} \nabla f\right) \geqslant f^{s}
$$

in an exterior domain of $\mathbb{R}^{N}$ is identically zero whenever $0<s \leqslant N \times$ $(p-1) /(N-p)$ (no condition on $s$ if $p \geqslant N$ ). Thus $f$ is constant, hence $u \geqslant W$. In particular, for any bounded open set $U \subset \mathbb{R}^{N}$, $\lim _{t \rightarrow 0} \int_{U} u(x, t) d x=\infty$, which means that $\mathscr{S}=\mathbb{R}^{N}$.

Remark 4.3. If we assume that the solution $u$ of (1.1) is such that $u(., t)$ is bounded in $\mathbb{R}^{N}$ for any $t>0$, then $u \leqslant W$, thus the statement of Theorem 4.2 becomes
(i) either $u \equiv W$,
(ii) or the initial trace of $u$ is a measure $\mu \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$.

### 4.3. The Case $q \leqslant 1<p-1, \Omega$ Bounded

Although the absorption term is dominated by the diffusion one, the fact that $\Omega$ is bounded, infers that the singular set of the initial trace of a solution may be nonempty, and, in such a case, it is whole $\Omega$.

Theorem 4.3. Assume $q \leqslant 1<p-1, \Omega$ is bounded with a smooth boundary and $u$ is a continuous and nonnegative solution of (1.1) in $(0, T) \times \Omega$. Then
(i) either $\mathscr{S}=\Omega$ and, more precisely,

$$
\liminf _{t \rightarrow 0} t^{1 /(p-2)} u(x, t) \geqslant v_{\Omega}(x), \quad \forall x \in \Omega
$$

(ii) or $\mathscr{S}=\emptyset$ and the initial trace of $u$ is a Radon measure $\mu \in \mathscr{M}^{+}(\Omega)$.

For the proof, we start from the following classical result.

Lemma 4.4. Assume $p>2$ and $\Omega$ is bounded, then there exists a unique nonzero function $v=v_{\Omega} \in W_{0}^{1, p}(\Omega) \cap C^{1}(\bar{\Omega})$ satisfying

$$
\left\{\begin{array}{l}
-\nabla \cdot\left(|\nabla v|^{p-2} \nabla v\right)=\frac{1}{p-2} v \quad \text { in } \Omega  \tag{4.12}\\
v \geqslant 0 \quad \text { in } \Omega
\end{array}\right.
$$

Moreover, $v>0$ in $\Omega$.
The existence of a nonnegative nontrivial solution is obtained by minimization. The positivity follows by the Hopf boundary lemma as in [41] and the regularity from [24]. This function $v_{\Omega}$ plays a fundamental role in describing the set of solutions of (1.12) in $\Omega \times(0, \infty)$ since the function

$$
\begin{equation*}
(x, t) \mapsto V_{\Omega}(x, t)=t^{-1 /(p-2)} v_{\Omega}(x) \tag{4.13}
\end{equation*}
$$

is a solution with full blow-up at $t=0$ and vanishes on $\partial \Omega \times(0, \infty)$.
Lemma 4.5. Assume $0<q \leqslant p-1, p>2$ and $\Omega$ is bounded. For $k>0$ and $b \in \Omega$ let $\tilde{w}=\tilde{w}_{k, b}$ be the solution of

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{w}-\nabla \cdot\left(|\nabla \tilde{w}|^{p-2} \nabla \tilde{w}\right)+\tilde{w}^{q}=0 \quad \text { in } \Omega \times(0, \infty),  \tag{4.14}\\
\tilde{w}=0 \quad \text { in } \partial \Omega \times(0, \infty), \\
\tilde{w}(., 0)=k \delta_{b}(\cdot) \quad \text { in } \Omega .
\end{array}\right.
$$

Then $\lim _{k \rightarrow \infty} \tilde{w}_{k, b}=\tilde{w}_{\infty, b}$ exists; $\tilde{w}_{\infty, b}$ is a solution of (1.1) dominated by $V_{\Omega}$, and

$$
\lim _{t \rightarrow 0} t^{1 /(p-2)} \tilde{w}_{\infty, b}(x, t)=v_{\Omega}(x)
$$

uniformly on $\Omega$.
Proof. The function $\tilde{w}$ is a limit of solutions of (1.1) with smooth initial data and zero lateral boundary data. Therefore $\tilde{w}_{k, b} \leqslant V_{\Omega}$ in $\Omega \times(0, \infty)$. The sequence $\left\{\tilde{w}_{k, b}\right\}$ is increasing and its limit $\tilde{w}_{\infty, b}$ is a solution of (1.1) in $\Omega \times(0, \infty)$. Moreover $\tilde{w}_{\infty, b}$ is dominated by $V_{\Omega}$. For $\ell>0$ we set

$$
S_{\ell}(\tilde{w})(x, t)=\ell^{1 /(p-2)} \tilde{w}(x, \ell t)
$$

Then $S_{\ell}(\tilde{w})=\tilde{w}^{\ell}$ satisfies

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{w}^{\ell}-\nabla \cdot\left(\left|\nabla \tilde{w}^{\ell}\right|^{p-2} \nabla \tilde{w}^{\ell}\right)+\ell^{(p-1-q) /(p-2)}\left(\tilde{w}^{\ell}\right)^{q}=0 \quad \text { in } \Omega \times(0, \infty)  \tag{4.15}\\
\tilde{w}^{\ell}=0 \quad \text { in } \partial \Omega \times(0, \infty), \\
\tilde{w}^{\ell}(., 0)=\ell^{1 /(p-2)} k \delta_{b}(.) \quad \text { in } \Omega .
\end{array}\right.
$$

Since

$$
\partial_{t} \tilde{w}^{\ell}-\nabla \cdot\left(\left|\nabla \tilde{w}^{\ell}\right|^{p-2} \nabla \tilde{w}^{\ell}\right)+\left(\tilde{w}^{\ell}\right)^{q}=\left(1-\ell^{(p-1-q) /(p-2)}\right)\left(\tilde{w}^{\ell}\right)^{q},
$$

$S_{\ell}(\tilde{w})$ is a subsolution for $\ell \geqslant 1$, and therefore

$$
S_{\ell}\left(\tilde{w}_{k, b}\right) \leqslant \tilde{w}_{\ell^{1 /(p-2)}, b, b}
$$

Letting $k \rightarrow \infty$ yields

$$
S_{\ell}\left(\tilde{w}_{\infty, b}\right) \leqslant \tilde{w}_{\ell^{1 /(p-2)} \infty, b}
$$

Taking $t=\tau / \ell$ (with $\tau>0$ arbitrary) and replacing $1 / \ell$ by $t / \tau \in(0,1]$ yields to

$$
\tilde{w}_{\infty, b}(x, t) \geqslant\left(\frac{t}{\tau}\right)^{-1 /(p-2)} \tilde{w}_{\infty, b}(x, \tau) \quad \text { in } \Omega \times(0, \tau]
$$

This implies that $t \mapsto t^{1 /(p-2)} \tilde{w}_{\infty, b}(x, t)$ is nonincreasing. Since $\tilde{w}_{\infty, b}(b, \tau)>0$ for $\tau$ small enough, there holds

$$
\Phi(x) \leqslant t^{1 /(p-2)} \tilde{w}_{\infty, b}(x, t) \leqslant v_{\Omega}(x)
$$

for some function $x \mapsto \Phi(x)$ which is positive at $x=b$ (if $0<q<1$ all the functions $\tilde{w}_{k, a}$ vanishes for $t$ large enough). Put

$$
\psi(x, s)=t^{1 /(p-2)} \tilde{w}_{\infty, b}(x, t) \quad \text { where } s=\ln t
$$

then

$$
\left\{\begin{array}{l}
\partial_{s} \psi-\frac{1}{p-2} \psi-\nabla \cdot\left(|\nabla \psi|^{p-2} \nabla \psi\right)+e^{s(p-1-q) /(p-2)} \psi^{q}=0 \quad \text { in } \Omega \times \mathbb{R}  \tag{4.16}\\
\psi=0 \quad \text { on } \partial \Omega \times \mathbb{R}
\end{array}\right.
$$

Moreover, $s \mapsto \psi(x, s)$ is nonincreasing and $\Phi(x) \leqslant \psi(x, s) \leqslant v_{\Omega}(x)$ for $s \leqslant \ln \tau$. Therefore $\lim _{s \rightarrow-\infty} \psi(x, s)=\omega(x)$ exists for any $x \in \bar{\Omega}$. Since $\psi(., s)$ is bounded in $L^{\infty}(\Omega)$ uniformly with respect to $s$, it is bounded in $C^{1}(\bar{\Omega})$, which infers $\lim _{s \rightarrow-\infty} \psi(., s)=\omega(\cdot)$ in $C^{1}(\bar{\Omega})$. Multiplying (4.16) by
$\xi \in W_{0}^{1, p}(\Omega)$ and integrating on $\Omega \times(n-1, n)(n<0)$ yields to

$$
\begin{aligned}
\int_{\Omega} & (\psi(n, x)-\psi(n-1, x)) \xi(x) d x \\
& -\int_{n-1}^{n} \int_{\Omega}\left(\frac{1}{p-2} \psi \xi+|\nabla \psi|^{p-2} \nabla \psi \cdot \nabla \xi\right) d x d s \\
& =\int_{n-1}^{n} \int_{\Omega} e^{s(p-1-q) /(p-2)} \psi^{q} \xi d x d s
\end{aligned}
$$

Since $\lim _{s \rightarrow-\infty} \psi(., s)$ is independent of $s$, we obtain

$$
\int_{\Omega}\left(\frac{1}{p-2} \omega \xi+|\nabla \omega|^{p-2} \nabla \omega \cdot \nabla \xi\right) d x d s=0
$$

Since $\omega \geqslant \Phi$, it infers $\omega=v_{\Omega}$ by Lemma 4.4, which ends the proof.
Proof of Theorem 4.3. As in the proof of Theorem 4.2, either (ii) holds, or there exists $b \in \Omega$ such that $u \geqslant \tilde{w}_{k, b}$ in $Q_{T}$ for any $k>0$. In that case, Lemma 4.4 implies that (i) holds.

Remark 4.4. The above technique also applies to Eq. (1.12). Actually, Lemma 4.5 is valid under the following form: Assume $p>2, \Omega$ is bounded, and let $\tilde{w}=\tilde{w}_{k, b}$ be the solution of

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{w}-\nabla \cdot\left(|\nabla \tilde{w}|^{p-2} \nabla \tilde{w}\right)=0 \quad \text { in } \Omega \times(0, \infty),  \tag{4.17}\\
\tilde{w}=0 \quad \text { on } \partial \Omega \times(0, \infty) \\
\tilde{w}(., 0)=k \delta_{b}(\cdot) \quad \text { in } \Omega
\end{array}\right.
$$

Then $\lim _{k \rightarrow \infty} \tilde{w}_{k, b}=V_{\Omega}$ (see the proof of Lemma 4.5). Actually, the proof is simpler since Eq. (1.12) is invariant under the similarity transformation $S_{\ell}$. As a consequence the following result holds.

Theorem 4.6. Assume $2<p, \Omega$ is bounded with a smooth boundary and $u$ is a continuous and nonnegative solution of (1.12) in $(0, T) \times \Omega$. Then
(i) either $u(x, t) \geqslant V_{\Omega}(x, t)=t^{-1 /(p-2)} v_{\Omega}(x, t)$,
(ii) or the initial trace of $u$ is a Radon measure $\mu \in \mathscr{M}^{+}(\Omega)$.

### 4.4. The Case $1<q \leqslant p-1, \Omega$ Bounded

In this range of exponents the function $W$ defined by (1.15) is a singular solution of (1.1). Therefore, if $1<q<p-1$ we have to compare $u$ with the two functions $V_{\Omega}$ and $W$. In the case $q=p-1, V_{\Omega}$ has to be replaced by $W_{\Omega}$
defined by

$$
\begin{equation*}
W_{\Omega}(x, t)=t^{-1 /(p-2)} w_{\Omega}(x) \tag{4.18}
\end{equation*}
$$

where $w_{\Omega}$ is a solution of

$$
\left\{\begin{array}{l}
-\nabla \cdot\left(|\nabla w|^{p-2} \nabla w\right)+w^{p-1}=\frac{1}{p-2} w \quad \text { in } \Omega  \tag{4.19}\\
w \geqslant 0 \quad \text { on } \partial \Omega
\end{array}\right.
$$

Notice that $t^{-1 /(p-2)}=t^{-1 /(q-1)}$ and that the constant solution $(1 /(p-2))^{1 /(p-2)}$ is a particular solution of (4.19).

Theorem 4.7. Assume $q=p-1>2, \Omega$ is bounded with a smooth boundary and $u$ is a continuous and nonnegative solution of (1.1) in $(0, T) \times \Omega$. Then
(i) either $u(x, t) \geqslant W_{\Omega}(x, t)=t^{-1 /(p-2)} w_{\Omega}(x, t)$,
(ii) or the initial trace of $u$ is a Radon measure $\mu \in \mathscr{M}^{+}(\Omega)$.

The proof is similar to the one of Theorem 3.1, by using the next lemmas.
Lemma 4.8. Assume $p>2$ and $\Omega$ is bounded, then there exists a unique nonzero function $w=w_{\Omega} \in W_{0}^{1, p}(\Omega) \cap C^{1}(\bar{\Omega})$ satisfying (4.19).

The proof is similar to the one of Lemma 4.4.
Lemma 4.9. Assume $p>2$ and $\Omega$ is bounded, and let $\tilde{w}=\tilde{w}_{k, a}$ (for $k>0$ and some $a \in \Omega$ ) be the solution of

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{w}-\nabla \cdot\left(|\nabla \tilde{w}|^{p-2} \nabla \tilde{w}\right)+\tilde{w}^{p-1}=0 \quad \text { in } \Omega \times(0, \infty)  \tag{4.20}\\
\tilde{w}=0 \quad \text { in } \partial \Omega \times(0, \infty) \\
\tilde{w}(., 0)=k \delta_{a}(.) \quad \text { in } \Omega
\end{array}\right.
$$

Then $\lim _{k \rightarrow \infty} \tilde{w}_{k, a}=W_{\Omega}$ uniformly on $\Omega$.
The proof is simpler than the one of Lemma 4.5 since Eq. (1.1) and domain $\Omega$ are invariant under the similarity transformation $S_{\ell}$ defined by $S_{\ell}(\tilde{w})(x, t)=\ell^{1 /(p-2)} \tilde{w}(x, \ell t)$. This implies

$$
S_{\ell}\left(\tilde{w}_{k, a}\right)=\tilde{w}_{\ell^{1 /(p-2)} k, a}, \quad \forall k, \ell>0 .
$$

Therefore $\lim _{k \rightarrow \infty} \tilde{w}_{k, a}=\tilde{w}_{\infty, a}$ is self-similar since $S_{\ell}\left(\tilde{w}_{\infty, a}\right)=\tilde{w}_{\infty, a}$. Then

$$
\tilde{w}_{\infty, a}(x ; t)=t^{-1 /(p-2)} \tilde{w}_{\infty, a}(x, 1) .
$$

Because $w_{\infty, a}(x, 1)$ is not identically zero and satisfies (4.19), $w_{\infty, a}(x, 1)=w_{\Omega}$.
The case $1<q<p-1$ appears more difficult to deal with, and up to now we do not have a full answer (notice that here $t^{-1 /(p-2)}<t^{-1 /(q-1)}$ near $t=0$ ).

Theorem 4.10. Assume $1<q<p-1, \Omega$ is bounded with a smooth boundary and $u$ is a continuous and nonnegative solution of (1.1) in $(0, T) \times \Omega$. Then
(i) either

$$
\liminf _{t \rightarrow 0} t^{1 /(p-2)} u(x, t) \geqslant v_{\Omega}(x), \quad \forall x \in \Omega
$$

(ii) or the initial trace of $u$ is a Radon measure $\mu \in \mathscr{M}^{+}(\Omega)$. Moreover, if $\lim \sup _{t \rightarrow 0} t^{1 /(p-2)} u_{\partial \Omega}(x, t)=c$, for some $c>0$, then

$$
\begin{equation*}
\limsup _{t \rightarrow 0} t^{1 /(p-2)} u(x, t) \leqslant v_{\Omega}(x)+c . \tag{4.21}
\end{equation*}
$$

Finally, if $\lim \inf _{t \rightarrow 0} t^{1 /(q-1)} u_{\partial \alpha \Omega}(x, t)=\rho>0$ and $\lim _{t \rightarrow 0} u(x, t)=\infty$, uniformly in $\Omega$, then

$$
\begin{equation*}
\liminf _{t \rightarrow 0} t^{1 /(q-1)} u(x, t) \geqslant \min \left(\rho,(1 /(q-1))^{1 /(q-1)}\right) \tag{4.22}
\end{equation*}
$$

uniformly on $\Omega$.
The next result is an extension of Lemma 4.4.
Lemma 4.11. Assume $p>2$ and $\Omega$ is bounded. Then for any $m \geqslant 0$ there exists a unique nonzero function $v=v_{\Omega, a} \in W_{0}^{1, p}(\Omega) \cap C^{1}(\bar{\Omega})$ satisfying

$$
\left\{\begin{array}{l}
-\nabla \cdot\left(|\nabla v|^{p-2} \nabla v\right)=\frac{1}{p-2}(v+m) \quad \text { in } \Omega  \tag{4.23}\\
v \geqslant 0 \quad \text { on } \partial \Omega
\end{array}\right.
$$

Proof. Uniqueness comes again from the argument of [24]. For existence we consider the functional $J$ on $W_{0}^{1, p}(\Omega)$ defined by

$$
J(\psi)=\int_{\Omega}\left(\frac{1}{p}|\nabla \psi|^{p}-\frac{1}{2(p-2)}\left(\psi^{+}\right)^{2}-\frac{m}{p-2} \psi^{+}\right) d x .
$$

Clearly,

$$
J(\psi) \rightarrow \infty \quad \text { as }\|\psi\|_{W_{0}^{1, p}} \rightarrow \infty
$$

and since for $t>0$,

$$
J(t \psi)=\frac{t^{p}}{p} \int_{\Omega}|\nabla \psi|^{p} d x-\frac{t^{2}}{2(p-2)} \int_{\Omega}\left(\psi^{+}\right)^{2} d x-\frac{m t}{p-2} \int_{\Omega} \psi^{+} d x
$$

there exists $t_{1}>0$ such that $J\left(t_{1} \psi\right)>0$. Therefore, the infimum of $J$ in $W_{0}^{1, p}(\Omega)$ is negative and achieved at some $v$ for which there holds

$$
-\nabla \cdot\left(|\nabla v|^{p-2} \nabla v\right)=\frac{1}{p-2}\left(v^{+}+m \chi_{u>0}\right) \quad \text { in } \Omega .
$$

From the maximum principle $u \geqslant 0$ and consequently (4.11) holds.
Remark 4.5. An equivalent formulation states that for any $d \geqslant 0$ there exists a unique $\tilde{v}$ in $W^{1, p}(\Omega)$ solution of

$$
\left\{\begin{array}{l}
-\nabla \cdot\left(|\nabla \tilde{v}|^{p-2} \nabla \tilde{v}\right)=\frac{1}{p-2} \tilde{v} \quad \text { in } \Omega  \tag{4.24}\\
\tilde{v}=d \quad \text { on } \partial \Omega \\
\tilde{v} \geqslant d \quad \text { in } \Omega
\end{array}\right.
$$

Actually $\tilde{v}=v_{\Omega}+m$ with $m=d(p-2)$.
Proof of Theorem 4.10. The dichotomy between cases (i) and (ii) is clear from the proof of the previous theorems, therefore let us assume

$$
\limsup _{t \rightarrow 0} t^{1 /(p-2)} u_{\partial \Omega}(x, t)=m
$$

For any $\varepsilon>0$ there exists $t_{\varepsilon}>0$ such that

$$
\begin{equation*}
u_{\partial \Omega}(x, t) \leqslant(m+\varepsilon) t^{-1 /(p-2)} \tag{4.25}
\end{equation*}
$$

on $\Omega \times\left(0, t_{\varepsilon}\right]$. Let $\Omega^{\prime}$ be a relatively compact smooth open domain containing $\bar{\Omega}$ and denote by $v_{\Omega^{\prime}}$ the solution of problem (4.12) relative to $\Omega^{\prime}$. From uniqueness and positivity $v_{\Omega^{\prime}}>v_{\Omega}$ on $\bar{\Omega}$. Consequently, for any $0<\delta<t_{\varepsilon}$, there holds

$$
u(x, t) \leqslant\left(v_{\Omega^{\prime}}(x)+m+\varepsilon\right)(t-\delta)^{-1 /(p-2)}
$$

on $\left(\delta, t_{\varepsilon}\right]$. We can take $\Omega^{\prime}$ to be such that $\operatorname{dist}\left(\Omega, \partial \Omega^{\prime}\right)<1 / n$ and let $n \rightarrow \infty$. Then the sequence $\left\{v_{\Omega^{\prime}}\right\}=\left\{v_{\Omega^{\prime}, n}\right\}$ decreases and converges to $v_{\Omega}$. Letting
successively $\delta \rightarrow 0, n \rightarrow \infty, t \rightarrow 0$ and $\varepsilon \rightarrow 0$ yields to

$$
\begin{equation*}
\limsup _{t \rightarrow 0} t^{1 /(p-2)} u(x, t) \leqslant v_{\Omega}(x)+m \tag{4.26}
\end{equation*}
$$

For the last assertion we just have to notice that for $\sigma>0$, and $0<\tau<\min \left(\rho,(q-1)^{-1 /(q-1)}\right)$, the function $t \mapsto \min \left(\rho-\tau,(q-1)^{-1 /(q-1)}-\right.$ $\tau)(t+\sigma)^{-1 /(q-1)}$ is a subsolution of (1.1) which is smaller than $u$ at $t=0$. Therefore there exists $t_{\tau}>0$, independent of $\sigma$ such that

$$
u(x, t) \geqslant \min \left(\rho-\tau,(q-1)^{-1 /(q-1)}-\tau\right)(t+\sigma)^{-1 /(q-1)}
$$

on $\Omega \times\left(t_{\tau}\right]$. Letting $\sigma \rightarrow 0$ yields to

$$
\liminf _{t \rightarrow 0} t^{1 /(q-1)} u(x, t) \geqslant \min \left(\rho-\tau,(1 /(q-1))^{1 /(q-1)}-\tau\right)
$$

which implies the claim by letting $\rho \rightarrow 0$.
Remark 4.6. If the function $u$ has zero boundary data, then the conclusion of Theorem 4.10 becomes more precise:
(i) either

$$
\liminf _{t \rightarrow 0} t^{1 /(p-2)} u(x, t)=v_{\Omega}(x), \quad \forall x \in \Omega
$$

(ii) or the initial trace of $u$ is a Radon measure $\mu \in \mathscr{M}^{+}(\Omega)$.

This conclusion also holds for Theorems 4.3, 4.6, and 4.7.

## 5. SOLUTION WITH A GIVEN INITIAL TRACE

In this section, we consider, the problem of the existence of a solution of (1.1) in $Q_{T}=\Omega \times(0, T)$ with a given initial trace

$$
\left\{\begin{array}{l}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+u^{q}=0 \quad \text { in } Q_{T}  \tag{5.1}\\
\operatorname{Tr}_{\Omega}(u)=v \in \mathscr{B}_{\mathrm{reg}}^{+}(\Omega) \approx(\mathscr{S}, \mu)
\end{array}\right.
$$

for a given couple $(\mathscr{S}, \mu) \in C M^{+}(\Omega)$. Let $\mathscr{R}=\mathbb{R}^{N} \backslash \mathscr{S}$.
Definition 5.1. A nonnegative function $u$ is a solution of (5.1) if $u$ is a weak solution of (1.1), and
(i) for any open subset $U$ such that $U \cap \mathscr{S} \neq \emptyset$,

$$
\lim _{t \rightarrow 0} \int_{U} u(x, t) d x=\infty,
$$

(ii) for any $0<\theta<T, u \in L^{\infty}\left((0, \theta) ; L_{\text {loc }}^{1}(\mathscr{R})\right),|\nabla u| \in L_{\mathrm{loc}}^{p-1}(\mathscr{R} \times[0, T))$, and $u$ satisfies (3.4) and (3.5).

There are many possibilities according to $p-1$ is smaller or larger than $q$, and $\Omega$ is bounded or equal to $\mathbb{R}^{N}$. Even in the case where the initial data is a Radon measure $\mu$, i.e. $\mathscr{S}=\emptyset$, some difficulties occur. The first one may come from the concentration of the measure whenever $1<p \leqslant p_{0}$. A detailed analysis of this problem will be made in a forthcoming paper. When $\Omega=\mathbb{R}^{N}$ and $p>2$, another difficulty comes from the growth of the measure at infinity as it was pointed out in [7,8] for Eq. (1.12). In view of all the difficulties mentioned above, we will concentrate essentially on the case $\Omega=\mathbb{R}^{N}$, and consider the general initial value problem associated to (1.1) in

$$
\begin{gather*}
Q_{\infty}=\mathbb{R}^{N} \times(0, \infty) \\
\left\{\begin{array}{l}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+u^{q}=0 \\
\operatorname{Tr}_{\mathbb{R}^{N}}(u)=v \in \mathscr{B}_{\text {reg }}^{+}\left(\mathbb{R}^{N}\right) \approx(\mathscr{S}, \mu),
\end{array} \quad \text { in } Q_{\infty},\right. \tag{5.2}
\end{gather*}
$$

for a given couple $(\mathscr{S}, \mu) \in C M^{+}\left(\mathbb{R}^{N}\right)$. We recall that $\mathscr{R}=\mathbb{R}^{N} \backslash \mathscr{S}$.
If $q \leqslant 1$, if the problem has a solution, then $\mathscr{S}=\emptyset$ and the initial trace is reduced to a Radon measure, from Theorems 3.3 and 4.1, at least if $u$ is continuous. If $1<q \leqslant p-1$, then $\mathscr{S}=\mathbb{R}^{N}$ or $\mathscr{S}=\emptyset$.

Thus the most interesting case occurs when

$$
\begin{equation*}
q>\max (1, p-1) \tag{5.3}
\end{equation*}
$$

where we can solve the problem for any Borel measure $v$. In that case, we have a pointwise estimate, which plays an important part: any nonnegative weak solution $u \in C\left(Q_{\infty}\right)$ of (1.1) in $Q_{\infty}$ satisfies

$$
\begin{equation*}
u(x, t) \leqslant W(t)=\left(\frac{1}{t(q-1)}\right)^{1 /(q-1)}, \quad \forall(x, t) \in Q_{\infty} \tag{5.4}
\end{equation*}
$$

In all the cases we have to solve first the initial value problem with a Radon measure as initial data. In the sequel, we shall distinguish according to the cases $q<q_{\mathrm{c}}$, with (5.3) called the subcritical case, and $q \geqslant q_{\mathrm{c}}$ with $q>1$ that we call the supercritical case. Notice that if $q \geqslant q_{\mathrm{c}}$ and $p>p_{0}$, then $q>\max (1, p-1)$.

### 5.1. The Subcritical Case with a Radon Measure

In the case of Eq. (1.12), with no absorption term, it was shown in [8] that a solution exists with any Radon measure as initial data, when $p_{0}<p<2$. But there are restrictions on the behavior of $\mu$ at infinity when $p>2$. Concerning (1.1), we can construct a solution when $q$ is subcritical and $p_{0}<p<2$, or $q \geqslant p-1$. In the case $p>p_{1}$ our proof relies on some arguments of [3], which were also used in [1, 17] for elliptic problems. We combine this with the regularity estimates for a power of $u$. This allows to reach the range $p>p_{0}$. Contrary to the proof of [8], we do not use the local boundedness of the solution, and actually, our result is more general, see Remark 5.1.

Theorem 5.1. Let $\mu \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$. Assume that

$$
\begin{gather*}
p>p_{0}\left(\text { or } p>1, \text { if } \mu \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N}\right)\right),  \tag{5.5}\\
p-1<q, \text { or } p<2 \tag{5.6}
\end{gather*}
$$

$$
\begin{equation*}
q<q_{\mathrm{c}} . \tag{5.7}
\end{equation*}
$$

Then there exists a solution $u \in C\left((0, \infty), L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N}\right)\right)$ to

$$
\left\{\begin{array}{l}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+u^{q}=0 \quad \text { in } Q_{\infty}  \tag{5.8}\\
u(., 0)=\mu \quad \text { in } \mathbb{R}^{N}
\end{array}\right.
$$

Proof. Let $\mu_{n} \in \mathscr{D}\left(\mathbb{R}^{N}\right)$ be nonnegative and converging to $\mu$ in the weak sense, for example $\mu_{n}=\left(\mu \chi_{B_{n-1}}\right) * \rho_{n}$, where $\left(\rho_{n}\right)$ is a regularizing sequence with support in $B_{1 / n}$; hence supp $\mu_{n} \subset B_{n}$. Then

$$
\begin{equation*}
\int_{B_{\rho}} \mu_{n}(x) d x \leqslant \int_{B_{\rho+1}} d \mu(x), \quad \forall \rho>0 . \tag{5.9}
\end{equation*}
$$

We consider the approximate problem

$$
\left\{\begin{array}{l}
\partial_{t} u_{n}-\nabla \cdot\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right)+u_{n}^{q}=0 \quad \text { in } Q_{\infty}  \tag{5.10}\\
u_{n}(0)=\mu_{n} \quad \text { in } \mathbb{R}^{N}
\end{array}\right.
$$

It admits a (unique) weak solution, such that $u_{n} \in L_{\mathrm{loc}}^{p}\left([0, \infty), W_{\mathrm{loc}}^{1, p}\left(\mathbb{R}^{N}\right)\right)$. And $u_{n}$ is a strong solution, namely $u_{n} \in C\left([0, \infty), L_{\mathrm{loc}}^{2}\left(\mathbb{R}^{N}\right)\right), u_{n}^{q} \in$ $C\left([0, \infty), L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N}\right)\right)$, and $\left|\nabla u_{n}\right| \in C\left([0, \infty), L_{\mathrm{loc}}^{p}\left(\mathbb{R}^{N}\right)\right)$ and $\partial_{t} u_{n} \in L^{2}\left(Q_{\infty}\right)$.

Step 1 (A Priori Estimates). Let $\rho, T>0$ be fixed, and

$$
Q_{\rho, T}=(0, T) \times B_{\rho}
$$

Let $\zeta \in \mathscr{D}\left(B_{2 \rho}\right)$ with values in $[0,1]$ such that $\zeta=1$ on $B_{\rho}$, and $\tau>0$ large enough. Set $\theta>0$. Applying (2.9) to $u_{n}$, and letting $t \rightarrow 0$ infers

$$
\begin{aligned}
& \frac{1}{4} \int_{\Omega} u_{n}(x, \theta) \zeta^{\tau}(x) d x+\frac{1}{2} \int_{0}^{\theta} \int_{\Omega} u_{n}^{q} \zeta^{\tau} d x d t \\
& \quad \leqslant C+\int_{\Omega} \mu_{n}(x) \zeta^{\tau}(x) d x+\tau \int_{0}^{\theta} \int_{\Omega} \zeta^{\tau-1}\left|\nabla u_{n}\right|^{p-1}|\nabla \zeta| d x d t
\end{aligned}
$$

First, we suppose that $q>p-1$. Then it implies estimate (3.11) and we get

$$
\begin{equation*}
\frac{1}{8} \int_{\Omega} u_{n}(x, \theta) \zeta^{\tau}(x) d x+\int_{0}^{\theta} \int_{\Omega} u_{n}^{q} \zeta^{\tau} d x d t \leqslant \int_{\Omega} \mu_{n}(x) \zeta^{\tau}(x) d x+C \tag{5.11}
\end{equation*}
$$

where $C=C(p, q, N, T, \zeta, \tau)$ is independent on $u_{n}$. Next we suppose that $q \leqslant p-1<1$. By (2.7), (2.8) and Young's inequality, for any $\alpha>0$,

$$
\begin{aligned}
& \tau \int_{0}^{\theta} \int_{\Omega} \zeta^{\tau-1}\left|\nabla u_{n}\right|^{p-1}|\nabla \zeta| d x d t \\
& \quad \leqslant \varepsilon \frac{|\alpha|}{2} \int_{0}^{\theta} \int_{\Omega}\left(1+u_{n}\right)^{\alpha-1}\left|\nabla u_{n}\right|^{p} \zeta^{\tau} d x d t+\varepsilon \int_{t}^{\theta} \int_{\Omega} u_{n} \zeta^{\tau} d x d t+C_{\varepsilon} \\
& \quad \leqslant 2 \varepsilon \int_{\Omega}\left(u_{n}(x, \theta)\right) \zeta^{\tau}(x) d x+\varepsilon \int_{t}^{\theta} \int_{\Omega} u_{n}^{q} \zeta^{\tau} d x d t+\varepsilon \int_{t}^{\theta} \int_{\Omega} u_{n} \zeta^{\tau} d x d t+C_{\varepsilon} .
\end{aligned}
$$

Hence

$$
\begin{align*}
& \frac{1}{8} \int_{\Omega} u_{n}(x, \theta) \zeta^{\tau}(x) d x+\frac{1}{4} \int_{0}^{\theta} \int_{\Omega} u_{n}^{q} \zeta^{\tau} d x d t \\
& \quad \leqslant C+\int_{\Omega} \mu_{n}(x) \zeta^{\tau}(x) d x+\varepsilon \int_{0}^{\theta} \int_{\Omega} u_{n} \zeta^{\tau} d x d t+C_{\varepsilon} \tag{5.12}
\end{align*}
$$

Integrating on $(0, T)$, and choosing $\varepsilon=\frac{1}{8} T$, we get

$$
\begin{equation*}
\int_{0}^{\theta} \int_{\Omega} u_{n} \zeta^{\tau} d x d t \leqslant C+C \int_{\Omega} \mu_{n}(x) \zeta^{\tau}(x) d x \tag{5.13}
\end{equation*}
$$

and then again (5.11) holds from (5.12) and (5.13). In both cases, $\left(u_{n}\right)_{n>\rho}$ is bounded in $L^{\infty}\left((0, T), L^{1}\left(B_{\rho}\right)\right), \quad\left(u_{n}^{q}\right)_{n>\rho}$ is bounded in $L^{1}\left(Q_{\rho, T}\right)$.

Besides this,

$$
\left(u_{n}^{p-1}\right)_{n>\rho} \text { is bounded in } L^{1}\left(Q_{\rho, T}\right)
$$

by the Hölder inequality, whenever $q>p-1$, or because $p-1<1$. Then, by Proposition 2.2,

$$
\begin{gather*}
\left(u_{n}^{\sigma}\right)_{n>\rho} \text { is bounded in } L^{1}\left(Q_{\rho, T}\right), \quad \forall \sigma \in\left(0, q_{\mathrm{c}}\right),  \tag{5.14}\\
\left(\left|\nabla u_{n}\right|^{r}\right)_{n>\rho} \text { is bounded in } L^{1}\left(Q_{\rho, T}\right), \quad \forall r \in\left(0, N q_{\mathrm{c}} /(N+1)\right) . \tag{5.15}
\end{gather*}
$$

Let $\alpha \in(\max (1-p,-1), 0)$ and

$$
\begin{equation*}
v_{n}(x, t)=\left(1+u_{n}(x, t)\right)^{\beta}, \quad \beta=(\alpha+p-1) / p \tag{5.16}
\end{equation*}
$$

Clearly,

$$
\left(v_{n}^{p+p / N \beta}\right)_{n>\rho} \text { is bounded in } L^{1}\left(Q_{\rho, T}\right),
$$

and it follows from (2.26) that

$$
\left(\left|\nabla v_{n}\right|^{p}\right)_{n>\rho} \text { is bounded in } L^{1}\left(Q_{\rho, T}\right)
$$

Consequently, $v_{n}$ is bounded in $L_{\mathrm{loc}}^{p}\left([0, \infty), W^{1, p}\left(B_{\rho}\right)\right)$. Next, we estimate the $t$-derivative of $v_{n}$ :

$$
\begin{aligned}
\partial_{t} v_{n} & =\beta\left(1+u_{n}\right)^{\beta-1} \partial_{t} u_{n}=\beta\left(1+u_{n}\right)^{\beta-1}\left(-u_{n}^{q}+\nabla \cdot\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right)\right. \\
& =\beta\left(H_{n}+\nabla \cdot K_{n}\right)
\end{aligned}
$$

where

$$
\begin{aligned}
& H_{n}=-\left(1+u_{n}\right)^{\beta-1} u_{n}^{q}-(\beta-1)\left(1+u_{n}\right)^{\beta-2}\left|\nabla u_{n}\right|^{p}, \\
& K_{n}=\left(1+u_{n}\right)^{\beta-1}\left|\nabla u_{n}\right|^{p-2} \nabla u_{n} .
\end{aligned}
$$

Because $\beta<1, H_{n}$ is bounded in $L^{1}\left(B_{\rho}\right)$, by (2.26). And $\left|K_{n}\right| \leqslant\left|\nabla u_{n}\right|^{p-1}$. Then from the estimate of $\left|\nabla u_{n}\right|^{r}, K_{n}$ is bounded in $L^{s}\left(Q_{\rho, T}\right)$ for any $s \in\left[1, s_{\mathrm{c}}\right)$, where

$$
\begin{equation*}
s_{\mathrm{c}}=1+1 /(N+1)(p-1) \tag{5.17}
\end{equation*}
$$

In particular $\nabla \cdot K_{n}$ is bounded in $L^{s}\left((0, T), W^{-1, s}(\Omega)\right)$. It implies that $\left(\partial_{t} v_{n}\right)_{n>\rho}$ is bounded in $L^{1}\left(B_{\rho}\right)+L^{s}\left((0, T), W^{-1, s}(\Omega)\right)$ for any $s \in\left[1, s_{\mathrm{c}}\right)$.

Step 2 (Convergence). For any $\rho>0$, and any $T>0$, the sequence $\left(v_{n}\right)_{n>\rho}$ is relatively compact in $L^{1}\left(Q_{\rho, T}\right)$. We choosing $\rho, T$ as integers. There exists a subsequence, still denoted $\left(v_{n}\right)$, such that $\left(v_{n}\right)$ converges a.e. in $Q_{\infty}$ toward
some locally integrable function $v$. Then $\left(u_{n}\right)$ converges a.e. in $Q_{\infty}$ to $u=v^{1 / \beta}-1$. In particular,

$$
\begin{equation*}
\left(u_{n}\right) \text { converges locally in measure to } u \text {. } \tag{5.18}
\end{equation*}
$$

Following [3], we deduce that

$$
\begin{equation*}
\left(\nabla u_{n}\right) \text { is locally a Cauchy sequence in measure. } \tag{5.19}
\end{equation*}
$$

Indeed let $\lambda, \varepsilon>0$ and set

$$
E_{n, m, \lambda}=\left\{\left|\nabla\left(u_{n}-u_{m}\right)\right|>\lambda\right\}=\left\{(x, t) \in Q_{\rho, T}:\left|\nabla u_{n}-\nabla u_{m}\right|>\lambda\right\}
$$

We have to prove that $\left|E_{n, m, \lambda}\right| \leqslant \varepsilon$ for $m, n$ large enough. We write $E_{n, m, \lambda} \subset$ $E_{n, m, \lambda}^{1} \cup E_{n, m, \lambda}^{2} \cup E_{n, m, \lambda}^{3}$ where

$$
\begin{aligned}
E_{n, m, \lambda}^{1}= & \left\{(x, t) \in Q_{\rho, T}:\left|\nabla u_{n}\right|>b\right\} \cup\left\{\left|\nabla u_{m}\right|>b\right\} \cup\left\{u_{n}>b\right\} \cup\left\{u_{m}>b\right\} \\
E_{n, m, \lambda}^{2}= & \left\{(x, t) \in Q_{\rho, T}:\left|u_{n}-u_{m}\right|>k\right\} \\
E_{n, m, \lambda}^{3}= & \left\{(x, t) \in Q_{\rho, T}:\left|\nabla u_{n}\right| \leqslant b,\left|\nabla u_{m}\right| \leqslant b, u_{n} \leqslant b, u_{m} \leqslant b,\right. \\
& \left.\left|u_{n}-u_{m}\right| \leqslant k,\left|\nabla\left(u_{n}-u_{m}\right)\right|>\lambda\right\}
\end{aligned}
$$

and $k, b>0$ are parameters. From the estimates on $u_{n}$ and $\nabla u_{n}$, we can choose $b=b_{\varepsilon}$ such that $\left|E_{n, m, \lambda}^{1}\right| \leqslant \varepsilon / 3$ for any $m, n \in \mathbb{N}$. Next, we have

$$
\begin{equation*}
\partial_{t}\left(u_{n}-u_{m}\right)-\nabla \cdot\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m}\right)+u_{n}^{q}-u_{m}^{q}=0 \tag{5.20}
\end{equation*}
$$

Set $\tau>1$ and $\Theta_{k}(s)=\int_{0}^{s} T_{k}(\theta) d \theta$, then $\left|\Theta_{k}(s)\right| \leqslant k|s|$. Multiplying (5.20) by $T_{k}\left(u_{n}-u_{m}\right) \zeta^{\tau}$, we get

$$
\begin{aligned}
& \int_{B_{2 \rho}} \Theta_{k}\left(u_{n}-u_{m}\right)(T) \zeta^{\tau} d x+\int_{0}^{T} \int_{B_{2 \rho}}\left(u_{n}^{q}-u_{m}^{q}\right) T_{k}\left(u_{n}-u_{m}\right) \zeta^{\tau} d x d t \\
& \quad+\int_{0}^{T} \int_{B_{2 \rho} \cap\left\{\left|u_{n}-u_{m}\right| \leqslant k\right\}}\left(\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m}\right)\left(\nabla u_{n}-\nabla u_{m}\right)\right) \zeta^{\tau} d x d t \\
& =\int_{B_{2 \rho}} \Theta_{k}\left(\mu_{n}-\mu_{m}\right) \zeta^{\tau} d x \\
& \quad+\int_{0}^{T} \int_{B_{2 \rho}}\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m}\right) T_{k}\left(u_{n}-u_{m}\right) \zeta^{\tau-1} \nabla \zeta d x d t
\end{aligned}
$$

Next from (5.9),

$$
\begin{equation*}
\int_{B_{2 \rho}} \Theta_{k}\left(\mu_{n}-\mu_{m}\right) \zeta^{\theta} d x \leqslant k\left(\int_{B_{2 \rho}}\left(\mu_{n}+\mu_{m}\right) d x\right) \leqslant 2 k C \int_{B_{2 \rho+1}} \mu d x=k C_{\rho} \tag{5.21}
\end{equation*}
$$

with $C_{\rho}=2 C \int_{B_{2 \rho}} \mu d x$, and from (5.15),

$$
\begin{aligned}
& \int_{0}^{T} \int_{B_{2 \rho}}\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m}\right) T_{k}\left(u_{n}-u_{m}\right) \zeta^{\tau-1} \nabla \zeta d x d t \\
& \quad \leqslant k \int_{0}^{T} \int_{B_{2 \rho}}\left(\left|\nabla u_{n}\right|^{p-1}+\left|\nabla u_{m}\right|^{p-1}\right) \zeta^{\tau-1}|\nabla \zeta| d x d t \leqslant k C_{\rho},
\end{aligned}
$$

with another $C_{\rho}>0$. Hence

$$
\int_{0}^{T} \int_{B_{\rho} \cap\left\{\left|u_{n}-u_{m}\right| \leqslant k\right\}}\left(\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m}\right)\left(\nabla u_{n}-\nabla u_{m}\right)\right) d x d t \leqslant k C_{\rho}
$$

When $p<2$, this implies

$$
\int_{0}^{T} \int_{B_{\rho} \cap\left|u_{n}-u_{m}\right| \leqslant k} \frac{\left|\nabla u_{n}-\nabla u_{m}\right|^{2}}{\left|\nabla u_{n}\right|^{2-p}+\left|\nabla u_{m}\right|^{2-p}} d x d t \leqslant k C_{\rho}
$$

hence $\left|E_{n, m, \lambda}^{3}\right| \leqslant 2 b_{\varepsilon}^{2-p} k C_{\rho} / \lambda^{2} \leqslant \varepsilon / 3$ as soon as $k \leqslant k_{\varepsilon}$ small enough. When $p \geqslant 2$, we have

$$
\int_{0}^{T} \int_{B_{\rho} \cap\left|u_{n}-u_{m}\right| \leqslant k}\left|\nabla u_{n}-\nabla u_{m}\right|^{p} d x d t \leqslant k C_{\rho}
$$

hence $\left|E_{n, m, \lambda}^{3}\right| \leqslant k C_{\rho} / \lambda^{p} \leqslant \varepsilon / 3$ as soon as $k \leqslant k_{\varepsilon}$ small enough. After having chosen such a $k$, we deduce that there exists $n_{\varepsilon} \in \mathbb{N}$ such that $\left|E_{n, m, \lambda}^{2}\right| \leqslant \varepsilon / 3$ for any $m, n \geqslant n_{\varepsilon}$, from (5.18). Hence $\left|E_{n, m, \lambda}\right| \leqslant \varepsilon$ for any $m, n \geqslant n_{\varepsilon}$, and (5.19) follows. This also implies that

$$
\begin{equation*}
\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right) \text { is locally a Cauchy sequence in measure. } \tag{5.22}
\end{equation*}
$$

But from (5.15),

$$
\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right)_{n>\rho} \text { is bounded in }\left(L^{s}\left(Q_{\rho, T}\right)\right)^{N}, \quad \forall s \in\left[1, s_{\mathrm{c}}\right)
$$

After extraction of another subsequence, there exists some $y=\left(y_{1}, \ldots, y_{N}\right)$, $w=\left(w_{1}, \ldots, w_{N}\right)$ such that

$$
\left(\nabla u_{n}\right) \rightarrow y \quad \text { a.e. in } Q_{\infty}, \quad\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right) \rightarrow w=|y|^{p-2} y \quad \text { a.e. in } Q_{\infty}
$$

and

$$
\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right) \rightarrow w \text { strongly in }\left(L^{s}\left(Q_{\rho, T}\right)\right)^{N}, \quad \forall s \in\left[1, s_{\mathrm{c}}\right) .
$$

Let us denote

$$
u_{n}^{k}=T_{k}\left(u_{n}\right) \quad \text { and } \quad u^{k}=T_{k}(u)
$$

Now for any fixed $k>0, u_{n}^{k}$ converges a.e. to $u^{k}$. Moreover, it converges weakly in the space $L^{p}\left((0, T), W^{1, p}\left(B_{\rho}\right)\right)$, and the limit function is $u^{k}$. In particular,

$$
\nabla u_{n}^{k} \rightarrow \nabla u^{k} \quad \text { weakly in }\left(L^{p}\left(Q_{\rho, T}\right)\right)^{N}
$$

But $\nabla u_{n}^{k}=\nabla u_{n} \times 1_{\left\{\left|u_{n}\right| \leqslant k\right\}}$, hence $\nabla u_{n}^{k}$ converges a.e. to $y \times 1_{\{|u| \leqslant k\}}$. Then $\nabla u^{k}=y \times 1_{\{|u| \leqslant k\}}$ a.e. in $Q_{\infty}$. Since we have defined the gradient of $u$ by $\nabla u=y$, it follows that

$$
\begin{equation*}
\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right) \rightarrow|\nabla u|^{p-2} \nabla u \quad \text { strongly in }\left(L^{s}\left(Q_{\rho, T}\right)\right)^{N}, \quad \forall s \in\left[1, s_{\mathrm{c}}\right) . \tag{5.23}
\end{equation*}
$$

Either $p>p_{0}$, hence $1<q_{\mathrm{c}}$, so that (5.14) implies that

$$
\begin{equation*}
\left(u_{n}\right) \rightarrow u \quad \text { strongly in } L^{\sigma}\left(Q_{\rho, T}\right), \quad \forall \sigma \in\left[1, q_{\mathrm{c}}\right) \tag{5.24}
\end{equation*}
$$

Or $\mu \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N}\right)$. Moreover $\mu_{n}$ converges strongly to $\mu$ in $L^{1}\left(B_{\rho}\right)$ for any $\rho>0$. Hence for any $\varepsilon>0$, and if $m, n \geqslant n(\varepsilon)$,

$$
\left.\frac{1}{k} \int_{B_{2 \rho}} \Theta_{k}\left(\mu_{n}-\mu_{m}\right) \zeta^{\theta} d x \leqslant \int_{B_{2 \rho}}\left|\mu_{n}-\mu_{m}\right|\right) d x \leqslant \varepsilon
$$

Since $\left|\Theta_{k}(s)\right| \geqslant k|s| / 2$ on the set $\{|s| \geqslant k\}$, we have for any $T>0$,

$$
\begin{aligned}
& \frac{1}{2} \int_{B_{\rho} \cap\left\{\left|u_{n}(., T)-u_{m}(., T)\right| \geqslant k\right\}}\left|u_{n}(x, T)-u_{m}(x, T)\right| d x \\
& \quad \leqslant \frac{1}{k} \int_{B_{2 \rho}} \Theta_{k}\left(u_{n}-u_{m}\right)(T) \zeta^{\tau} d x
\end{aligned}
$$

Hence

$$
\begin{align*}
& \frac{1}{2} \int_{B_{\rho} \cap\left\{\left|u_{n}(\cdot, T)-u_{m}(., T)\right| \geqslant k\right\}}\left|u_{n}(x, T)-u_{m}(x, T)\right| d x \\
& \quad+\frac{1}{k} \int_{0}^{T} \int_{B_{\rho} \cap\left\{\left|u_{n}-u_{m}\right| \leqslant k\right\}}\left(\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m}\right)\left(\nabla u_{n}-\nabla u_{m}\right)\right) d x d t \\
& \quad \leqslant \varepsilon+\frac{1}{k} \int_{0}^{T} \int_{B_{2 \rho}}\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m}\right) T_{k}\left(u_{n}-u_{m}\right) \zeta^{\tau-1} \nabla \zeta d x d t \\
& \quad \leqslant \varepsilon+C_{\rho} \int_{0}^{T} \int_{B_{2 \rho}} \|\left.\nabla u_{n}\right|^{p-2} \nabla u_{n}-\left|\nabla u_{m}\right|^{p-2} \nabla u_{m} \mid d x d t \leqslant 2 \varepsilon \tag{5.25}
\end{align*}
$$

for $m, n \geqslant n^{\prime}(\varepsilon)$ large enough, independent on $k$. Letting $k \rightarrow 0$, this proves that

$$
\begin{equation*}
\left(u_{n}\right) \rightarrow u \text { strongly in } C\left([0, T], L^{1}\left(B_{\rho}\right)\right) \tag{5.26}
\end{equation*}
$$

In any case, from (5.24) or (5.26),

$$
\begin{equation*}
\left(u_{n}\right) \rightarrow u \text { strongly in } L^{1}\left(Q_{\rho, T}\right) \tag{5.27}
\end{equation*}
$$

Next, we show that the limit function admits $\mu$ as an initial trace, since $q<q_{\mathrm{c}}$. We have the property

$$
\begin{equation*}
\left(u_{n}^{q}\right) \rightarrow u^{q} \text { strongly in } L^{\tau}\left(Q_{\rho, T}\right), \quad \forall \tau \in\left[1, q_{\mathrm{c}} / q\right) \tag{5.28}
\end{equation*}
$$

For any $\xi \in C_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{N}\right)$ and $t>0$, we have

$$
\begin{align*}
& \int_{\mathbb{R}^{N}} u_{n}(x, t) \xi(x) d x-\int_{\mathbb{R}^{N}} \mu_{n}(x) \xi(x) d x \\
& \quad=\int_{0}^{t} \int_{\Omega}\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n} \nabla \xi+u_{n}^{q} \xi\right) d x d t \tag{5.29}
\end{align*}
$$

Up to the extraction of a subsequence, we can pass to the limit in each term, for almost any $t>0$, and get

$$
\int_{\mathbb{R}^{N}} u(x, t) \xi(x) d x-\int_{\mathbb{R}^{N}} \xi(x) d \mu(x)=\int_{0}^{t} \int_{\Omega}\left(|\nabla u|^{p-2} \nabla u \nabla \xi+u^{q} \xi\right) d x d t
$$

from (5.23), (5.28). Since the right-hand side tends to 0 as $t$ goes to 0 , hence

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u(x, t) \xi(x) d x=\int_{\mathbb{R}^{N}} \xi(x) d \mu(x) . \tag{5.30}
\end{equation*}
$$

For any $\varphi \in C_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{N} \times[0, \infty)\right)$ and $\theta>0$, we have

$$
\begin{aligned}
& \int_{0}^{\theta} \int_{\mathbb{R}^{N}}\left(-u_{n} \varphi_{t}+\left|\nabla u_{n}\right|^{p-2} \nabla u_{n} \nabla \varphi+u_{n}^{q} \varphi\right) d x d t \\
& \quad=\int_{\mathbb{R}^{N}} \varphi(0, x) \mu_{n}(x) d x-\int_{\mathbb{R}^{N}} u_{n}(x, \theta) \varphi(x, \theta) d x
\end{aligned}
$$

for any $n>\rho$, by using also (5.27). Then $u$ satisfies (3.5) in $\mathscr{R}=\mathbb{R}^{N}$.
It remains to prove that $u$ is a weak solution. First suppose that $\mu \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N}\right)$. Going to the limit in (5.25) when $m \rightarrow \infty$ from the Fatou Lemma and Lebesgue theorem, we get for any $T>0$

$$
\begin{aligned}
& \frac{1}{k} \int_{0}^{T} \int_{B_{\rho} \cap\left\{\left|u_{n}-u\right| \leqslant k\right\}}\left(\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-|\nabla u|^{p-2} \nabla u\right)\left(\nabla u_{n}-\nabla u\right)\right) d x d t \\
& \quad \leqslant \varepsilon+\left.C_{\rho} \int_{0}^{T} \int_{B_{2 \rho}}| | \nabla u_{n}\right|^{p-2} \nabla u_{n}-|\nabla u|^{p-2} \nabla u \mid d x d t \leqslant 2 \varepsilon
\end{aligned}
$$

Moreover,

$$
\begin{aligned}
X_{k, n} & =\int_{0}^{T} \int_{B_{\rho}}\left(\left(\left|\nabla u_{n}^{k}\right|^{p-2} \nabla\left(u_{n}^{k}\right)-\left|\nabla u^{k}\right|^{p-2} \nabla u^{k}\left(\nabla u_{n}^{k}-\nabla u^{k}\right) d x d t\right.\right. \\
& \leqslant \int_{0}^{T} \int_{B_{\rho} \cap\left\{\left|u_{n}-u\right| \leqslant 2 k\right\}}\left(\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}-|\nabla u|^{p-2} \nabla u\right)\left(\nabla u_{n}-\nabla u\right)\right) d x d t
\end{aligned}
$$

Hence for fixed $k, X_{k, n} \rightarrow 0$. Therefore,

$$
\int_{0}^{T} \int_{B_{\rho}}\left|\nabla u_{n}^{k}\right|^{p} d x d t \rightarrow \int_{0}^{T} \int_{B_{\rho}}\left|\nabla u^{k}\right| \cdot{ }^{p} d x d t
$$

which implies

$$
\nabla u_{n}^{k} \rightarrow \nabla u^{k} \text { strongly in } L^{p}\left(Q_{\rho, T}\right)
$$

Next, we consider the general case where $\mu$ is a measure. Then for almost all $\tau>0$ and $\rho>0 u_{n}(., \tau)$ converges strongly to $u(., \tau)$ in $L^{1}\left(B_{\rho}\right)$. Let $0<\tau<T$; if we replace the interval $[0, T]$ by $[\tau, T]$, we deduce as above that $u \in C\left((0, \infty), L^{1}\left(B_{\rho}\right)\right),\left(u_{n}\right) \rightarrow u$ strongly in $C\left([\tau, T], L^{1}\left(B_{\rho}\right)\right)$, and

$$
\begin{equation*}
\nabla u_{n}^{k} \rightarrow \nabla u^{k} \text { strongly in } L^{p}\left((\tau, T) \times B_{\rho}\right) \tag{5.31}
\end{equation*}
$$

In fact this happens for any $0<\tau<T$. Let $\varphi \in C_{\mathrm{c}}^{\infty}\left(Q_{\infty}\right)$. Consider $0<\tau<T$ and $\rho>0$ such that the support of $\varphi$ is contained in $(\tau, T) \times B_{\rho}$. Let $h \in C^{1}(\mathbb{R}) \cap W^{1, \infty}(\mathbb{R})$ where $H^{\prime}(r)=h(r)$, and $h(r)$ constant outside $[-k, k]$
for some $k>0$. Multiplying Eq. (5.10) by $h\left(u_{n}\right) \varphi$, we get

$$
\begin{align*}
0= & \int_{0}^{\infty} \int_{\mathbb{R}^{v}}\left(-H\left(u_{n}\right) \partial_{t} \varphi+\left|\nabla u_{n}\right|^{p-2} \nabla\left(h\left(u_{n}\right) \varphi\right)+h\left(u_{n}\right) u_{n}^{q} \varphi\right) d x d t \\
= & \int_{0}^{\infty} \int_{\mathbb{R}^{v}}\left(-H\left(u_{n}\right) \partial_{t} \varphi+\left|\nabla u_{n}^{k}\right|^{p} h^{\prime}\left(u_{n}\right) \varphi+h\left(u_{n}\right)\left|\nabla u_{n}\right|^{p-2} \nabla u_{n} \cdot \nabla u \varphi\right. \\
& \left.+h\left(u_{n}\right) u_{n}^{q} \varphi\right) d x d t . \tag{5.32}
\end{align*}
$$

Owing to (5.28), (5.27) and (5.31), we can pass to the limit in each term. Thus the proof is complete.

Remark 5.1. In the case $q>\max (1, p-1)$, the proof is much shorter, since (5.4) implies that $\left\{u_{n}\right\}$ is locally bounded, uniformly with respect to $n$; therefore the use of the $T_{k}$-truncature of $u_{n}$ is useless. Moreover, because

$$
\partial_{t} u_{n}-\nabla \cdot\left(\left|\nabla u_{n}\right|^{p-2} \nabla u_{n}\right)+u_{n}^{q}=0
$$

holds in $Q_{\infty}$ and $\left\{u_{n}^{q}\right\}$ is locally bounded uniformly with respect to $n$, the sequence $\left\{u_{n}\right\}$ is equicontinuous in the local uniform topology of $Q_{\infty}$, hence, up to a subsequence, it converges uniformly on any compact subset of $Q_{\infty}$. Thus in particular $u$ is continuous on $Q_{\infty}$. Actually, for a given open bounded domain $\Omega$ in $\mathbb{R}^{N}$, the regularity results of $[5,9]$, assert that, any $\left.v \in C\left((0, T), L^{2}(\Omega)\right)\right) \cap L^{2}\left((0, T), W^{1, p}(\Omega)\right) \cap L^{\infty}\left(Q_{T}\right)$ solution of

$$
\begin{equation*}
\partial_{t} v-\nabla \cdot\left(|\nabla v|^{p-2} \nabla v\right)=h \tag{5.33}
\end{equation*}
$$

in $Q_{T}=\Omega \times(0, T)$, with $h \in L^{\infty}\left(Q_{T}\right)$, is Hölder continuous, and the same is true with $\nabla v$ : there exists $\alpha \in(0,1)$ such that for any compact set $K \subset \Omega$, and any $\left(x_{i}, t_{i}\right) \in K \times[T / 2,2 T / 3]$.

$$
\left|v\left(x_{1}, t_{1}\right)-v\left(x_{2}, t_{2}\right)\right| \leqslant \gamma\left(\left|x_{1}-x_{2}\right|+\left|t_{1}-t_{2}\right|^{1 / p}\right)^{\alpha},
$$

where $\gamma=\gamma\left(N, p,\|h\|_{L^{\infty}(\Omega \times(0, T))},\|v\|_{L^{\infty}(G \times(0, T))}, \operatorname{dist}(K, \partial \Omega)\right)$.
Remark 5.2. If we replace $u^{q}$ by $|u|^{q-1} u$ in (5.8) the existence result of changing sign solution with a given signed Radon measure $\mu$ as initial data holds under some minor modifications in the proof. Also, since the proof does not use local boundedness of the solutions, we can treat where there is a forcing term $f$ in the equation. Consider the problem

$$
\left\{\begin{array}{l}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+|u|^{q-1} u=f \quad \text { in } \mathbb{R}^{N} \times(0, T),  \tag{5.34}\\
u_{n}(0)=\mu \quad \text { in } \mathbb{R}^{N},
\end{array}\right.
$$

where $f$ and $\mu$ are given Radon measures in $\mathbb{R}^{N} \times(0, T)$ and $\mathbb{R}^{N}$, respectively. Under the assumptions of Theorem 5.1 on $p$ and $q$, there exists a solution of problem (5.34) in the sense of distributions in $\mathbb{R}^{N} \times(0, T)$, and such that (5.30) holds for any $\xi \in C_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{N}\right)$. If moreover $f \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N} \times(0, T)\right)$, then $u$ is a weak solution as before, that means for any $\varphi \in C_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{N} \times(0, T)\right)$,

$$
\int_{0}^{T} \int_{\mathbb{R}^{N}}\left(-H(u) \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u \cdot \nabla(h(u) \varphi)+h(u)\left(|u|^{q-1} u-f\right) \varphi\right) d x d t=0 .
$$

In the same way, we can prove the existence of a solution of the CauchyDirichlet problem

$$
\left\{\begin{array}{l}
\partial_{t} u-\nabla \cdot\left(|\nabla u|^{p-2} \nabla u\right)+|u|^{q-1} u=f \quad \text { in } \Omega \times(0, T)  \tag{5.35}\\
u=0 \quad \text { on } \partial \Omega \times(0, T), \\
u(0)=\mu \quad \text { in } \Omega
\end{array}\right.
$$

in any bounded regular domain $\Omega$, for any bounded Radon measures $f$ and $\mu$ in $\Omega \times(0, T)$ and $\Omega$, respectively. The Dirichlet condition on $\partial \Omega \times(0, \infty)$ is given in the sense $T_{k}(u) \in L_{\text {loc }}^{p}\left((0, \infty), W_{0}^{1, p}(\Omega)\right.$. In case $f \in L^{1}(\Omega \times(0, T))$, $\mu \in L^{1}(\Omega)$, one finds again the result of [15].

### 5.2. Constructive Solutions

Up to now, no uniqueness result is known for a measure, except when it is a (sum of) Dirac measure(s). That is the reason why we define a notion of solutions corresponding to the previous construction which will permit some operations and comparison between them.

Definition 5.2. Let $p>1, q>0$, and $\mu \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$. Suppose that there exists a weak nonnegative solution $u$ of (1.1) in $Q_{\infty}$ with initial trace $\mu$. Then $u$ is called a constructive solution if there exists a sequence $\left\{\mu_{n}\right\}$ of continuous functions with compact support converging weakly to $\mu$, such that the corresponding sequence $\left\{u_{n}\right\}$ of solutions to (5.8) with initial data $\mu_{n}$ converges a.e. to $u$. We shall denote by $\mathscr{M}^{*+}\left(\mathbb{R}^{N}\right)$ the set of initial traces of all the constructive solutions.

Remark 5.3. Notice that the solutions of (5.8) constructed in Theorem 5.1 are constructive, hence in that case $\mathscr{M}^{*+}\left(\mathbb{R}^{N}\right)=\mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$. Now assume (5.5) and $q \geqslant q_{\mathrm{c}}$ (hence $q>p-1$ ), then all the estimates of Theorem 5.1 are still valid, and all the convergences up to (5.27), with the noticeable exception of (5.28). When $q>1$, the limit function $u$ is a solution of (1.1), by (5.4) and Remark 5.1. The function $u$ does admits an initial trace $\mu^{\prime}$, which satisfies $\mu^{\prime} \leqslant \mu$ instead of equality. Actually, $u$ is a weak solution of (1.1) with
initial trace $\mu$ if and only if for any $\varphi \in C_{\mathrm{c}}\left(\mathbb{R}^{N} \times[0, \infty)\right)$,

$$
\begin{equation*}
\int_{0}^{T} \int_{\mathbb{R}^{N}} u_{n}^{q} \varphi d x d t \rightarrow \int_{0}^{T} \int_{\mathbb{R}^{N}} u^{q} \varphi d x d t \tag{5.36}
\end{equation*}
$$

or equivalently, if and only if

$$
\begin{equation*}
u_{n}^{q} \rightarrow u^{q} \text { strongly in } L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N} \times[0, \infty)\right) \tag{5.37}
\end{equation*}
$$

since $u_{n}^{q} \rightarrow u^{q}$ a.e. in $Q_{\infty}$. Indeed if (5.36) holds, then for any $\varphi \in C_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{N} \times[0, \infty)\right)$ we can pass to the limit in (5.29) and get (5.30), and in (5.32). Then $u$ is a weak solution of (1.1) with initial trace $\mu$. Conversely, if $u$ is such a solution, then (5.29) and (5.30) hold. Therefore (5.36) holds by difference for any $\varphi \in C_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{N} \times[0, \infty)\right.$ ), and then, by density, for any $\varphi \in C_{\mathrm{c}}\left(\mathbb{R}^{N} \times[0, \infty)\right)$. As a consequence a function $u$ is constructive if and only if there exists a sequence $\left\{\mu_{n}\right\}$ of continuous functions with compact support converging weakly to $\mu$, such that the corresponding sequence $\left\{u_{n}\right\}$ of solutions to (5.8) with initial data $\mu_{n}$ converges a.e. to $u$ and (5.36) holds.

Next we give some useful properties of the constructive solutions. They are settled upon the following measure theory result the proof of which, due to E. Lesigne, is given in the appendix.

Lemma 5.2. Let $\tilde{\mu}, \mu \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$ with $\tilde{\mu} \leqslant \mu$. Let $\mu_{n} \in C_{\mathrm{c}}\left(\mathbb{R}^{N}\right)$ be a sequence of nonnegative functions converging weakly to $\mu$. Then there exists a sequence of nonnegative functions $\tilde{\mu}_{n} \in C_{\mathrm{c}}\left(\mathbb{R}^{N}\right)$ such that

$$
\tilde{\mu}_{n}(x) \leqslant \mu_{n}(x), \quad \forall x \in \mathbb{R}^{N} \quad \text { and } \quad \forall n \in \mathbb{N}^{*}
$$

converging weakly to $\tilde{\mu}$.
Proposition 5.3. Let $\tilde{\mu}, \mu \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$ with $\tilde{\mu} \leqslant \mu$.
(i) Assume (5.5)-(5.7). Then there exist constructive solutions $\tilde{u}$, $u$ with respective initial traces $\tilde{\mu}, \mu$, such that $\tilde{u} \leqslant u$ a.e. in $Q_{\infty}$.
(ii) Assume (5.5) and $q \geqslant q_{\mathrm{c}}$, that is $q \geqslant q_{\mathrm{c}}>1$. If there exists a constructive solution $u$ with initial trace $\mu$, then there exists a constructive solution $\tilde{u}$ with initial trace $\tilde{\mu}$ such that $\tilde{u} \leqslant u$ a.e. in $Q_{\infty}$. Moreover $\mathscr{M}^{*+}\left(\mathbb{R}^{N}\right)$ is a positive cone in $\mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$.

Proof. (i) Let $\tilde{\mu}_{n}, \mu_{n}$, be defined in Lemma 5.2, and $\tilde{u}_{n} \leqslant u_{n}$ the corresponding solutions. Then $\tilde{u}_{n} \leqslant u_{n}$ in $Q_{\infty}$ from the maximum principle, hence $\tilde{u} \leqslant u$ a.e. in $Q_{\infty}$.
(ii) Let $u$ be a constructive solution with initial trace $\mu$, and $\tilde{\mu}_{n}, \mu_{n}$, and $\tilde{u}_{n} \leqslant u_{n}$ be defined as above. Then $u_{n}$ converges to $u$ a.e., in $Q_{\infty}$ and $\tilde{u}_{n}$ converges to some $\tilde{u}$ a.e. in $Q_{\infty}$. And $u_{n}^{q} \rightarrow u^{q}$, strongly in $L_{\text {loc }}^{1}\left(\mathbb{R}^{N} \times[0, \infty)\right.$ ),
hence also $\tilde{u}_{n}^{q} \rightarrow \tilde{u}^{q}$ from a variant of the Lebesgue theorem, see for example [26], so that $\tilde{u}$ is constructive.

Set $\mu \in \mathscr{M}^{*+}\left(\mathbb{R}^{N}\right), \lambda>0$ and let $u$ be a constructive solution with initial trace $\mu$. If $\lambda<1$, then $\lambda \mu \leqslant \mu$, hence from above, there exists a constructive solution $\tilde{u}$ with initial trace $\lambda \mu$ such that $\tilde{u} \leqslant u$ a.e. in $Q_{\infty}$. Next, we assume $\lambda>1$ and let $\mu_{n} \in C_{\mathrm{c}}\left(\mathbb{R}^{N}\right)$ converging weakly to $\mu$ and $u_{n}$ be the solution with initial trace $\mu_{n}$. We perform the change of scale

$$
\begin{equation*}
S_{\lambda}\left(u_{n}\right)(x, t)=\lambda u_{n}\left(x, \lambda^{p-2} t\right) \tag{5.38}
\end{equation*}
$$

Then $v_{n}=S_{\lambda}\left(u_{n}\right)$ satisfies

$$
\partial_{t} v_{n}-\nabla \cdot\left(\left|\nabla v_{n}\right|^{p-2} \nabla v_{n}\right)+v_{n}^{q}=\left(1-\lambda^{p-1-q}\right) v_{n}^{q}
$$

in $Q_{\infty}$, hence $v_{n}$ is a supersolution, with initial trace $\lambda \mu$. If $w_{n}$ denotes the solution with initial trace $\lambda \mu_{n}$, then $w_{n} \leqslant v_{n}$ from the comparison principle, and $w_{n}$ converges to some $w$ a.e. in $Q_{\infty}$. Moreover, $u_{n}^{q} \rightarrow u^{q}$, strongly in $L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{N} \times[0, \infty)\right.$ ), hence also $v_{n}^{q}=\left(S_{\lambda}\left(u_{n}\right)\right)^{q} \rightarrow\left(S_{\lambda}(u)\right)^{q}$, thus also $w_{n}^{q} \rightarrow w^{q}$. This proves that $w$ is a constructive solution with initial trace $\lambda \mu$.

Remark 5.5. Although we believe that $\mathscr{M}^{*+}\left(\mathbb{R}^{N}\right)$ is stable by addition, we have not been able to give a proof to this property.

Remark 5.6. Under the assumptions of Theorem 5.1, suppose that $\mu \in L_{\mathrm{loc}}^{1}(\Omega)$. Let $u$ be a constructive solution. Then $u \in C\left([0, T], L_{\mathrm{loc}}^{1}(\Omega)\right)$, and $u$ is in the class of existence and uniqueness introduced in [15]. Indeed for any open sets $U \subset \subset U^{*} \subset \subset \Omega$, and any $k>0$, taking $\xi \in C_{\mathrm{c}}^{\infty}(\Omega)$, with values in the interval $[0,1], 1$ in $U$, and 0 outside of $U^{*}$, and $h\left(u_{n}\right)=T_{k+1}\left(u_{n}\right)-T_{k}\left(u_{n}\right)$ in (2.4), we get for any $0<t<\theta<T$,

$$
\begin{aligned}
& \int_{t}^{\theta} \int_{U \cap\{k \leqslant u \leqslant k+1\}}\left|\nabla u_{n}\right|^{p} d x d t+\int_{t}^{\theta} \int_{U \cap\left\{u_{n} \geqslant k+1\right\}} u_{n}^{q} d x d t \\
& \quad \leqslant C \int_{t}^{\theta} \int_{U^{*} \cap\{u \geqslant k\}}|\nabla u|^{p-1} d x d t+\int_{U^{*}} H(u(x, t)) d x \\
& \quad \leqslant C \int_{0}^{\theta} \int_{U^{*} \cap\{u \geqslant k\}}|\nabla u|^{p-1} d x d t+\int_{U^{*}}(u(x, t)-k)^{+} d x .
\end{aligned}
$$

Letting $t$ go to 0 ,

$$
\begin{aligned}
\int_{0}^{\theta} \int_{U \cap\{k \leqslant u \leqslant k+1\}}\left|\nabla u_{n}\right|^{p} d x d t \leqslant & C \int_{0}^{\theta} \int_{U^{*} \cap\left\{u_{n} \geqslant k\right\}}\left|\nabla u_{n}\right|^{p-1} d x d t \\
& +\int_{U^{*}}(\mu-k)^{+} d x
\end{aligned}
$$

Letting $n$ go to $\infty$,

$$
\begin{aligned}
& \int_{0}^{\theta} \int_{U \cap\{k \leqslant u \leqslant k+1\}}|\nabla u|^{p} d x d t \\
& \quad \leqslant C \int_{0}^{\theta} \int_{U^{*} \cap\{u \geqslant k\}}|\nabla u|^{p-1} d x d t+\int_{U^{*}}(\mu-k)^{+} d x
\end{aligned}
$$

Finally, letting $k$ go to $\infty$,

$$
\lim _{k \rightarrow \infty} \int_{0}^{\theta} \int_{U \cap\{k \leqslant u \leqslant k+1\}}|\nabla u|^{p} d x d t=0
$$

### 5.3. The Subcritical Case with a Generalized Borel Measure

The main result of this section is the following.
Theorem 5.4. Let $\Omega=\mathbb{R}^{N}$ and

$$
\begin{equation*}
\max (1, p-1)<q<q_{\mathrm{c}} \tag{5.39}
\end{equation*}
$$

Then for any $v \in \mathscr{B}_{\text {reg }}^{+}\left(\mathbb{R}^{N}\right)$, there exists at least one solution to (5.2), and $u \in C\left(Q_{\infty}\right)$.

Under (5.39) we recall that there exist singular solutions $\tilde{w}_{k}$ of (1.16) with initial data $k \delta_{0}$ for any $k>0$. When $k \rightarrow \infty,\left\{\tilde{w}_{k}\right\}$ increases and converges to $\tilde{w}_{\infty}$, which is a singular solution of (1.1) invariant under the similarity transformations $N_{\ell}$ defined in (4.7). Therefore, it takes the form

$$
\begin{equation*}
\tilde{w}_{\infty}(x, t)=t^{-1 /(q-1)} f\left(t^{-\gamma} x\right) \tag{5.40}
\end{equation*}
$$

where $\gamma=(q+p-1) /(p(q-1))>0$, and $f$ is radial and the unique nontrivial nonnegative solution of the problem

$$
\left\{\begin{array}{l}
r^{1-N}\left(r^{N-1}\left|f^{\prime}\right|^{p-2} f^{\prime}\right)^{\prime}+\gamma r f^{\prime}+\frac{1}{q-1} f-f^{q}=0 \quad \text { in }(0, \infty)  \tag{5.41}\\
f^{\prime}(0)=0 \\
\lim _{r \rightarrow \infty} r^{p /(q+1-p)} f(r)=0
\end{array}\right.
$$

see [24,38] for the case $p>2$, where $f$ has a compact support, and [22] for the case $p<2$. The following result points out the pointwise blow-up over a singular point in the subcritical case.

Lemma 5.5. Assume (5.39) and let $u \in C\left(\mathbb{R}^{N} \times(0, T)\right)$ be a nonnegative weak solution of (1.1) in $\mathbb{R}^{N} \times(0, T)$ with initial trace

$$
\begin{align*}
& \operatorname{tr}_{\mathbb{R}^{N}}(u)=(\mathscr{S}, \mu) \in C M^{+}\left(\mathbb{R}^{N}\right) . \text { If } y \in \mathscr{S} \text {, then } \\
& \qquad u(x, t) \geqslant \tilde{w}_{\infty}(x-y, t), \quad \forall(x, t) \in \mathbb{R}^{N} \times(0, T) . \tag{5.42}
\end{align*}
$$

Proof. The proof is a variant of the one of Theorem 4.1 and is based on a construction due to Marcus and Véron [33,34]. Assuming that $y \in \mathscr{S}$ infers that for any open neighborhood $U$ of $y$,

$$
\lim _{t \rightarrow 0} \int_{U} u(x, t) d x=\infty
$$

Therefore, for any $k>0$ there exist two sequences $\left\{t_{n}\right\}$ and $\left\{r_{n}\right\}$ decreasing to 0 such that

$$
\int_{B_{r_{n}(y)}} u\left(x, t_{n}\right) d x=k .
$$

Then $u \geqslant \tilde{w}_{k, y}$ where $\tilde{w}_{k, y}(x, t)=\tilde{w}_{k}(x-y, t)$ is the fundamental solution of (1.1) with initial data $k \delta_{y}($.$) . Letting k \rightarrow \infty$ implies the claim.

Remark 5.7. If $\mathbb{R}^{N}$ is replaced by a general open subset $\Omega$, the lower estimate on $u$ takes the following form:

$$
u(x, t) \geqslant \tilde{w}_{\infty, R}(x-y, t)
$$

where $\tilde{w}_{\infty, R}$ is the increasing limit as $k \rightarrow \infty$ of the solution of $\tilde{w}=\tilde{w}_{k, R}$

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{w}-\nabla \cdot\left(|\nabla \tilde{w}|^{p-2} \nabla \tilde{w}\right)+\tilde{w}^{q}=0 \quad \text { in } B_{R} \times(0, \infty)  \tag{5.43}\\
\tilde{w}=0 \quad \text { in } \partial B_{R} \times(0, \infty) \\
\tilde{w}(., 0)=k \delta_{0}(.) \quad \text { in } B_{R}
\end{array}\right.
$$

and $R>0$ is chosen in such a way that $\bar{B}_{R}(y) \subset \Omega$.
Proof of Theorem 5.4. Suppose $v=(\mathscr{S}, \mu)$, and let $\left\{a_{k}\right\}_{k \in \mathbb{N}^{*}}$ be a countable dense subset of $\mathscr{S}$. If $n \in \mathbb{N}^{*}$, we define $\mu_{k} \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$ by

$$
\mu_{k}=\mu+k \sum_{j=1}^{k} \delta_{a_{j}}
$$

By Proposition 5.3, there exists a sequence $\left\{u_{k}\right\}$ of constructive continuous solutions of problem (5.8) with initial data $\mu_{k}$, such that

$$
0 \leqslant \tilde{w}_{a_{j}, k} \leqslant u_{k} \leqslant u_{k+1} \leqslant W, \quad \forall k>0 \quad \text { and } j=1, \ldots, k
$$

where $\tilde{w}_{a_{j}, k}$ is the solution of (5.8) with initial data $k \delta_{a_{j}}$. If $k \rightarrow \infty,\left\{u_{k}\right\}$ converges to some function $u$ which satisfies

$$
0 \leqslant \tilde{w}_{a_{j}, \infty} \leqslant u \leqslant W
$$

in $Q_{\infty}$, hence $u$ is locally bounded in $Q_{\infty}$. The techniques developed in Theorem 5.1 on $(0, T)$ apply here on $(\tau, T)$ for any $0<\tau<T$, and infer that $u$ is a weak solution of (1.1) in $Q_{\infty}$. Moreover, as in Remark 5.1, the sequence $\left\{u_{k}\right\}$ is equicontinuous in the local uniform topology of $Q_{\infty}$. Up to a subsequence it converges uniformly on any compact subset of $Q_{\infty}$. Thus $u$ is continuous on $Q_{\infty}$. An easy calculation shows that for any $\rho>0$ and $j \geqslant 1$,

$$
\lim _{t \rightarrow 0} \int_{B_{\rho}\left(a_{j)}\right.} w_{a_{j}, \infty} d x=\infty
$$

Because $\left\{a_{k}\right\}$ is dense in $\mathscr{S}$, the singular set of the initial trace of $u$ contains $\mathscr{S}$. But on the other hand, for any open subsets $V$ $\subset \subset V^{*} \subset \subset \mathscr{R}=\mathbb{R}^{N} \backslash \mathscr{S}$, if we take a test function $\zeta$ with support in $V^{*}$ in the proof of Theorem 5.1, we obtain estimates (5.14) and (2.22) of Lemma 2.2 in $V$ for $u_{k}$. They also hold for $u$, since $\mu_{k}$ and $\mu$ have the same restriction to $\mathscr{R}$. Therefore the regular set of the initial trace of $u$ contains $\bar{V}$. Finally, for any $\theta>0$, letting $k \rightarrow \infty$ in the equality

$$
\begin{aligned}
& \left.\int_{0}^{\theta} \int_{V^{*}}\left(-u_{k} \partial_{t} \varphi+|\nabla u|^{p-2} \nabla u_{k} \cdot \nabla u_{k} \varphi\right)+u_{k}^{q} \varphi\right) d x d t \\
& \quad=\int_{V^{*}} \varphi(x, 0) d \mu_{k}-\int_{V^{*}} u_{k}(x, \theta) \varphi(x, \theta) d x
\end{aligned}
$$

where $\varphi \in C_{\mathrm{c}}^{\infty}\left(V^{*} \times[0, \infty)\right)$, implies (3.5) in $V^{*}$, since $q<q_{\mathrm{c}}$. This proves that the regular part of the initial trace of $u$ is $\mu$ and consequently $\operatorname{Tr}_{\mathbb{R}^{n}}(u)=v \approx(\mathscr{S}, \mu)$.

Remark 5.8. If we endow the set $\mathscr{B}_{+}^{\text {reg }}\left(\mathbb{R}^{N}\right)$ of the following order relation:

$$
v_{1} \tilde{\leqslant} v_{2} \Leftrightarrow\left\{\begin{array}{l}
\mathscr{S}_{1} \subseteq \mathscr{S}_{2} \\
\mu_{1_{\mathscr{S}_{2}}} \leqslant \mu_{2}
\end{array} \quad \text { if } v_{i} \approx\left(\mathscr{S}_{i}, \mu_{i}\right) \quad \text { with } \mathscr{R}_{i}=\mathbb{R}^{N} \backslash \mathscr{S}_{i}\right.
$$

the solutions $u_{i}$ of (5.2) with respective initial trace $v_{i}$ satisfy $u_{1} \leqslant u_{2}$ in $Q_{\infty}$.

### 5.4. The Super-Critical Case

In this section we assume

$$
\begin{equation*}
q \geqslant q_{\mathrm{c}}>1 . \tag{5.44}
\end{equation*}
$$

In that case it is important to notice that neither every measure in an open subset of $\mathbb{R}^{N}$, nor every closed subset of $\mathbb{R}^{N}$ are eligible for being, respectively, the regular part and the singular part of the initial trace of a positive solution of (1.1) in $Q_{\infty}$. The examples of the Dirac measure or the
pointwise singular set show this fact. The conditions for a measure to be eligible should probably be expressed in terms of capacities, but the theory is not known up to now.

Given an open subset $\mathscr{R}$ of $\mathbb{R}^{N}$ and $\varepsilon>0$, we denote

$$
\mathscr{S}=\mathbb{R}^{N} \backslash \mathscr{R}, \quad \mathscr{S}^{\varepsilon}=\left\{x \in \mathbb{R}^{N}: \operatorname{dist}(x, \mathscr{S}) \leqslant \varepsilon\right\}, \quad \mathscr{R}^{\varepsilon}=\mathbb{R}^{N} \backslash \mathscr{S}^{\varepsilon}
$$

If $\mu \in \mathscr{M}^{+}(\mathscr{R})$ we define a measure $\mu_{\varepsilon} \in \mathscr{M}^{+}\left(\mathbb{R}^{N}\right)$ by

$$
\mu_{\varepsilon}(E)=\mu\left(E \cap \mathscr{R}^{\varepsilon}\right), \quad \forall E \subset \mathbb{R}^{N}, E \text { Borel set. }
$$

Definition 5.3. Let $p>1, q \geqslant q_{\mathrm{c}}$ and $\mathscr{R}$ be an open subset of $\mathbb{R}^{N}$. A Radon measure $\mu \in \mathscr{M}^{+}(\mathscr{R})$ is called a $(p, q)$-trace if for any $\varepsilon>0$ there exists a nonnegative constructive solution $u_{\mu_{\varepsilon}}$ of (1.1) in $Q_{\infty}$ with initial trace $\mu_{\varepsilon}$. We denote by $\mathscr{M}^{p, q}(\mathscr{R})$ the set of all $(p, q)$-traces on $\mathscr{R}$.

Remark 5.8. It follows from Proposition 5.3 that if $\mu \in \mathscr{M}^{p, q}(\mathscr{R})$ and $\tilde{\mu} \in \mathscr{M}^{+}(\mathscr{R})$ are such that $\tilde{\mu} \leqslant \mu$, then $\tilde{\mu} \in \mathscr{M}^{p, q}(\mathscr{R})$. Moreover, $\mathscr{M}^{p, q}(\mathscr{R})$ is a positive cone in $\mathscr{M}^{+}(\mathscr{R})$.

Lemma 5.6. Assume $q \geqslant q_{\mathrm{c}}>1$ and let $\mathscr{R}$ be an open subset of $\mathbb{R}^{N}$ and $\mu \in \mathscr{M}^{p, q}(\mathscr{R})$. For $0<\varepsilon^{\prime}<\varepsilon$ there exist constructive solutions of (1.1) corresponding to $\mu_{\varepsilon}$ and $\mu_{\varepsilon^{\prime}}$ satisfying

$$
\begin{equation*}
u_{\mu_{\varepsilon}} \leqslant u_{\mu_{e^{\prime}}} . \tag{5.45}
\end{equation*}
$$

Moreover, $u_{\mu}=\lim _{\varepsilon \rightarrow 0} u_{\mu_{\varepsilon}}$ is a solution of (1.1) in $Q_{\infty}$ with initial data $\mu$ on $\mathscr{R}$.
Proof. Since $\varepsilon^{\prime}<\varepsilon$, we have $\mathscr{R}^{\varepsilon} \subset \mathscr{R}^{\varepsilon^{\prime}}$, hence $\mu_{\varepsilon} \leqslant \mu_{\varepsilon^{\prime}}$, hence by Proposition 5.3, there exist nonnegative constructive corresponding solutions such that

$$
u_{\mu_{\varepsilon}} \leqslant u_{\mu_{e^{\prime}}} \leqslant W
$$

As in Theorem 5.4, they converge a.e. to $u_{\mu}=\sup u_{\mu_{\varepsilon}}$, and $u_{\mu}$ is a weak solution of (1.1) in $Q_{\infty}$, and it converges uniformly on any compact subset of $Q_{\infty}$, hence $u_{\mu}$ is continuous on $Q_{\infty}$. The fact that $u_{\mu}$ has initial data $\mu$ on $\mathscr{R}$ is proved in the following way. Let $\varphi \in C_{\mathrm{c}}^{\infty}(\mathscr{R} \times[0, \infty))$ be nonnegative. For $\varepsilon$ small enough, the support of $\zeta(., t)$ lies in a compact subset of $\mathscr{R}^{\varepsilon}$ independently of $t$. Since $u_{\mu_{\varepsilon}}$ admits $\mu_{\varepsilon}$ as initial trace, we have
for almost all $\theta>0$

$$
\begin{aligned}
\int_{0}^{\theta} \int_{\mathscr{R}} u_{\mu_{\varepsilon}}^{q} \varphi d x d t= & \int_{0}^{\theta} \int_{\mathscr{R}}\left(u_{\mu_{\varepsilon}} \partial_{t} \varphi-\left|\nabla u_{\mu_{\varepsilon}}\right|^{p-2} \nabla u_{\mu_{\varepsilon}} \cdot \nabla \varphi\right) d x d t \\
& +\int_{\mathscr{R}} \varphi(x, 0) d \mu_{\varepsilon}-\int_{\mathscr{R}} \varphi(x, \theta) u_{\mu_{\varepsilon}}(x, \theta) d x .
\end{aligned}
$$

But $\int_{\mathscr{R}} \varphi(x, 0) d \mu_{\varepsilon}=\int_{\mathscr{R}} \varphi(x, 0) d \mu$, and we can pass to the limit in the right-hand side. Then by the Beppo-Levi theorem,

$$
\int_{0}^{\theta} \int_{\mathscr{R}} u_{\mu_{\varepsilon}}^{q} \zeta d x d t \rightarrow \int_{0}^{\theta} \int_{\mathscr{R}} u_{\mu}^{q} \zeta d x d t<\infty
$$

Then

$$
u_{\mu_{e}}^{q} \rightarrow u_{\mu}^{q} \quad \text { strongly in } L_{\mathrm{loc}}^{1}(\mathscr{R} \times[0, \infty))
$$

Hence as in Remark 5.3,

$$
\begin{aligned}
\int_{0}^{\theta} \int_{\mathscr{R}} u_{\mu}^{q} \varphi d x d t= & \int_{0}^{\theta} \int_{\mathscr{R}}\left(u_{\mu} \partial_{t} \varphi-\left|\nabla u_{\mu}\right|^{p-2} \nabla u_{\mu} \cdot \nabla \varphi\right) d x d t \\
& +\int_{\mathscr{R}} \varphi(x, 0) d \mu-\int_{\mathscr{R}} \varphi(x, \theta) u_{\mu}(x, \theta) d x
\end{aligned}
$$

Moreover, for almost all $t>0$, and any $\xi \in C_{\mathrm{c}}^{\infty}(\mathscr{R})$,

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} u_{\mu}(x, t) \xi(x) d x-\int_{\mathbb{R}^{N}} \xi(x) d \mu(x) \\
& \quad=\int_{0}^{t} \int_{\Omega}\left(\left|\nabla u_{\mu}\right|^{p-2} \nabla u_{\mu} \cdot \nabla \xi+u_{\mu}^{q} \xi\right) d x d t
\end{aligned}
$$

by passing to the limit in the corresponding equality for $u_{\mu_{e}}$ which implies

$$
\lim _{t \rightarrow 0} \int_{\mathscr{R}^{\ell}} u_{\mu}(x, t) \xi(x) d x=0
$$

By definition, this means that the solution $u_{\mu}$ admits $\mu$ as initial trace in $\mathscr{R}$.

Lemma 5.7. Assume $q \geqslant q_{\mathrm{c}}>1$ and let $\mathscr{R}$ be an open subset of $\mathbb{R}^{N}$ and $\mathscr{S}=\mathbb{R}^{N} \backslash \mathscr{R}$. For $\varepsilon, k>0$, let $u_{k, \mathscr{S}^{\varepsilon}}$ be the solution of (1.1) in $Q_{\infty}$ with initial data $k \chi_{\mathscr{G}^{8} \cap B_{k}}$. Then $k \mapsto u_{k, \mathscr{G}^{\varepsilon}}$ is increasing and

$$
u_{\mathscr{C}^{\varepsilon}}=\lim _{k \rightarrow \infty} u_{k, \mathscr{G}^{\varepsilon}}
$$

is a solution of $(1.1)$ in $Q_{\infty}$ with initial trace $v_{\mathscr{V}^{e}} \approx\left(\mathscr{S}^{\varepsilon}, 0\right) \in \mathscr{B}_{+}^{\mathrm{reg}}\left(\mathbb{R}^{N}\right)$. Moreover,

$$
\begin{equation*}
\lim _{t \rightarrow 0} t^{1 /(q-1)} u_{\mathscr{S}^{\varepsilon}}(x, t)=\left(\frac{1}{q-1}\right)^{1 /(q-1)}, \quad \forall x \text { interior to } \mathscr{S}^{\varepsilon} \tag{5.46}
\end{equation*}
$$

and this limit holds uniformly on any compact subset interior to $\mathscr{S}^{\varepsilon}$. Moreover, for $0<\varepsilon^{\prime}<\varepsilon$ there holds

$$
\begin{equation*}
u_{\mathscr{P}^{e^{\prime}}} \leqslant u_{\mathscr{C}^{\mathrm{e}}} \tag{5.47}
\end{equation*}
$$

Proof. The existence and uniqueness of $u_{k, \mathscr{S}^{e}}$ follows from the classical theory for the Cauchy problem for Eq. (1.1), and from the maximum principle, $k \mapsto u_{k, \mathscr{G}^{e}}$ is increasing. Since there always holds

$$
\begin{equation*}
u_{k, \mathscr{Y}^{\varepsilon}}(x, t) \leqslant W(t) \tag{5.48}
\end{equation*}
$$

in $Q_{\infty}$, then $u_{\mathscr{C}^{\varepsilon}}=\lim _{k \rightarrow \infty} u_{k, \mathscr{Y}^{\varepsilon}}$ exists. As above, it is a weak solution of (1.1) in $Q_{\infty}$. Moreover, inequality (5.47) holds clearly for $0<\varepsilon^{\prime}<\varepsilon$ as a consequence of the approximation process.

In order to prove (5.46), we consider a ball $B_{r}$. For $\sigma>0$ small enough the functional $H_{\sigma}$ defined on $W_{0}^{1, p}\left(B_{r}\right) \cap L^{q+1}\left(B_{r}\right)$ by

$$
H_{\sigma}(\varphi)=\int_{B_{r}}\left(\frac{\sigma}{p}|\nabla \varphi|^{p}+\frac{1}{q+1}|\varphi|^{q+1}-\frac{1}{2} \varphi^{2}\right) d x
$$

achieves a negative minimal value. Let $\Phi_{r, \sigma}$ be a positive minimizer, solution of the problem

$$
\left\{\begin{array}{l}
-\sigma \nabla \cdot\left(\left|\nabla \Phi_{r, \sigma}\right|^{p-2} \nabla \Phi_{r, \sigma}\right)+\Phi_{r, \sigma}^{q}=\Phi_{r, \sigma} \quad \text { in } B_{r},  \tag{5.49}\\
\Phi_{r, \sigma}=0 \quad \text { on } \partial B_{r} .
\end{array}\right.
$$

From the maximum principle

$$
0<\Phi_{r, \sigma} \leqslant 1 \quad \text { in } B_{r}
$$

Put $w(x, t)=W(t) \Phi_{r, \sigma}$. Then

$$
\begin{aligned}
\partial_{t} w & -\nabla \cdot\left(|\nabla w|^{p-2} \nabla w\right)+w^{q} \\
& =-W^{p-1} \nabla \cdot\left(\left|\nabla \Phi_{r, \sigma}\right|^{p-2} \nabla \Phi_{r, \sigma}\right)+U^{q}\left(\Phi_{r, \sigma}^{q}-\Phi_{r, \sigma}\right) \\
& =\left(W^{q}-\frac{1}{\sigma} W^{p-1}\right)\left(\Phi_{r, \sigma}^{q}-\Phi_{r, \sigma}\right)
\end{aligned}
$$

Set $t_{\sigma}=(q-1) \sigma^{(q-1) /(q+1-p)}$. For $0<t \leqslant t_{\sigma}$, we have

$$
W(t)^{q}-\frac{1}{\sigma} W(t)^{p-1} \geqslant 0
$$

Since $\Phi_{r, \sigma}^{q}-\Phi_{r, \sigma} \leqslant 0$, the function $w$ is a subsolution. Now consider $b \in \mathscr{S}_{\varepsilon}$ such that the ball $B_{r}(b)$ be compactly imbedded in $\mathscr{S}_{\varepsilon}$. Replacing $t$ by $t+\delta$ in the definition of $w$, and $B_{r}$ by a ball $B_{r}(b)$, and letting $\delta$ go to 0 , yields to the lower estimate

$$
\begin{equation*}
u_{\mathscr{S}^{\varepsilon}}(x, t) \geqslant W(t) \Phi_{r, \sigma}(x-b) \quad \text { in } B_{r}(b) \times\left(0, t_{\sigma}\right] \tag{5.50}
\end{equation*}
$$

Using the scaling invariance of the equation, the transformed function $N_{\ell}\left(u_{\mathscr{C}^{\varepsilon}}\right)$ defined by (4.7) satisfies (1.1), and

$$
\begin{equation*}
W(t) \Phi_{r, \sigma}\left(\ell^{\beta}(x-b)\right) \leqslant N_{\ell}\left(u_{\mathscr{S}^{\varepsilon}}\right)(x, t) \leqslant W(t) \tag{5.51}
\end{equation*}
$$

in $B_{r \ell^{-\beta}}(b) \times\left(0, t_{\sigma} \ell^{-1}\right]$. By the previous estimates and the local regularity theory, see Remark 5.3, there exists a sequence $\left\{\ell_{n}\right\}$ converging to 0 such that $N_{\ell_{n}}\left(u_{g^{\varepsilon}}\right)$ converges to a function $U$ solution of (1.1) in $Q_{\infty}$. Since $\Phi_{r, \sigma}$ $\geqslant \theta>0$ on $B_{r / 2}$ for some $\theta>0$, the function $U$ satisfies

$$
\begin{equation*}
\theta W(t) \leqslant U(x, t) \leqslant W(t) \tag{5.52}
\end{equation*}
$$

in $Q_{\infty}$. Because $W(t+\lambda)<U(x, t) \leqslant W(t)$ for any $t>0$ and $\lambda>0$ it follows classically $U=W$ by letting $\lambda$ go to 0 . This equality implies also

$$
\lim _{\ell \rightarrow 0} N_{\ell}\left(u_{\mathscr{S}^{\varepsilon}}\right)=W .
$$

Taking $t=1$ and $x=b$ yields to

$$
\lim _{t \rightarrow 0} t^{1 /(q-1)} u_{\mathscr{C}^{\varepsilon}}(b, t)=\left(\frac{1}{q-1}\right)^{1 /(q-1)}
$$

and this holds for any $b$ interior to $\mathscr{S}^{2}$. Since the equation is invariant by $x$-translations, the uniformity on any compact interior to $\mathscr{S}^{\varepsilon}$ follows easily by contradiction.

Lemma 5.8. Under the assumptions of Lemmas 5.6 and 5.7, let $\partial_{\mu} \mathscr{S}$ be the singular set of the initial trace of $u_{\mu}$, and

$$
u_{\mathscr{I}}=\lim _{\varepsilon \rightarrow 0} u_{\mathscr{S}^{\varepsilon}} .
$$

Then $u_{\mathscr{S}}$ is a solution of (1.1) in $Q_{\infty}$. If we denote by $\mathscr{S}_{p, q}^{*}$ the singular set of its initial trace, there always holds

$$
\mathscr{S}_{p, q}^{*} \cup \partial_{\mu} \mathscr{S} \subset \mathscr{S}
$$

Proof. Since the initial trace $u_{\mu}$ in $\mathscr{R}$ is equal to $\mu$, we have $\partial_{\mu} \mathscr{S} \subset \mathscr{S}$. Since $\varepsilon \mapsto u_{\mathscr{C}^{\varepsilon}}$ is increasing, $u_{\mathscr{C}}=\lim _{\varepsilon \rightarrow 0} u_{\mathscr{C}^{\varepsilon}}$ exists and is a solution of (1.1) in $Q_{\infty}$. Since $\mathscr{R}^{\varepsilon}$ is the regular set of the initial trace of $u_{\mathscr{P}^{\varepsilon}}$, we deduce that for any $\varepsilon>0$,

$$
\begin{aligned}
& \left.\int_{0}^{\theta} \int_{\mathscr{R}^{\varepsilon}}\left(-u_{\mathscr{S}} \partial_{t} \varphi+\left|\nabla u_{\mathscr{S}}\right|^{p-2} \nabla u_{\mathscr{S}} \cdot \nabla u_{\mathscr{S}} \varphi\right)+u_{\mathscr{S}}^{q} \varphi\right) d x d t \\
& \quad=-\int_{\mathscr{R}^{\varepsilon}} u_{\mathscr{C}}(x, \theta) \varphi(x, \theta) d x
\end{aligned}
$$

for any $\theta>0$ and $\varphi \in C_{\mathrm{c}}^{\infty}\left(\mathscr{R}^{\varepsilon} \times[0, \infty)\right)$, by passing to the limit in the corresponding relation for $u_{\mathscr{C}^{\varepsilon}}$. Similarly we get, for almost all $t>0$ and any $\xi \in C_{\mathrm{c}}^{\infty}\left(\mathscr{R}^{\varepsilon}\right)$,

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} u_{\mathscr{L}}(x, t) \xi(x) d x-\int_{\mathbb{R}^{N}} \xi(x) d \mu(x) \\
& \quad=\int_{0}^{t} \int_{\Omega}\left(\left|\nabla u_{\mathscr{S}}\right|^{p-2} \nabla u_{\mathscr{S}} \nabla \xi+u_{\mathscr{S}}^{q} \xi\right) d x d t
\end{aligned}
$$

which implies

$$
\lim _{t \rightarrow 0} \int_{\mathscr{R}^{R}} u_{\mathscr{S}}(x, t) \xi(x) d x=0 .
$$

Then $u_{\mathscr{S}}$ admits 0 as an initial trace in $\mathscr{R}^{\varepsilon}$. Hence $\mathscr{R}^{\varepsilon} \subset \mathbb{R}^{N} \backslash \mathscr{S}_{p, q}^{*}$, for any $\varepsilon>0$, hence $\mathscr{R} \subset \mathbb{R}^{N} \backslash \mathscr{S}_{p, q}^{*}$, i.e. $\mathscr{S}_{p, q}^{*} \subset \mathscr{S}$.

The main result in this section is the following.
TheOrem 5.9. Assume $q \geqslant q_{\mathrm{c}}>1$ and let $v \approx(\mathscr{S}, \mu) \subset \mathscr{B}_{+}^{\mathrm{reg}}\left(\mathbb{R}^{N}\right)$. Then a sufficient condition for the existence of a nonnegative solution $u$ of (1.1) in $Q_{\infty}$ with initial trace $v$, is the following:
(i) $\mu \in \mathscr{M}^{p, q}(\mathscr{R})$.
(ii) There holds

$$
\begin{equation*}
\mathscr{S}_{p, q}^{*} \cup \partial_{\mu} \mathscr{S}=\mathscr{S} . \tag{5.53}
\end{equation*}
$$

Proof. (i) We first assume that $p>2$. For $\varepsilon>\eta>0$ and $k>0$ we consider the functions $u_{\mu_{\varepsilon}}$ and $u_{k, \mathscr{G}^{\eta}}$ introduced in Lemmas 5.6 and 5.7. The support of $u_{\mu_{\varepsilon}}(., 0)$ and $u_{k, \mathscr{G}^{n}}(., 0)$ are disjoint. Since $p>2$ the speed of propagation of the support of $u_{\mu_{\varepsilon}}$ and $u_{k, \mathscr{S}^{\eta}}$ is finite and depends locally on the amount of mass concentrated near the free boundary [32]. Consequently, there exists $t_{\varepsilon, \eta, k}>0$ such that, for $0<t<t_{\varepsilon, \eta, k}$ the support of $u_{\mu_{\varepsilon}}(., t)$ and $u_{k, \mathscr{S}^{\eta}}(., t)$ are disjoint. It implies that $u_{\mu_{\varepsilon}}+u_{k, \mathscr{S}^{\eta}}$ is a solution of (5.8) on $Q_{t_{\varepsilon, n, k}}=\mathbb{R}^{N} \times\left(0, t_{\varepsilon, \eta, k}\right)$ with initial data $\mu_{\varepsilon}+k \chi_{\mathscr{S}^{\eta} \cap B_{k}}$. Therefore, we can
define a solution $u_{\mu_{e}, k, \mathscr{S}^{\eta}}$ by

$$
u_{\mu_{\varepsilon}, k, \mathscr{S}_{\eta}}(x, t)= \begin{cases}u_{\mu_{\varepsilon}}(x, t)+u_{k, \mathscr{S}^{\eta}}(x, t) & \text { in } \mathbb{R}^{N} \times\left(0, t_{\varepsilon, \eta, k}\right),  \tag{5.54}\\ v_{k, \varepsilon, \eta}\left(x, t-t_{\varepsilon, \eta, k}\right) & \text { in } \mathbb{R}^{N} \times\left(t_{\varepsilon, \eta, k}, \infty\right),\end{cases}
$$

where $v_{k, \varepsilon, \eta}$ is a solution of (5.8) in $Q_{\infty}$ with initial data $u_{\mu_{\varepsilon}}\left(., t_{\varepsilon, \eta, k}\right)+u_{k, \mathscr{G}^{\eta}}\left(., t_{\varepsilon, \eta, k}\right)$. By Theorem 5.1 it exists since the initial data is locally integrable. Moreover, since $u_{\mu_{\varepsilon}}$ and $u_{k, \mathscr{S}^{\eta}}$ are constructive, we can construct $v_{k, \varepsilon, \eta}$ such that

$$
v_{k, \varepsilon, \eta} \geqslant \sup \left(u_{\mu_{\varepsilon}}, u_{k, \mathscr{S}^{\eta}}\right) \quad \text { in } \mathbb{R}^{N} \times\left(t_{\varepsilon, \eta, k}, \infty\right),
$$

hence

$$
u_{\mu_{e}, k, \mathscr{S}^{\eta}} \geqslant \sup \left(u_{\mu_{\varepsilon}}, u_{k, \mathscr{S}^{\eta}}\right) \quad \text { in } Q_{\infty} .
$$

Letting $k \rightarrow \infty$ infers that $u_{\mu_{e}, k, \mathscr{S}^{n}}$ increases and converges to a solution $u_{\mu_{\varepsilon}, \mathscr{Y}^{\eta}}$ of (1.1). Therefore,

$$
u_{\mu_{\varepsilon}, \mathscr{Y}^{\eta}} \geqslant \sup \left(u_{\mu_{\varepsilon}}, u_{\mathscr{S}^{\eta}}\right) \quad \text { in } Q_{\infty} .
$$

Because of Lemmas 5.6 and 5.7 the function $u_{\mu_{\varepsilon}, k, \mathscr{P}^{\eta}}$ is monotone decreasing with respect to $\varepsilon$ and increasing with respect to $k$ and $\eta$ (such are the initial data). Therefore $u_{\mu_{\varepsilon}, \mathscr{S}^{\eta}}$ is monotone decreasing with respect to $\varepsilon$ and increasing with respect to $\eta$. We let successively $\eta$ and $\varepsilon$ go to 0 . Then there exists

$$
u_{\mu, \mathscr{S}}=\lim _{\varepsilon \rightarrow 0} \lim _{\eta \rightarrow 0} u_{\mu_{\varepsilon}, \mathscr{S}^{\eta}}
$$

As in the previous limit process, $u_{\mu, \mathscr{S}}$ is solution of (1.1) in $Q_{\infty}$, and

$$
\sup \left(u_{\mu}, u_{\mathscr{S}}\right) \leqslant u_{\mu, \mathscr{S}} \leqslant W
$$

If we call $\tilde{\mathscr{S}}$ the singular part of the initial trace of $u_{\mu, \mathscr{\mathscr { L }}}$, then by definition of a singular set,

$$
\mathscr{S}_{p, q}^{*} \cup \partial_{\mu} \mathscr{S}=\mathscr{S} \subset \tilde{\mathscr{S}} .
$$

But, for any $b \in \mathscr{R}$, if $0<\varepsilon_{0}<\operatorname{dist}(b, \mathscr{S}) / 3$, then for any $T>0, \int_{B_{\varepsilon_{0}}(b)} u_{\mu_{\varepsilon}, \mathscr{S}^{\eta}} \times$ $(x, \theta) d x$ remains bounded for $\theta \in(0, T]$ independently of $0<\eta<\varepsilon<\varepsilon_{0}$. Let $\zeta \in C_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{N}\right)$ with support in $\mathscr{R}^{\varepsilon_{0}}$ such that $\zeta=1$ in $B_{\varepsilon_{0}}(b)$, and $\tau>0$ large enough. By estimate (5.11), we derive

$$
\frac{1}{8} \int_{\Omega} u_{\mu_{\varepsilon}, \mathscr{G}^{\eta}}(x, \theta) \zeta^{\tau}(x) d x+\int_{0}^{\theta} \int_{\Omega} u_{\mu_{\varepsilon}, \mathscr{G}^{\eta}}^{q} \zeta^{\tau} d x d t \leqslant \int_{\Omega} \mu(x) \zeta^{\tau}(x) d x+C
$$

with $C=C=C(p, q, N, T, \zeta, \tau)$. Consequently, the same holds for $\int_{B_{\varepsilon_{0}}(b)} u_{\mu, \mathscr{\mathscr { C }}}(x, \theta) d x$, and this implies $B_{\varepsilon_{0}}(b) \subset \tilde{\mathscr{R}}=\mathbb{R}^{N} \backslash \tilde{\mathscr{S}}$. By complement $\tilde{\mathscr{S}}$ $\subset{ }^{{ }^{\circ} \mathscr{S}}$ and finally

$$
\mathscr{S}=\tilde{\mathscr{S}} .
$$

For the regular part of the initial trace of $u_{\mu, \mathscr{S}}$, we prove that it coincides with $\mu$ in $\mathscr{R}$, as in Lemma 5.6. It follows that the initial trace of $u_{\mu, \mathscr{\mathscr { C }}}$ is the couple ( $\mathscr{R}, \mu$ ).
(ii) Now assume that $1<p<2$. Then the speed of propagation of the support of the solutions to (1.1) is infinite. However, for any $\kappa<p-1$ and $\gamma>0$, the following equation:

$$
\begin{equation*}
\partial v-\nabla \cdot\left(|\nabla v|^{p-2} \nabla v\right)+v^{q}+\gamma v^{k}=0 \tag{5.55}
\end{equation*}
$$

has the finite speed of propagation property, [25]. Moreover the positive classical solutions depends monotonically of $\gamma$. Therefore, we follow the construction of the case $p>2$ in constructing first the functions $u_{\mu_{\varepsilon}}^{\gamma}$ and $u_{k, S n}^{\eta}$ solutions of (5.55) with respective initial data $\mu_{\varepsilon}$ and $k \chi_{\text {Sn }^{\eta} B_{k}}$. Those solutions exist and are dominated by $u \gamma_{\mu_{\varepsilon}}$ and $u \gamma_{k, S^{n}}$, respectively. We define $u_{\mu_{c}, k, \mathscr{S}^{n}}^{\gamma}$ by additivity in a similar way as in (5.54). Now $u_{\mu, \mathscr{S}}$ is well defined by the expression

$$
u_{\mu, \mathscr{G}}=\lim _{\varepsilon \rightarrow 0} \lim _{\eta \rightarrow 0} \lim _{k \rightarrow \infty} \lim _{\gamma \rightarrow 0} u_{\mu_{c}, k, \mathscr{S}^{\prime \prime}}^{\gamma} .
$$

The remaining of the proof is as in the first case.
Remark 5.6. In the case $p=2$, it is proved in [33] that the two assertions (i) and (ii) in Theorem 5.9 are necessary and sufficient conditions for the existence of a solution with initial trace $v \approx(\mathscr{S}, \mu)$. Moreover, the solution which is constructed is a maximal solution. Finally, conditions (i) and (ii) are expressed in terms of $N$-dimensional Bessel capacities $C_{2 / q, q}$ :

$$
\mu \in \mathscr{M}^{2, q}(\mathscr{R}) \Leftrightarrow \mu(E)=0, \quad \forall \text { Borel set } E \text { with } C_{2 / q, q^{\prime}}(E)=0 .
$$

And the set $\mathscr{S}_{p, q}^{*}=\mathscr{S}_{q}^{*}$ is the "nonremovable part" of $\mathscr{S}$, and more precisely,

$$
\mathscr{S}_{q}^{*}=\left\{x \in \mathbb{R}^{N}: C_{2 / q, q^{\prime}}(\mathscr{S} \cap U)>0, \forall U \text { open neighborhood of } x\right\},
$$

and $\partial_{\mu} \mathscr{S}$ is the set of blow-up points of $\mu$, that is

$$
\partial_{\mu} \mathscr{S}=\{x \in \mathscr{S}: \mu(U \cap \mathscr{S})=\infty, \forall U \text { open neighborhood of } x\} .
$$

## APPENDIX

We prove below Lemma 5.2. We assume that $\left\{\mu_{n}\right\} \subset C_{\mathrm{c}}\left(\mathbb{R}^{N}\right)$ are nonnegative and converges to $\mu$ in the sense of measures. Since $\tilde{\mu} \leqslant \mu$, the Radon-Nikodym theorem applies and there exists a function $g \in L_{\mu}^{1}\left(\mathbb{R}^{N}\right)$ such that

$$
d \tilde{\mu}=g d \mu
$$

Moreover,

$$
0 \leqslant g \leqslant 1, \quad \mu \text {-a.e. in } \mathbb{R}^{N}
$$

Now there exists a sequence $\left\{g_{n}\right\} \subset C_{\mathrm{c}}\left(\mathbb{R}^{N}\right)$ of nonnegative functions, such that $g_{n} \rightarrow g$ in $L_{\mu}^{1}\left(\mathbb{R}^{N}\right)$. By truncation it can also be assumed that $0 \leqslant g_{n} \leqslant 1, \forall n \in \mathbb{N}$. From the assumption, for any $\varphi \in C_{\mathrm{c}}\left(\mathbb{R}^{N}\right)$ and any $k \in$ $\mathbb{N}^{*}$ there exists $n_{k} \in \mathbb{N}^{*}$ such that $\forall n \geqslant n_{k}$

$$
\left|\int_{\mathbb{R}^{N}} \varphi g_{k} \mu_{n} d x-\int_{\mathbb{R}^{N}} \varphi g_{k} d \mu\right| \leqslant \frac{1}{k}
$$

We define a new sequence $\left\{g_{n}^{\prime}\right\}$ by setting

$$
g_{n}^{\prime}=g_{k} \quad \text { if } n_{k} \leqslant n<n_{k+1}
$$

Then

$$
\begin{aligned}
& \left|\int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} \mu_{n} d x-\int_{\mathbb{R}^{N}} \varphi d \tilde{\mu}\right| \\
& \quad \leqslant\left|\int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} \mu_{n} d x-\int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} d \mu\right|+\left|\int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} d \mu-\int_{\mathbb{R}^{N}} \varphi g d \mu\right|
\end{aligned}
$$

and, for $n \geqslant n_{k}$,

$$
\left|\int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} \mu_{n} d x-\int_{\mathbb{R}^{N}} \varphi d \tilde{\mu}\right| \leqslant \frac{1}{k}+\left|\int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} d \mu-\int_{\mathbb{R}^{N}} \varphi g d \mu\right| .
$$

Since, by Lebesgue's theorem,

$$
\lim _{n^{\prime} \rightarrow \infty} \int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} d \mu-\int_{\mathbb{R}^{N}} \varphi g d \mu=0
$$

we conclude that for any continuous function $\varphi$ with compact support, there exists a subsequence $\left\{g_{n}^{\prime}\right\}$ of $\left\{g_{n}\right\}$ such that

$$
\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N}} \varphi g_{n}^{\prime} \mu_{n} d x-\int_{\mathbb{R}^{N}} \varphi d \tilde{\mu}=0
$$

Since the set $C_{\mathrm{c}}\left(\mathbb{R}^{N}\right)$ admits a countable dense subset $\left\{\varphi_{j}\right\}_{j \in \mathbb{N}}$ for the topology of uniform convergence on compact subsets of $\mathbb{R}^{N}$ (which is defined by the countable set of semi-norms $\|\eta\|_{m}=\max _{|x| \leqslant m}|\eta(x)|$, for $m \in$ $\mathbb{N}_{*}$ and $\eta \in C_{\mathbf{c}}\left(\mathbb{R}^{N}\right)$ ), there exists a diagonal sequence $\left\{\tilde{\mu}_{n}\right\}$ extracted from $\left\{g_{n}^{\prime} \mu_{n}\right\}$ such that

$$
\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N}} \varphi_{j} \tilde{\mu}_{n} d x=\int_{\mathbb{R}^{N}} \varphi_{j} d \tilde{\mu}, \quad \forall j \in \mathbb{N}
$$

The conclusion follows by density.

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